



STELLAR AFTERLIFE

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SIRIUS

One of the most famous and influential binary star systems is **Sirius A & B**.

As the brightest star in the sky, Sirius A has been known since antiquity.

The **gravitational effect** of Sirius B on Sirius A was noted by Bessel in 1844, and Sirius B was observed in 1862 by Clark.

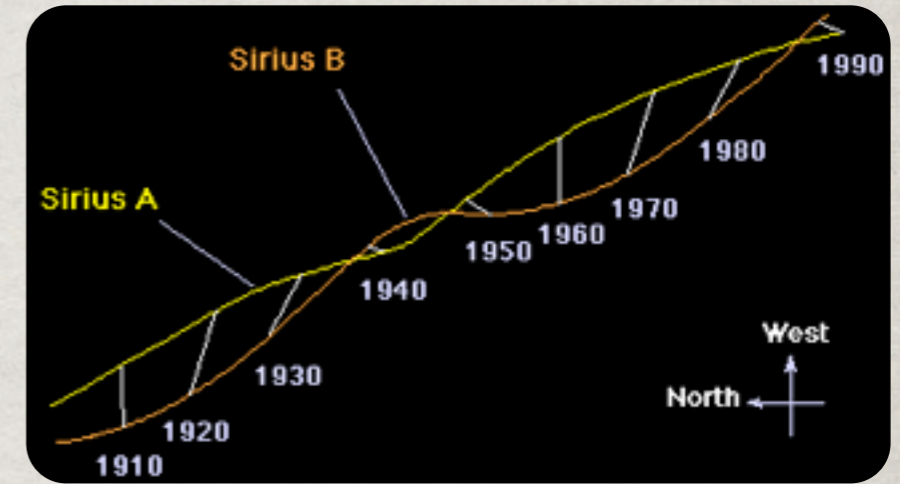
With a much smaller luminosity ($L_B = 0.024 L_\odot$) than Sirius A ($L_A = 26 L_\odot$), it was very surprising when a spectrum of Sirius B taken in 1915, revealed a color similar to Sirius A. Sirius B was the 2nd **white dwarf** identified.

Using modern $T_A = 9900$ K & $T_B = 24800$ K.

$$\frac{R_B}{R_A} = \left(\frac{L_B}{L_A} \right)^{\frac{1}{2}} \left(\frac{T_A}{T_B} \right)^2 = 0.005$$

$$\text{For } R_A = 1.71 R_\odot$$

$$R_B = 0.008 R_\odot = 0.9 R_\oplus$$



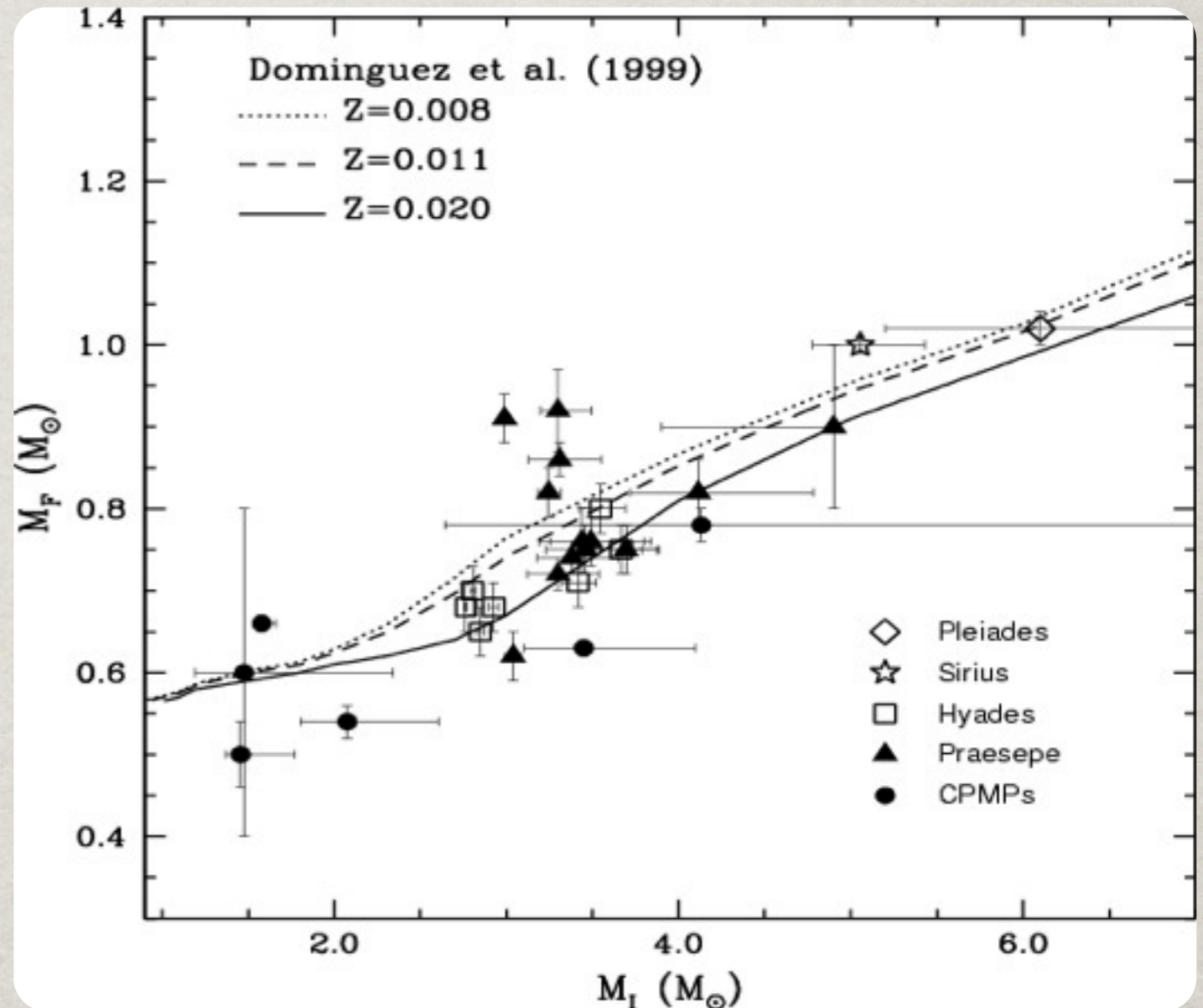
MAKING WHITE DWARVES

The final white dwarf product of a star's evolution depends on the star's **mass**, as well as its metallicity and solitude.

In general, more massive stars are **less efficient** at growing white dwarves.

A star of 7-8 M_{\odot} leaves a **C-O** white dwarf $\sim 1.1 M_{\odot}$.

Any **larger white dwarf** is either composed of O-Ne or gained more mass after it formed.



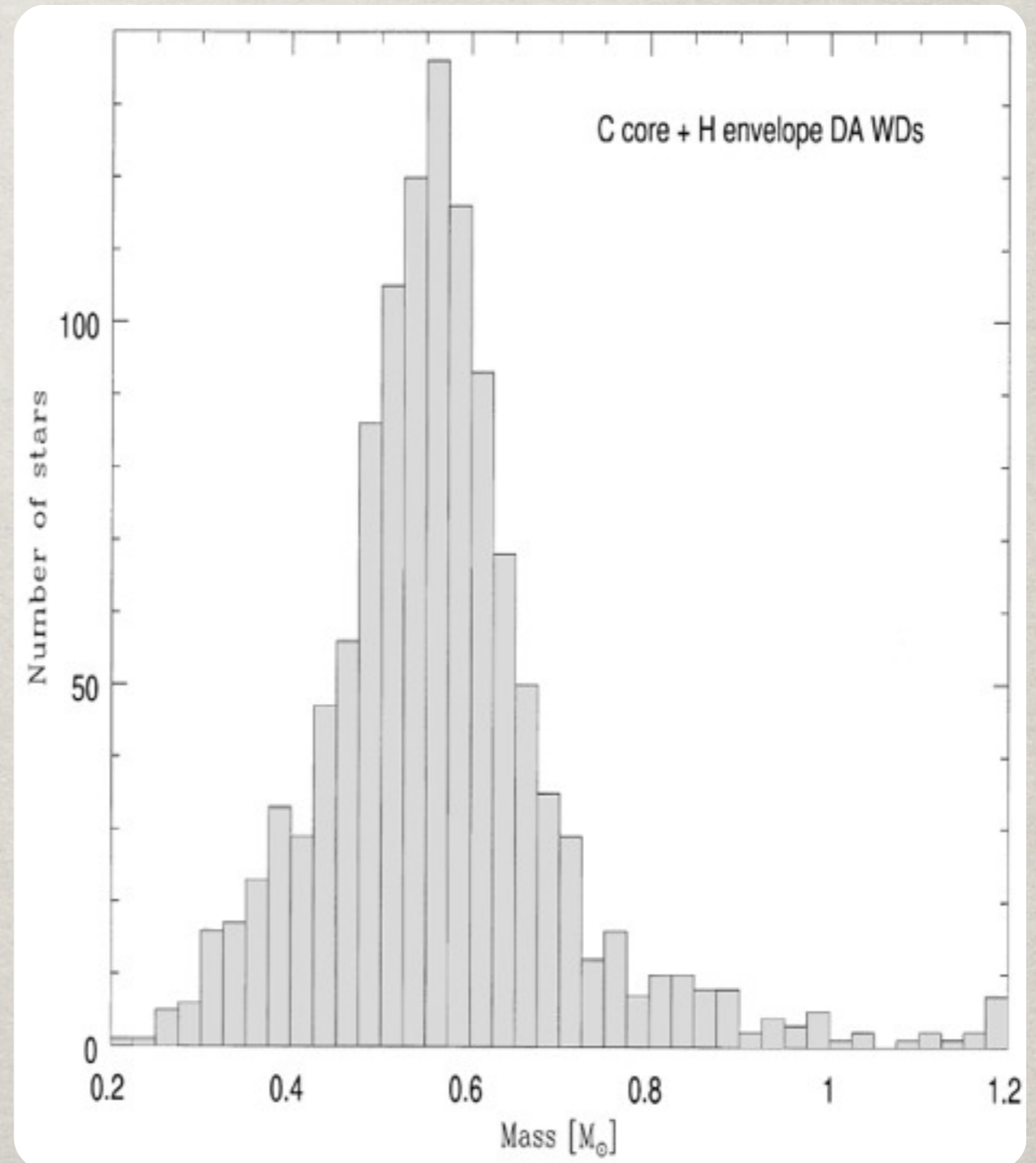
WHITE DWARF ACCOUNTING

Stellar censuses which account for the dimness of white dwarves find that typically **10% of all stars** are white dwarves.

The current distribution of white dwarf masses involves a competition between the **higher frequency** and **longer lifetimes** of lower mass main sequence stars.

It also depends on changes due to **binary evolution**.

The observed distribution shows a **peak around $0.6 M_{\odot}$** .



DEGENERACY PRESSURE

Thermal electrons exert a pressure

$$P_{th} = n_e kT \sim n_e m_e v_{th}^2$$

A similar pressure can be computed from the “Heisenberg speed”

$$P_{deg} \sim n_e m_e v_H^2 \sim n_e m_e \left(\frac{\hbar n_e^{1/3}}{m_e} \right)^2 \sim \hbar^2 \frac{n_e^{5/3}}{m_e}$$

It is instructive to rewrite this degeneracy pressure by separating out a **proportionality** to the number of electrons contributing, as is clearly present in the thermal pressure.

$$P_{deg} \sim n_e \frac{\hbar^2}{m_e} n_e^{2/3} \sim n_e E_f$$

Thus $P_{deg} > P_{th}$ as long as the gas is **degenerate** ($E_f > kT$).

MASS-RADIUS

Degeneracy pressure leads to an unusual relation between mass and radius.

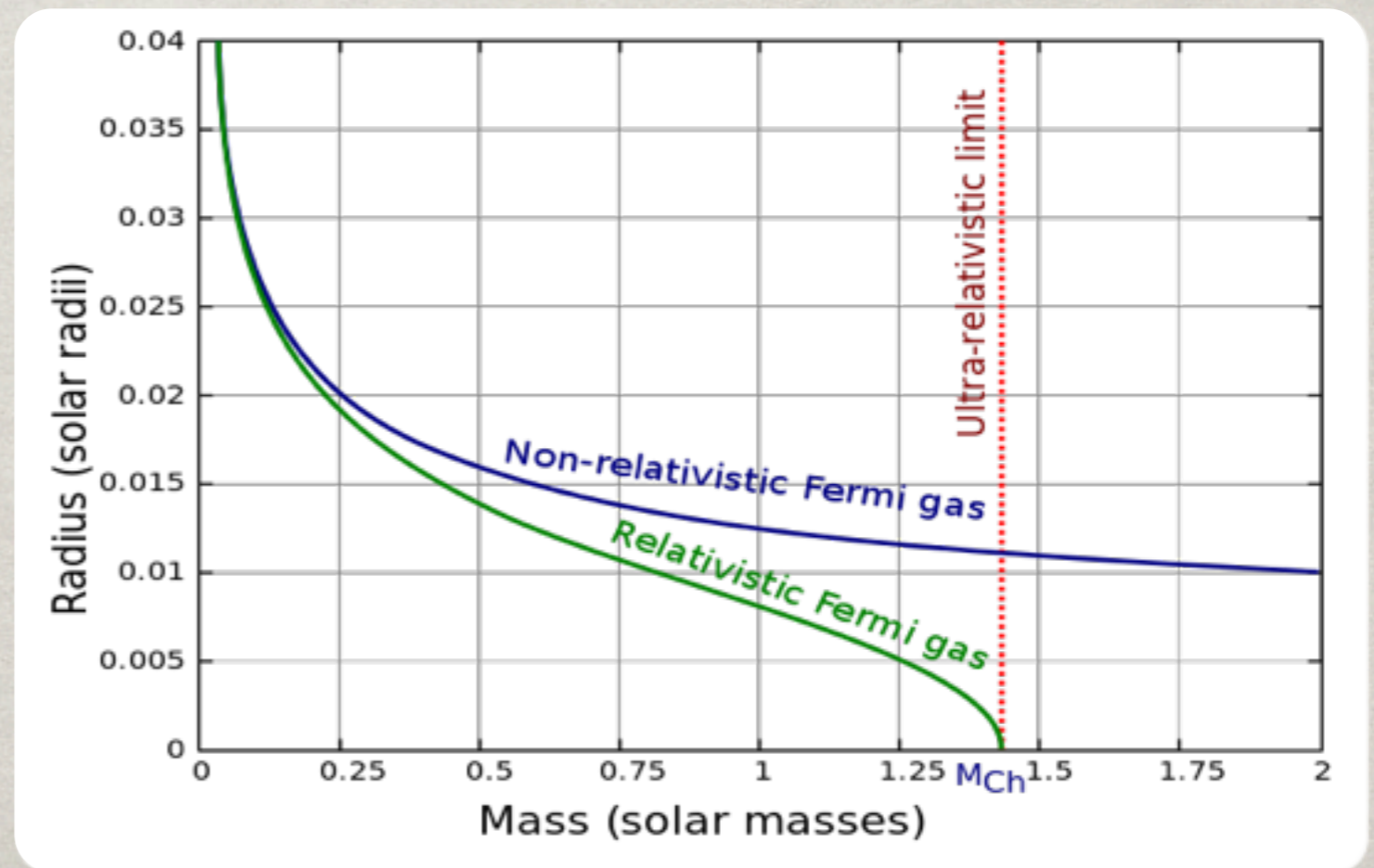
$$R \sim \frac{\hbar^2}{Gm_e m_p^{5/3}} M^{-1/3} \approx 0.01 R_\odot \left(\frac{M}{0.7 M_\odot} \right)^{-1/3}$$

Subrahmanyan Chandrasekhar derived this relation in 1931 for the more general case of special relativity and a star whose density varies with radius, finding the star's radius goes to zero at a limiting mass of $M_{Ch} = 5.76 Y_e^2 M_\odot$.

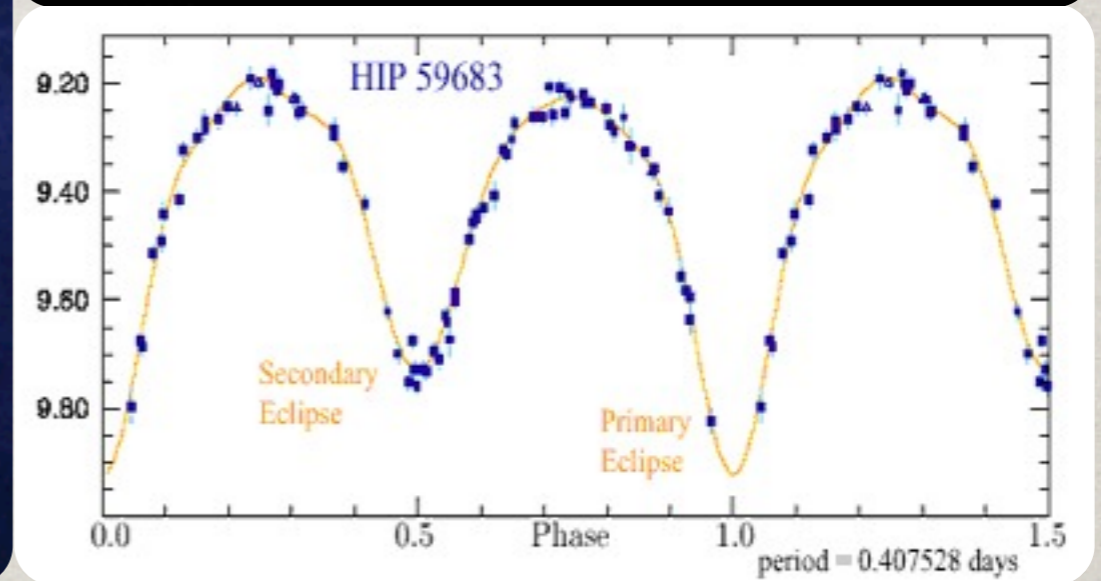
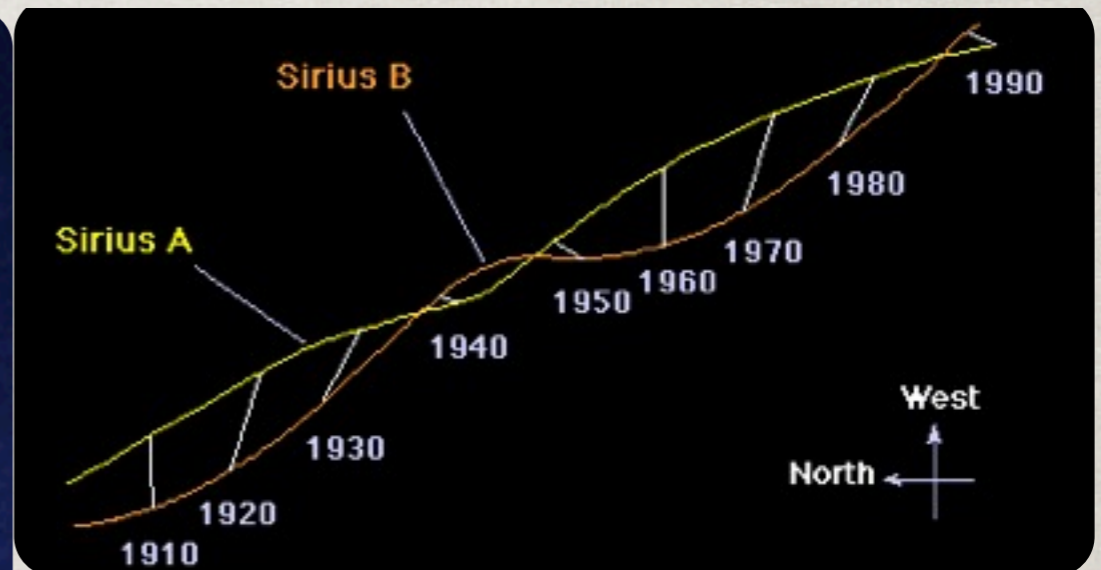
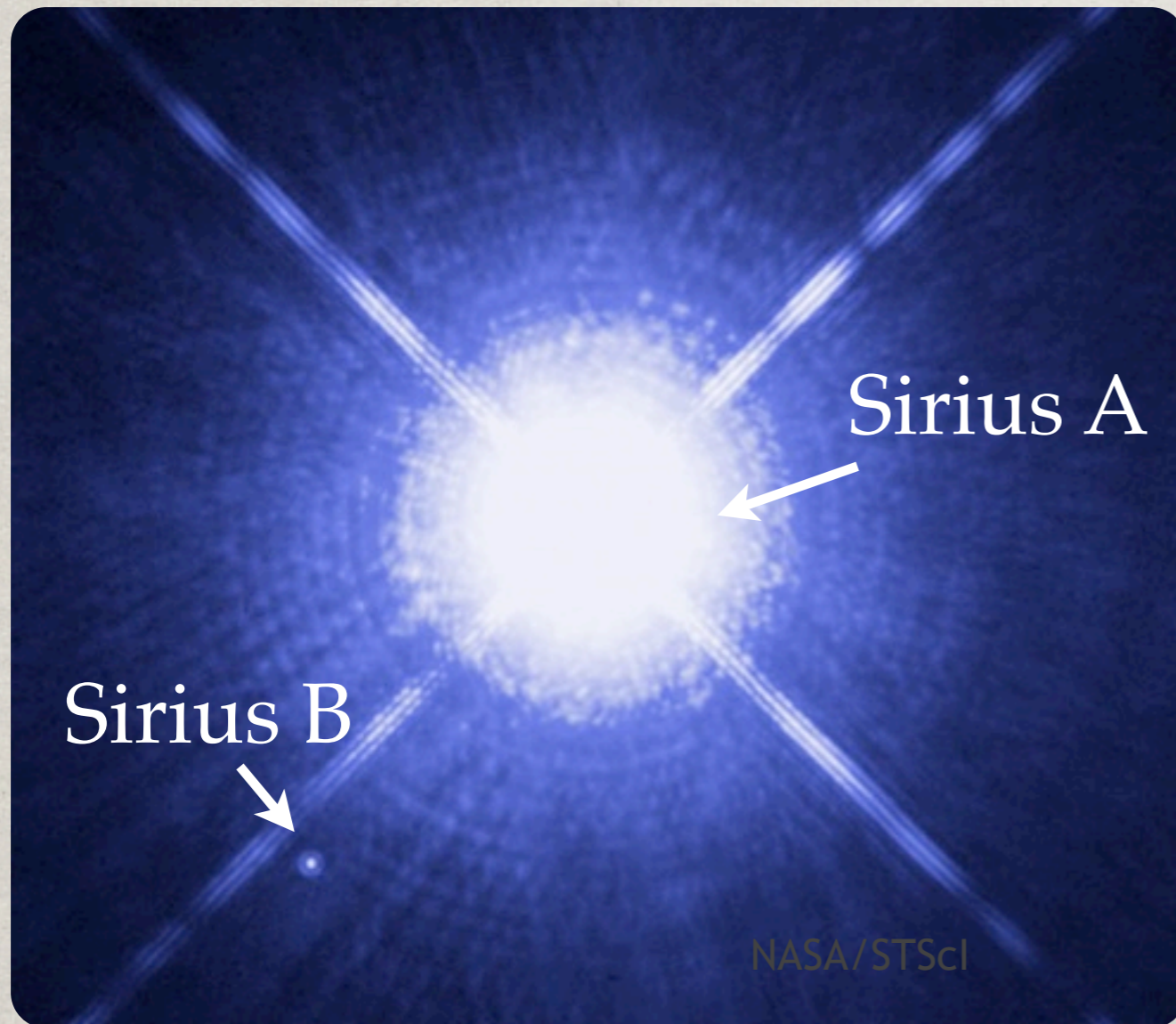
Y_e is the electron fraction, the number of electrons per nucleon, $Y_e = \langle Z/A \rangle$.

For pure C+O, $Y_e = .5$.

A small amount of ^{22}Ne present makes $Y_e = .498$.



STARS DON'T ALWAYS LIVE ALONE



Many (most?) stars are members of binary systems.

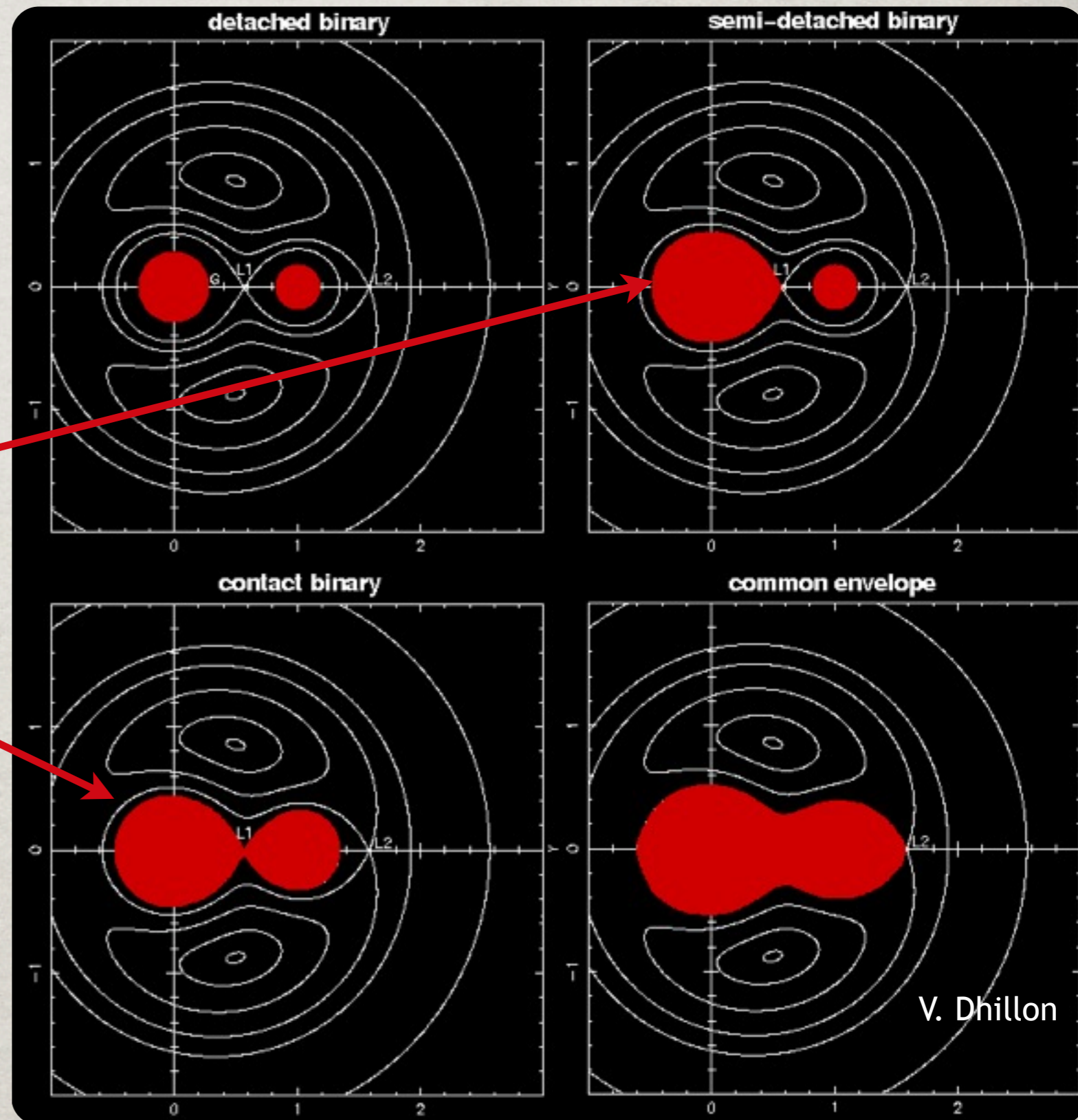
Binaries can be identified by their orbits effect on the **position** (astrometric), **light curve (eclipsing)** or **spectra** (spectroscopic) of the stars.

EQUIPOTENTIALS

Presence of companion star results in **non-spherical equipotential surfaces**.

When a giant star expands, it can fill its **Roche Lobe** and preferentially lose mass to the companion.

This mass, with significant **angular momentum**, accretes on to the companion.



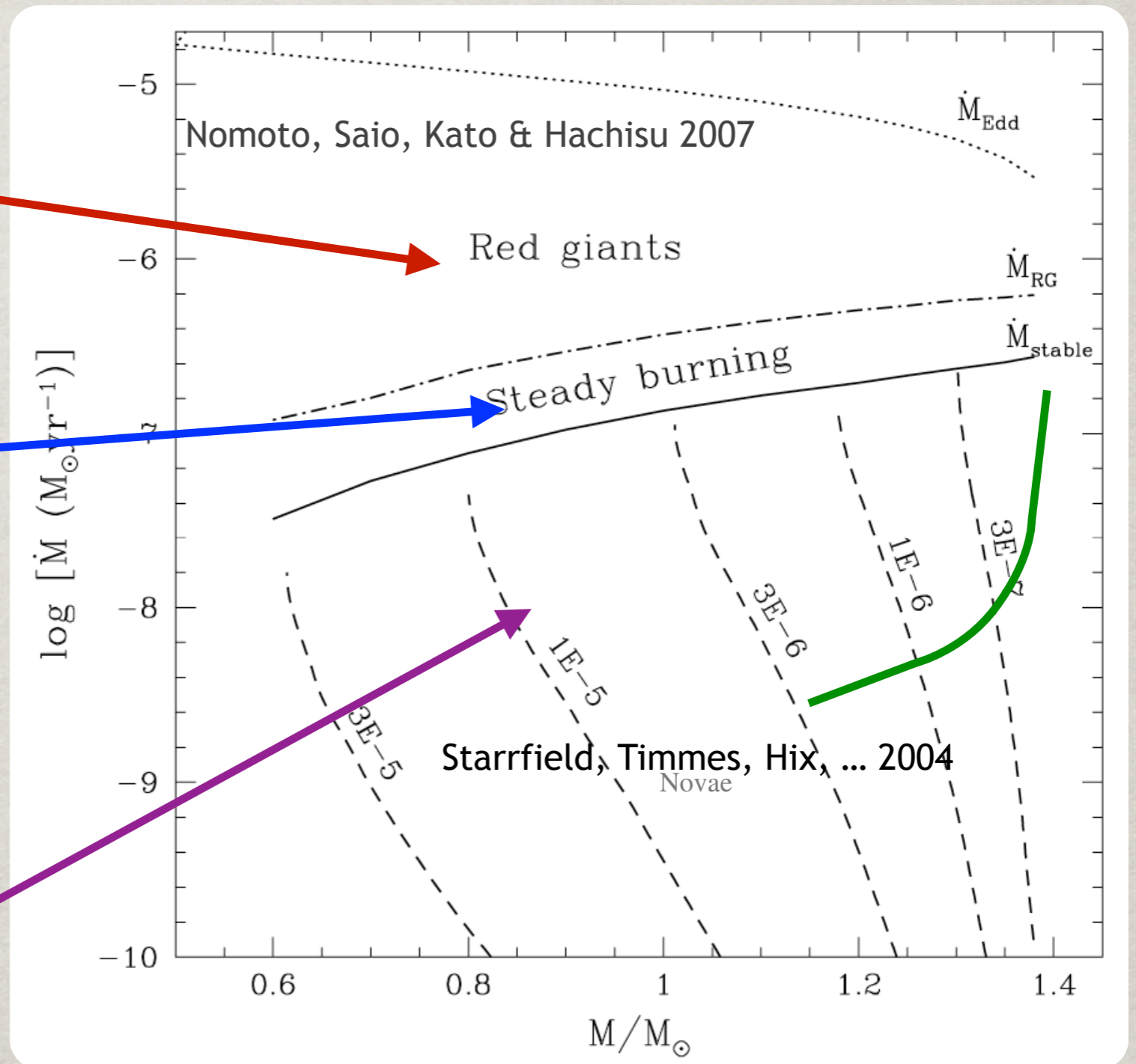
LOADING A WHITE DWARF

For high accretion rates, a **Red Giant**-like envelope re-forms.

For the right accretion rate, **H and He burn steadily** to C & O as they fall onto the white dwarf, causing it to grow.

For lesser accretion rates, **a layer of H** builds on the surface.

Accreting on a **hot white dwarf** may broaden range of accretion where the accreted matter burns to CO.

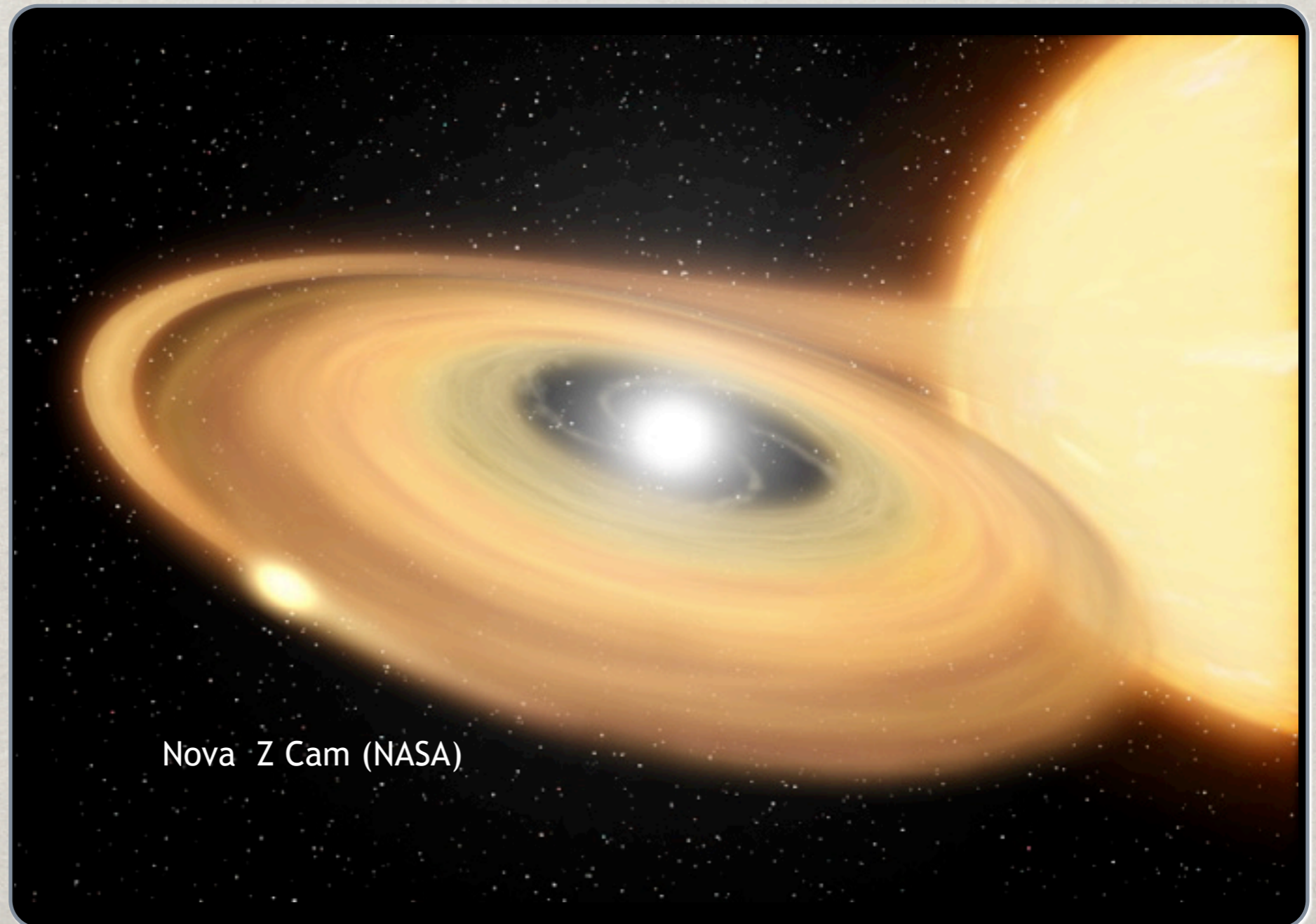


TYPES OF NOVAE

Observational Novae are categorized into 3 types based on recurrence timescale; **dwarf**, **recurrent** & **classical**.

Dwarf and some recurrent seem to be due to an **accretion disk instability**.

Others are **thermo-nuclear explosion** with recurrence time related to WD mass and accretion rate.



Nova Z Cam (NASA)

CLASSICAL NOVAE

Star Brightens a million-fold.

10^{38} J (10^{22} Megaton) Hydrogen bomb!

Delphinus

Sagitta

Aquila

< Nova V1494 Aql

T. Credner 12/99

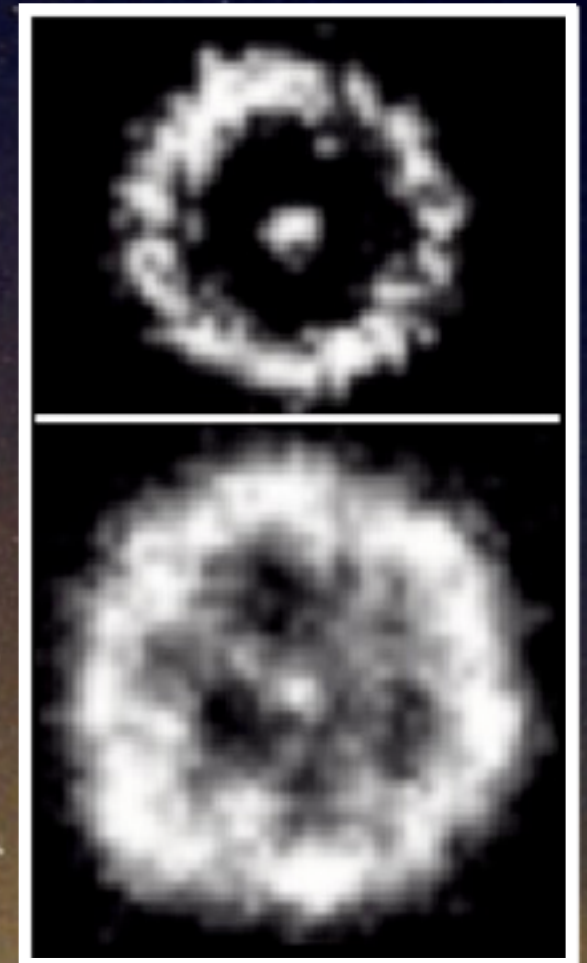
About 40 novae each year in our galaxy.

Frequently discovered by amateurs.

Mars > The binary systems that make up these cataclysmic variables are typically very tight ($a \sim R_{\odot}$) with periods of < 1 day.

Eject dust grains of material into space.

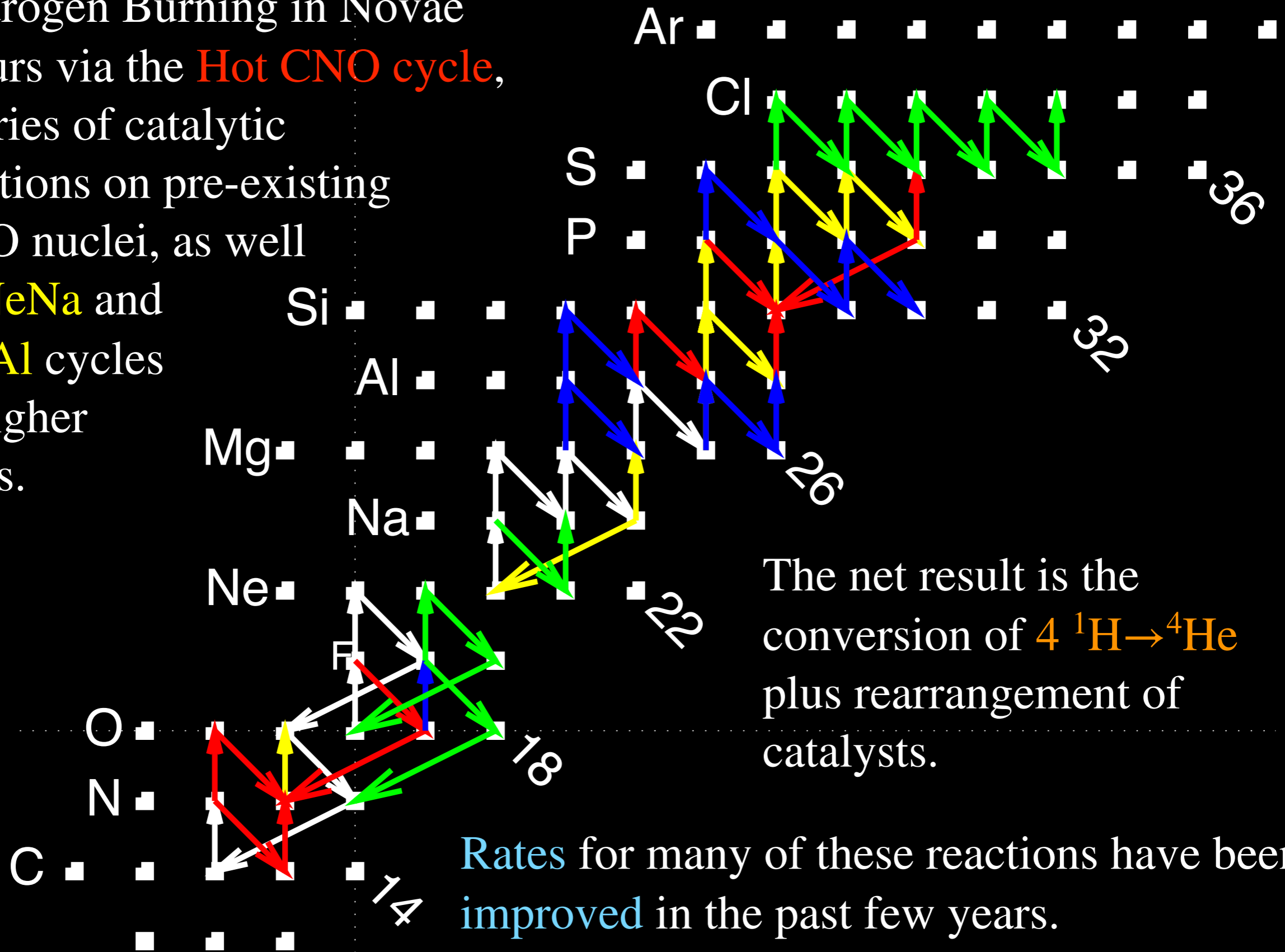
Ejecta includes White Dwarf material.



Nova T Pyxidis (HST/NASA)

NUCLEAR REACTIONS IN NOVAE

Hydrogen Burning in Novae occurs via the **Hot CNO cycle**, a series of catalytic reactions on pre-existing CNO nuclei, as well as **NeNa** and **MgAl** cycles at higher mass.



The net result is the conversion of $4\ ^1\text{H} \rightarrow\ ^4\text{He}$ plus rearrangement of catalysts.

Rates for many of these reactions have been improved in the past few years.

UPDATING NUCLEAR DATA

Models using newer rates show **significant variations** in bulk properties, like **luminosity**.

Nucleosynthesis products change by factors of two or more.

For example,

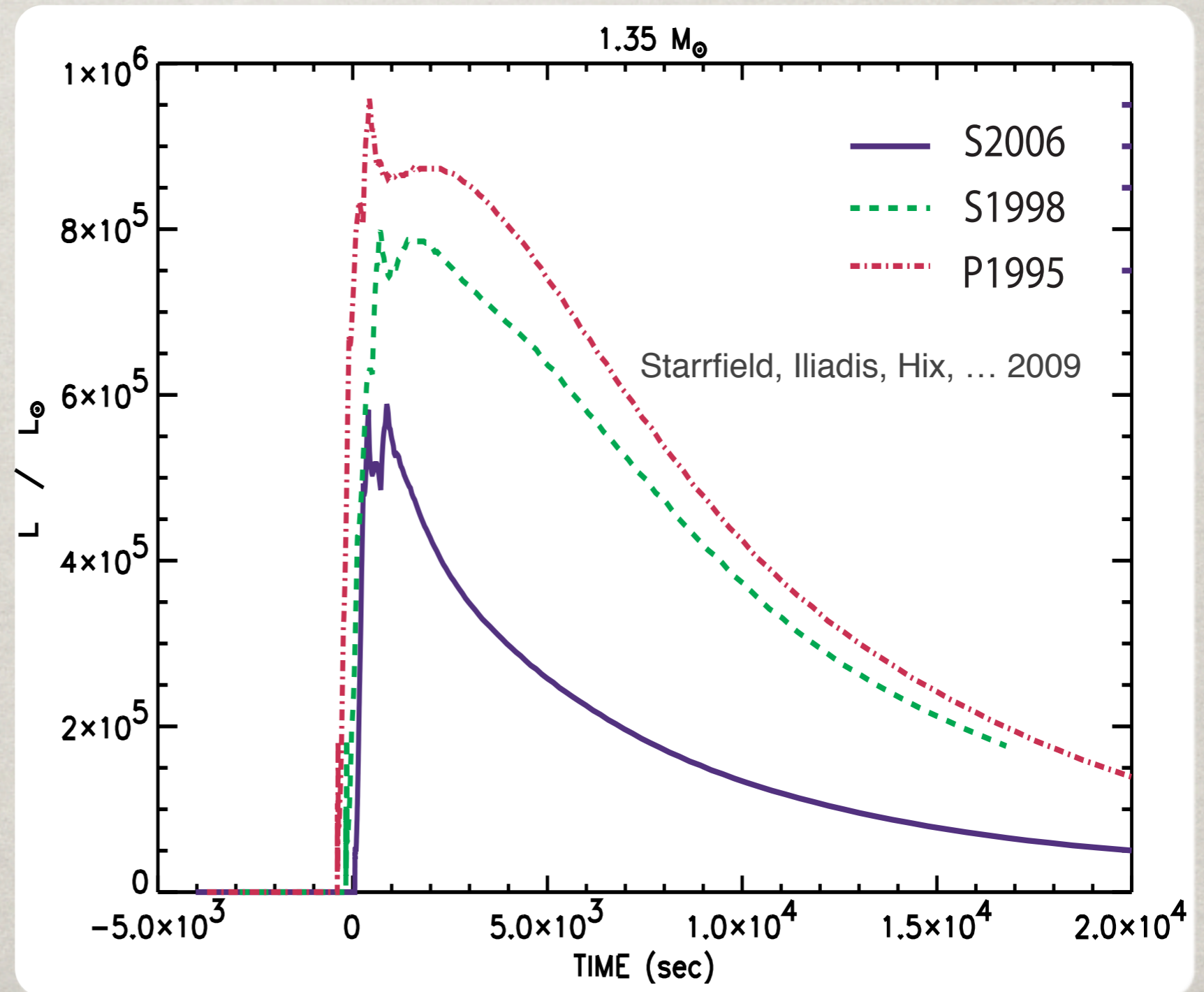
^{13}C (-17%)

^{15}N (-83%)

^{17}O (-64%)

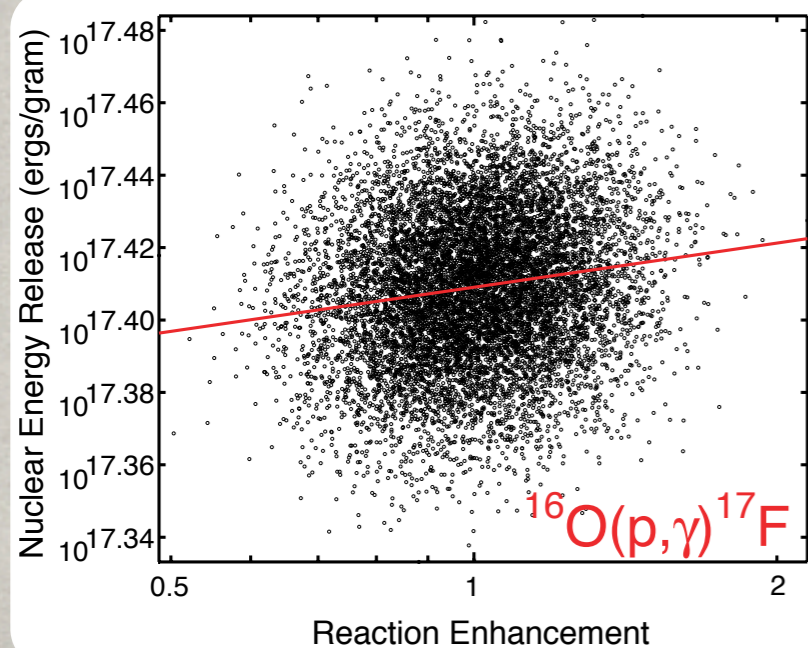
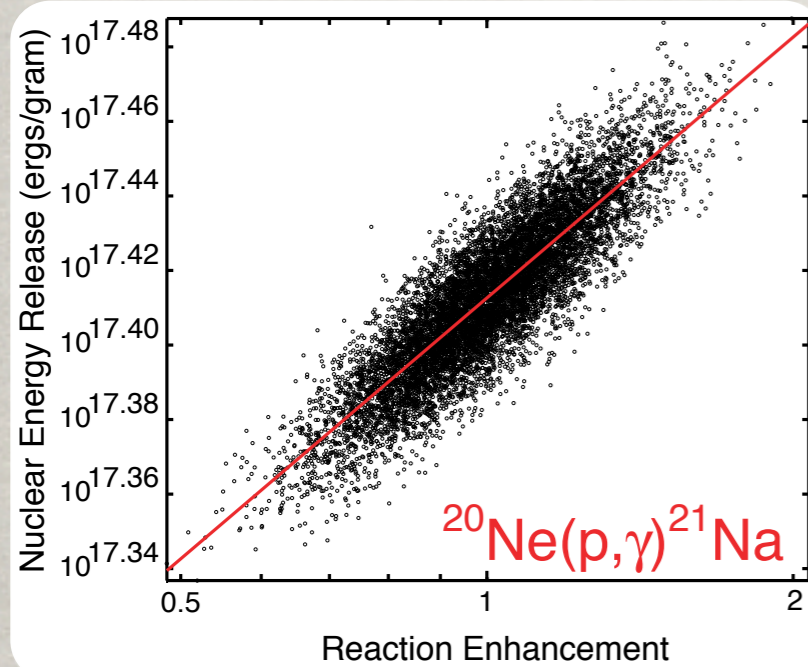
^{22}Na (-52%)

^{26}Al (+7%)

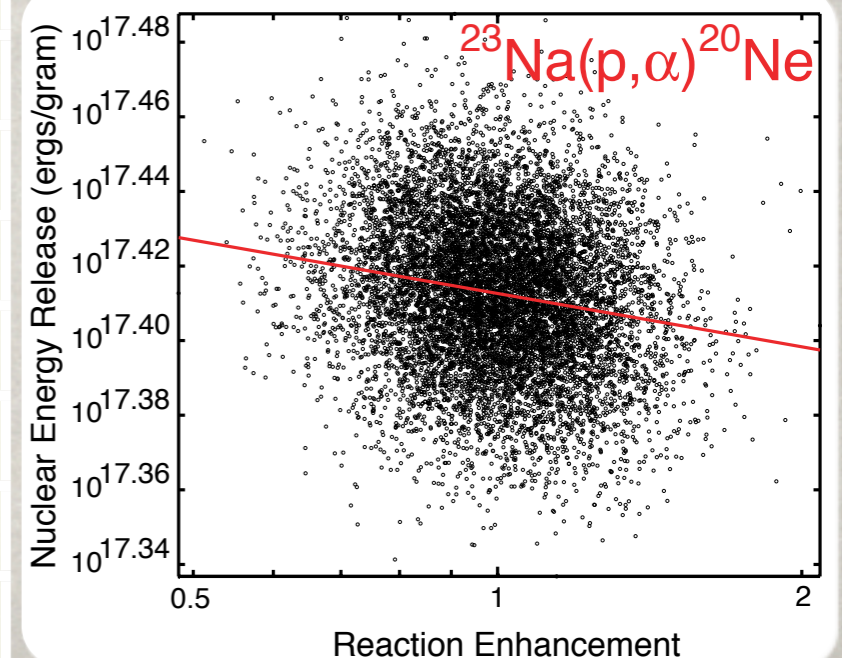
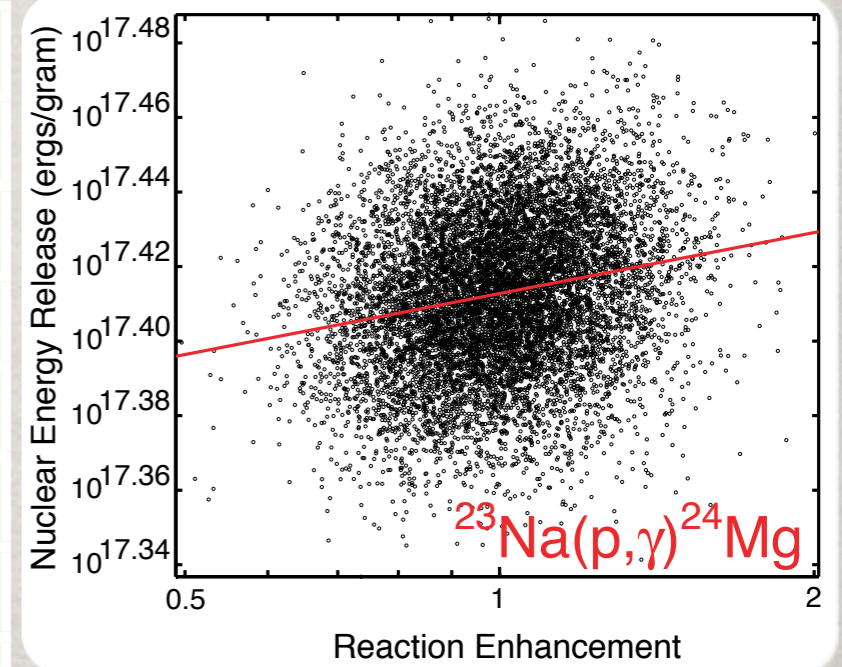


FINDING SENSITIVITIES

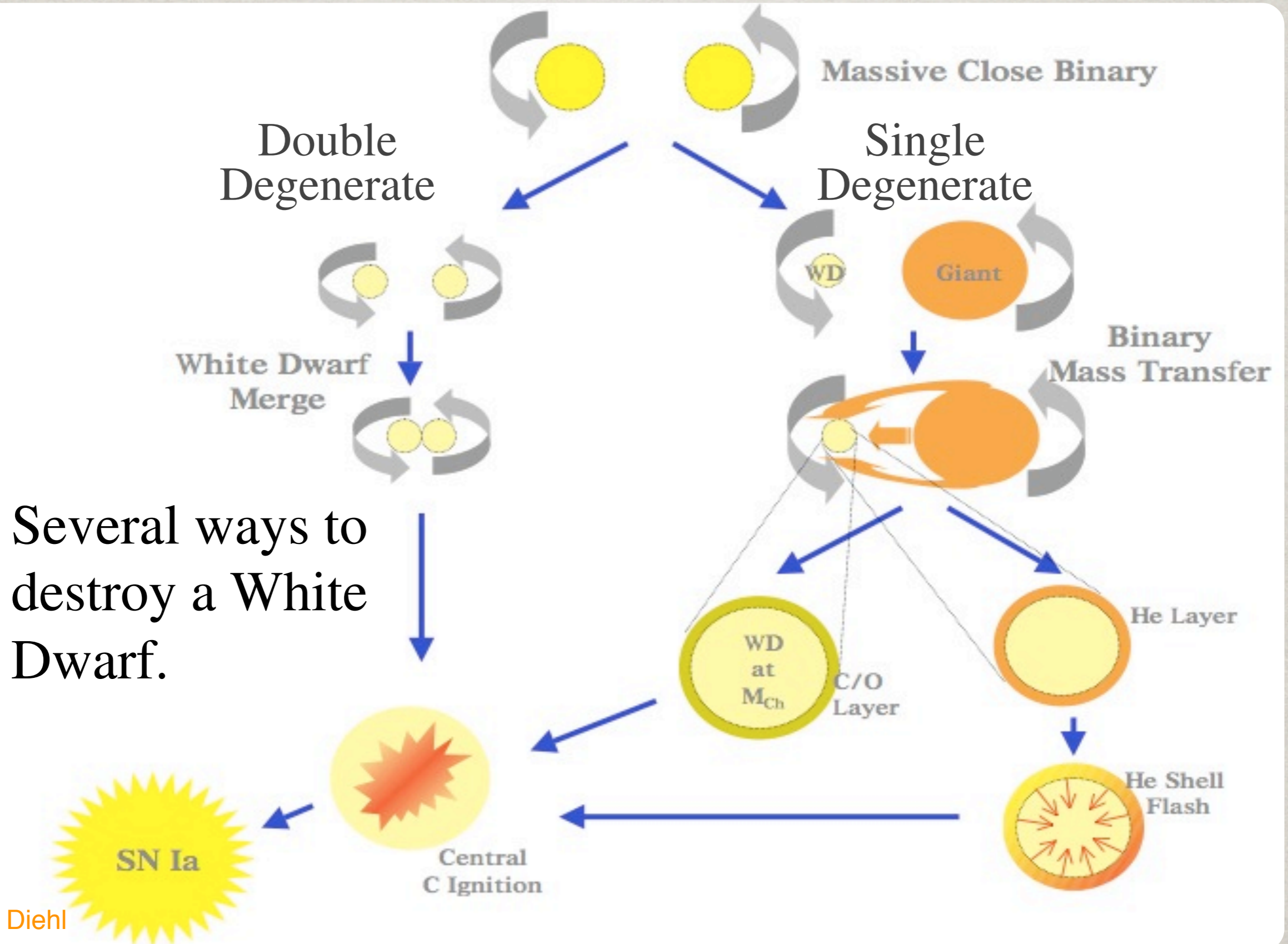
For a nova on a 1.25 solar mass WD, **Monte Carlo sensitivity analysis** indicates these reactions most strongly impact the energy generation.



Reaction	Slope
$^{20}\text{Ne}(p,\gamma)^{21}\text{Na}$	0.233 ± 0.001
$^{23}\text{Na}(p,\gamma)^{24}\text{Mg}$	0.054 ± 0.003
$^{23}\text{Na}(p,\alpha)^{20}\text{Ne}$	$-.047\pm 0.003$
$^{16}\text{O}(p,\gamma)^{17}\text{F}$	0.041 ± 0.003
$^{28}\text{Si}(p,\gamma)^{29}\text{P}$	0.025 ± 0.003
$^{23}\text{Mg}(p,\gamma)^{24}\text{Al}$	0.024 ± 0.001
$^{14}\text{N}(p,\gamma)^{15}\text{O}$	0.021 ± 0.003
$^{17}\text{F}(p,\gamma)^{18}\text{Ne}$	0.021 ± 0.001
$^{25}\text{Al}(p,\gamma)^{26}\text{Si}$	0.011 ± 0.001
$^{27}\text{Si}(p,\gamma)^{28}\text{P}$	0.010 ± 0.001
$^{25}\text{Mg}(p,\gamma)^{26}\text{Al}$	0.009 ± 0.003
$^{30}\text{P}(p,\gamma)^{31}\text{S}$	0.009 ± 0.001
$^{26}\text{Al}(p,\gamma)^{27}\text{Si}$	0.007 ± 0.001
$^{13}\text{N}(p,\gamma)^{14}\text{O}$	0.004 ± 0.001



THERMONUCLEAR SN MECHANISM



TNSN PROGENITORS

One approach to understanding the relative frequency of the different potential mechanisms is to search for the binary companions. Unfortunately, the results have been mixed.

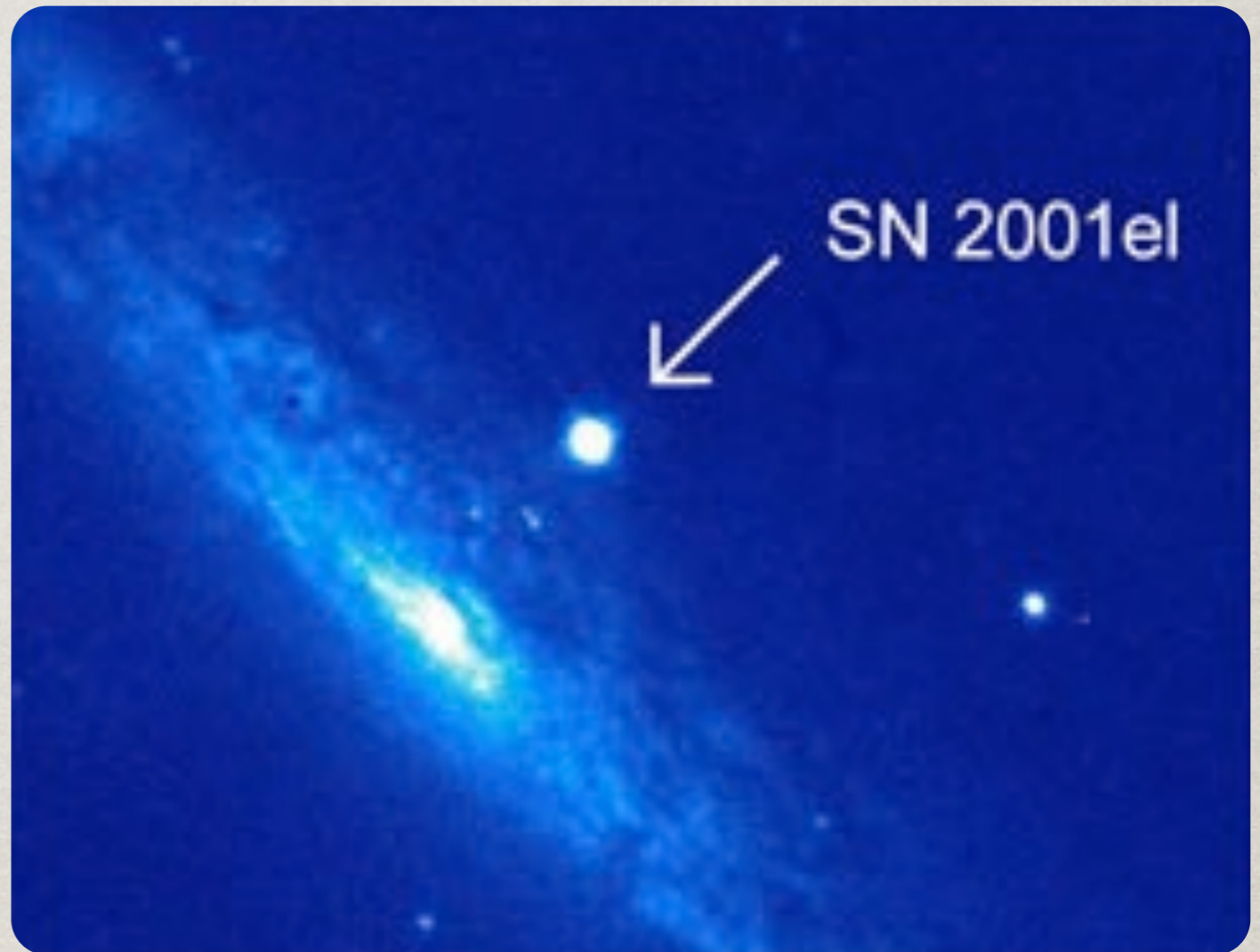
Supernova PTF 11kx, discovered 1/16/11, showed evidence of **interaction between the supernova ejecta and nova ejecta**, suggesting that the WD had previously experienced a nova outburst. This supports the single degenerate scenario, since novae have hydrogen-rich companions.

Observations of Supernova 2011fe, discovered 8/24/11 in nearby M101, lack features expected for a red giant or white dwarf companion, favoring a **main sequence star**.

Deep observations of SNR 0509–67.5 find **no suitable ex-companion star**, supporting the double degenerate scenario.

VISIBLE AT GREAT DISTANCE

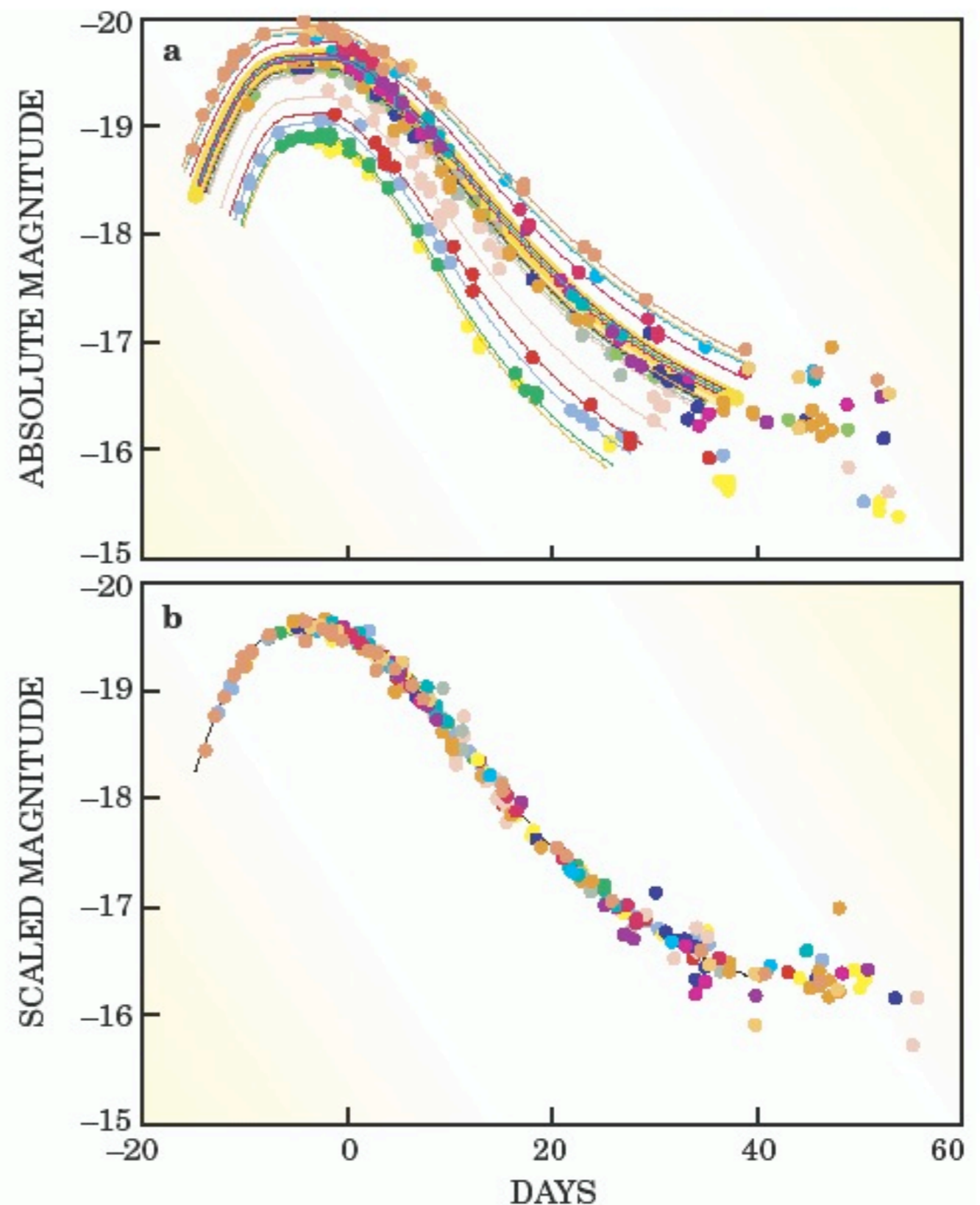
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This is particularly true for Type Ia SN, whose luminosity can be calibrated from their light curve.

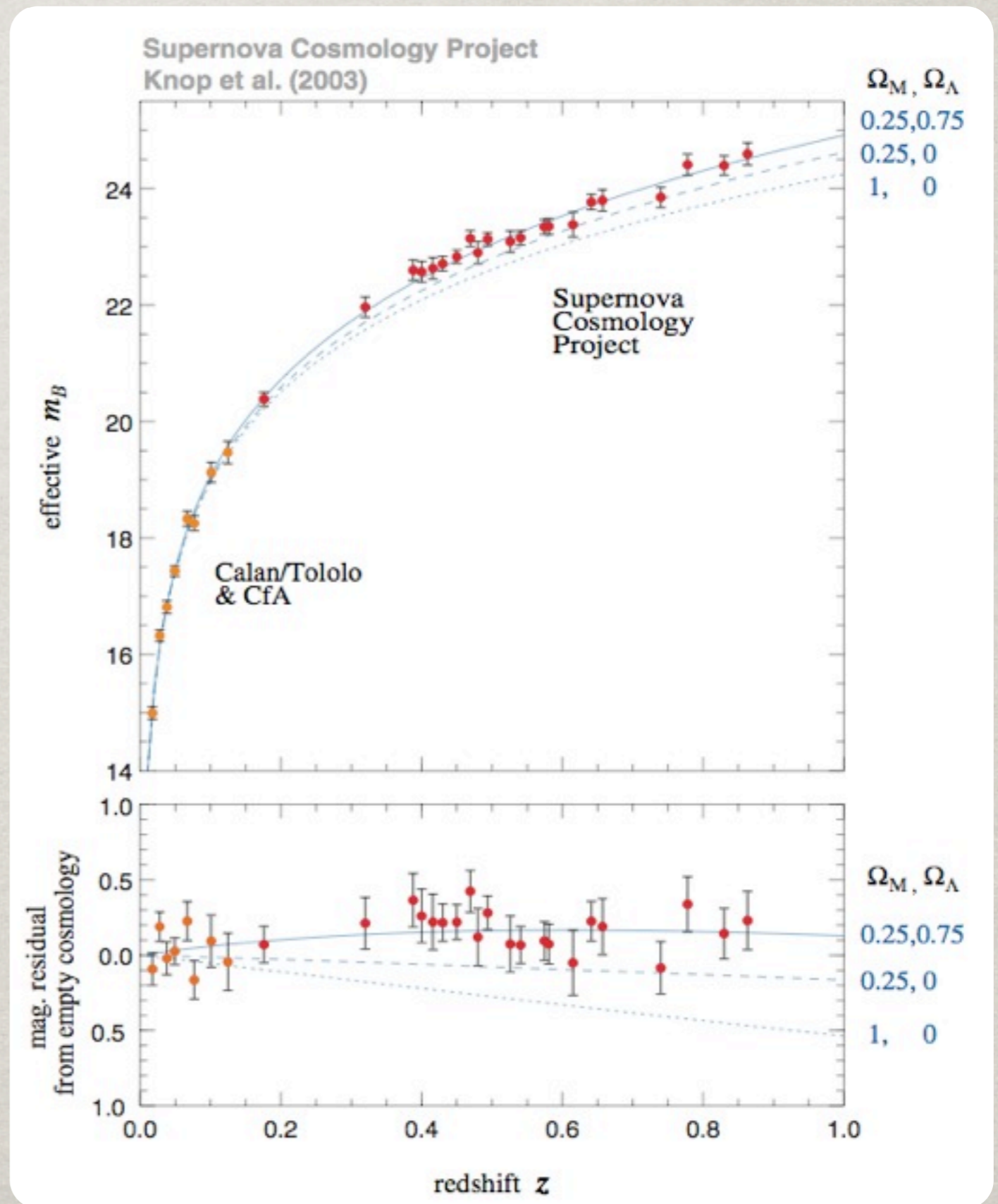


VISIBLE AT GREAT DISTANCE

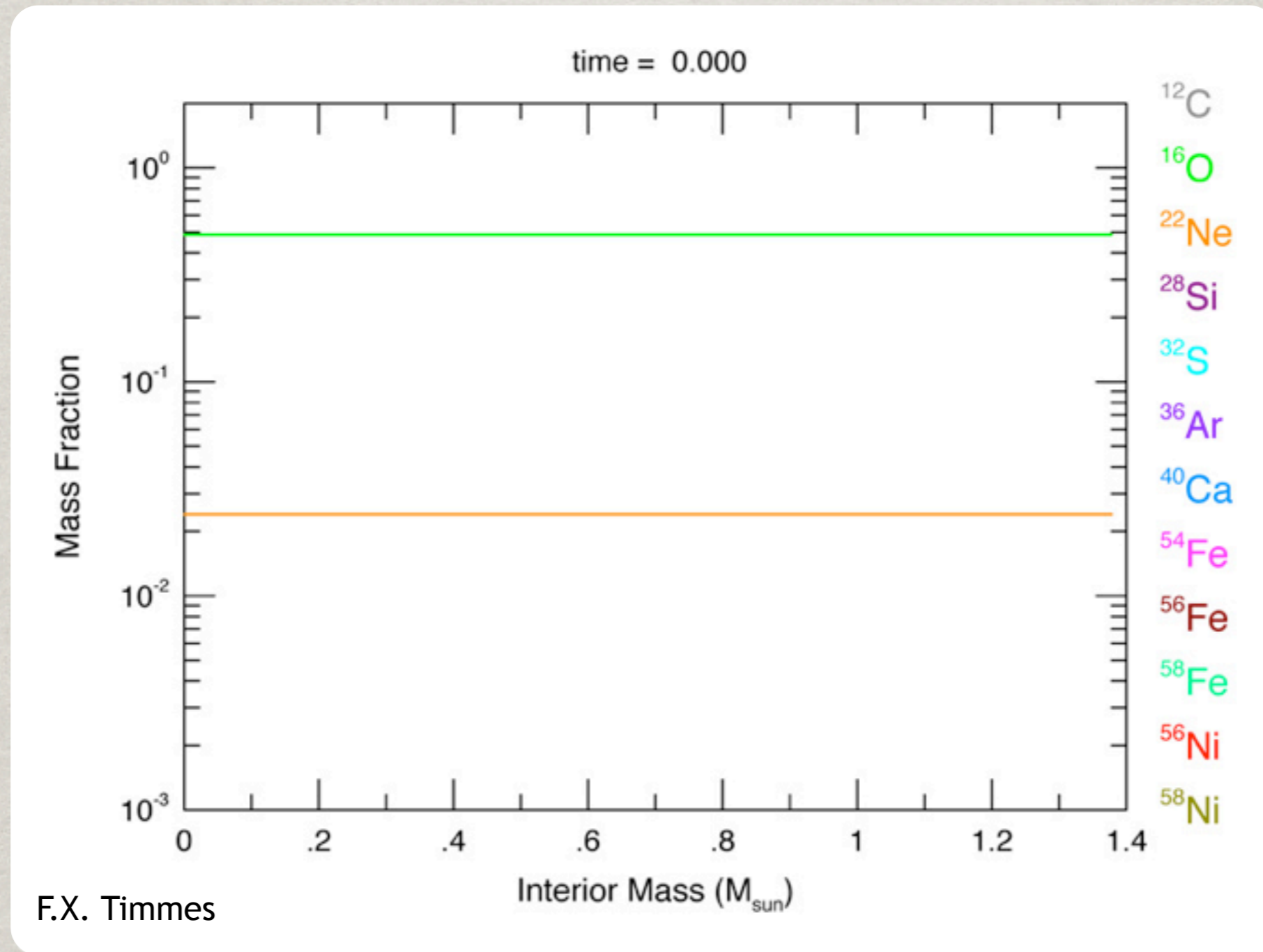
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This led to the discovery that the expansion of our universe is accelerating.

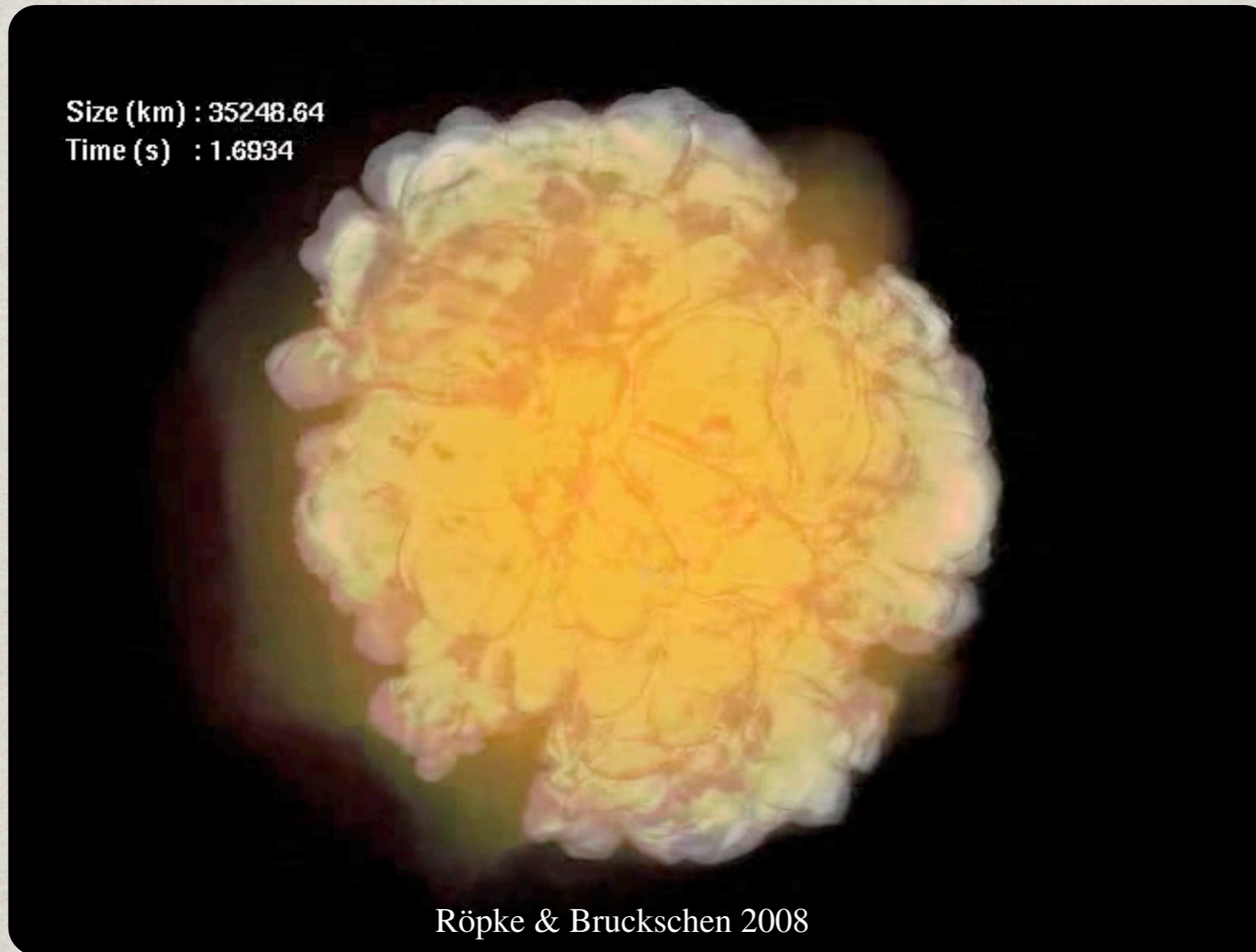


NUCLEOSYNTHESIS IN THERMONUCLEAR SN



As the central density rises, a **thermonuclear flame** is ignited which eventually **propagates throughout the star**.

THE MULTI-D VIEW

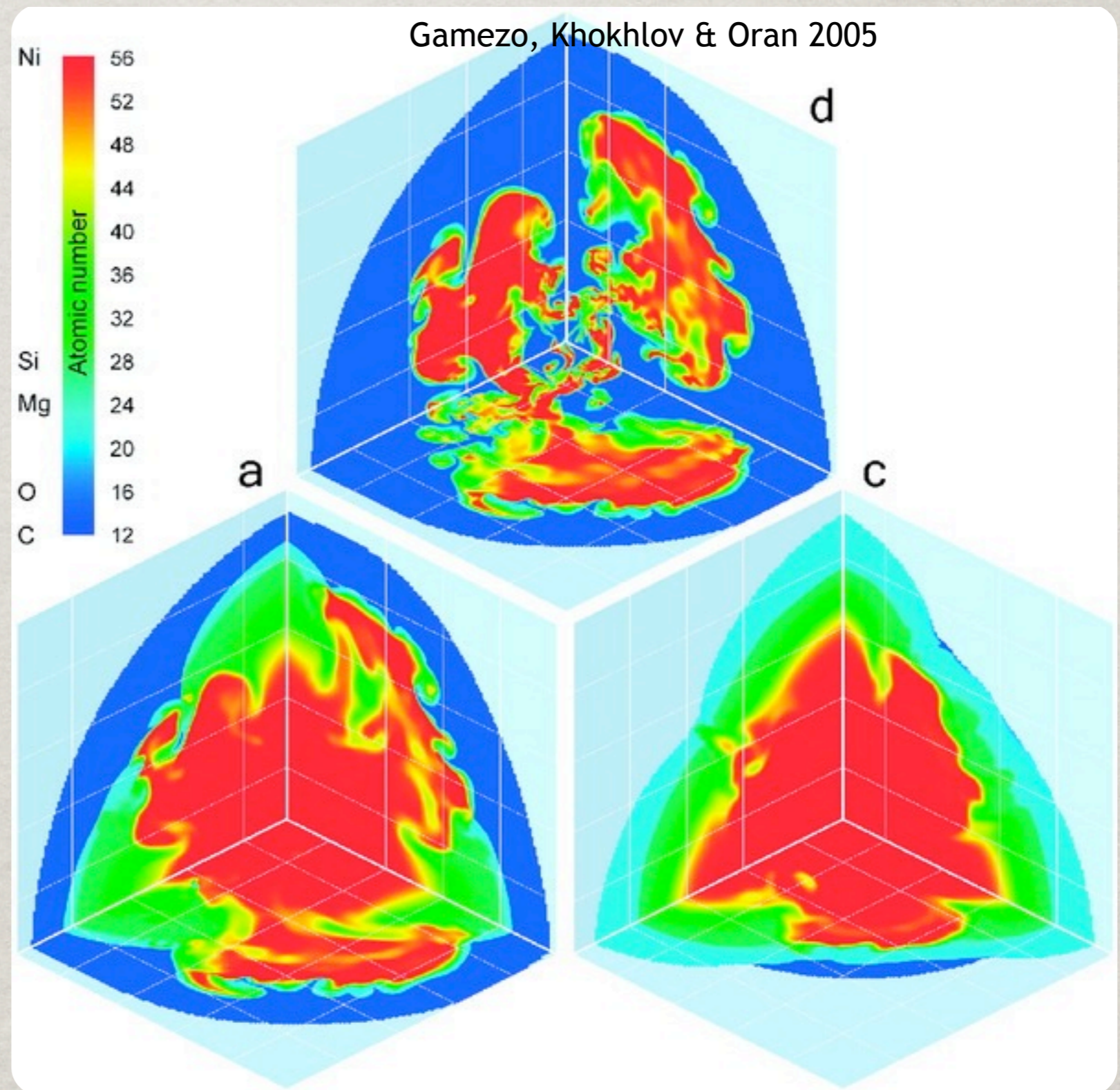


In reality, the **flame propagation** is much more **complex** than 1D implies, with turbulence and other hydrodynamic instabilities shaping the flame.

DEFLAGRATION/DETONATION

Multi-D
deflagrations
leave **pockets**
of unburned
material
behind, but
observations do
not show this.

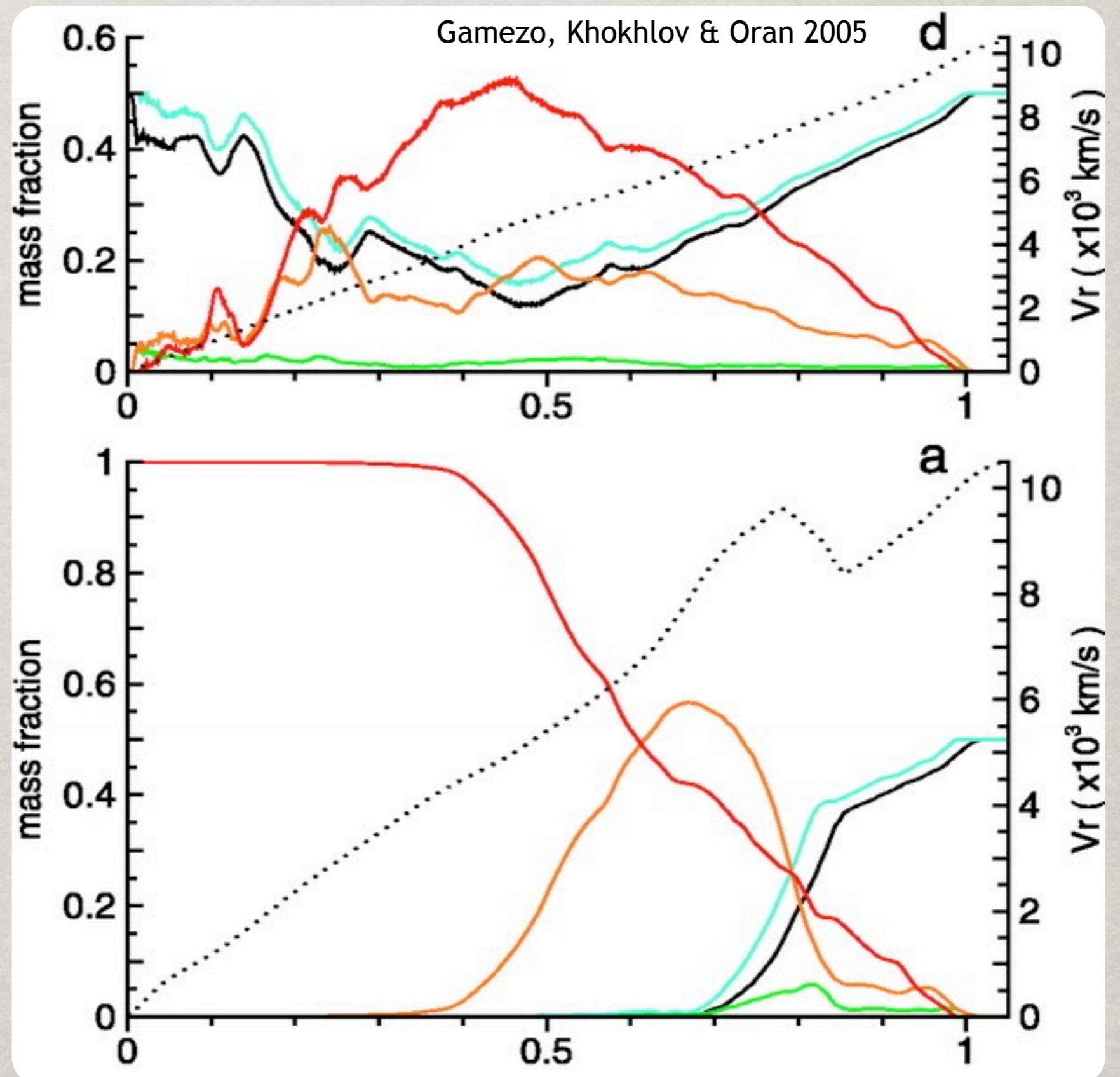
A **Deflagration**
to **Detonation**
transition
(**DDT**) must
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ONE WAY TO DDT?

White Dwarf Deflagration

Resolution: 6 km

Initial Bubble Radius: 18 km

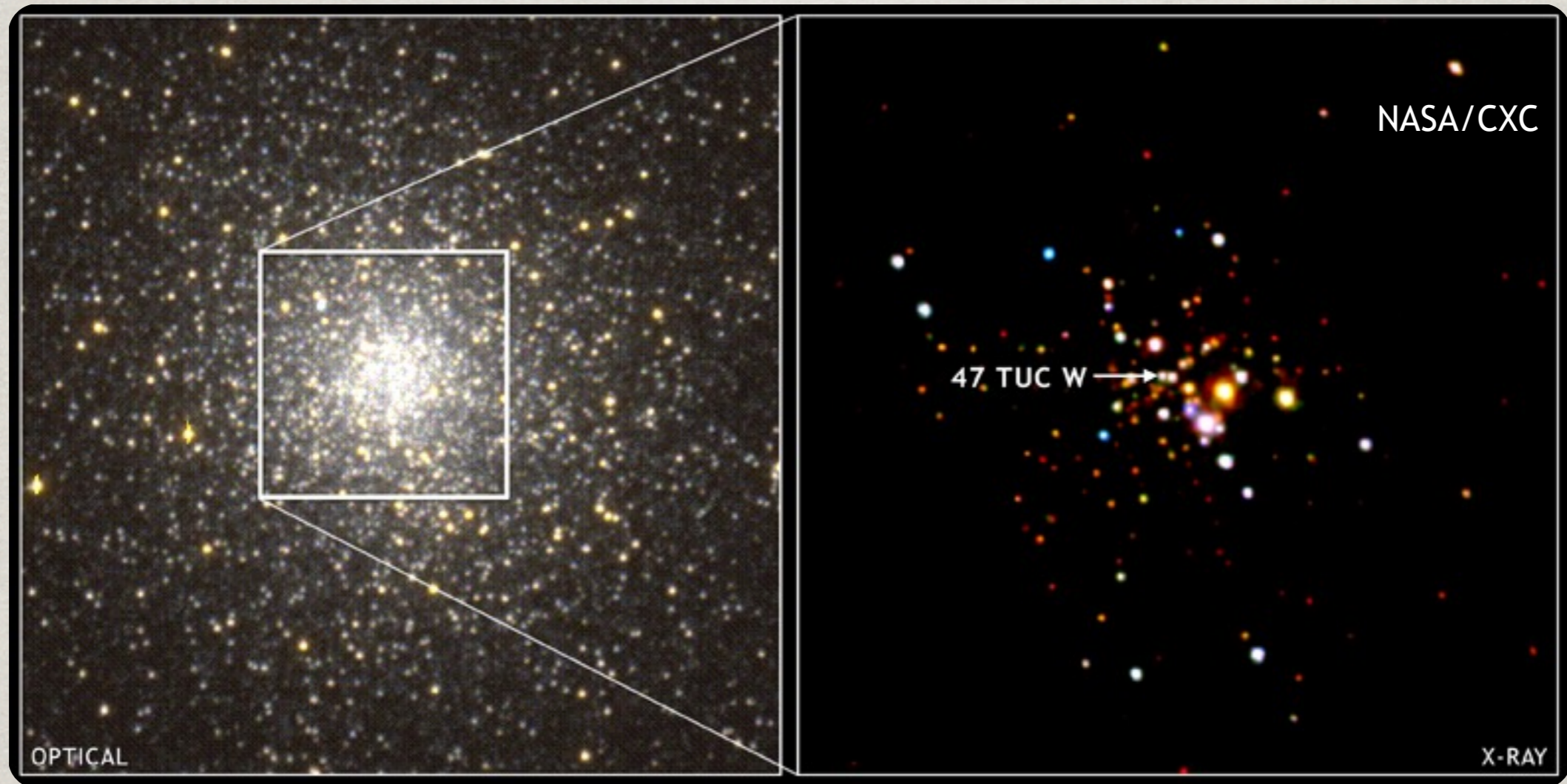
Ignition Offset: 42 km

Variable 1: Density [1.5e+07 - 2.0e+07]

Variable 2: Reaction Progress [0.0 - 1.0]

In terrestrial conditions, like a back-firing engine, **DDT** is often due to [geometry/confinement](#). Perhaps this is also true for Thermonuclear SN.

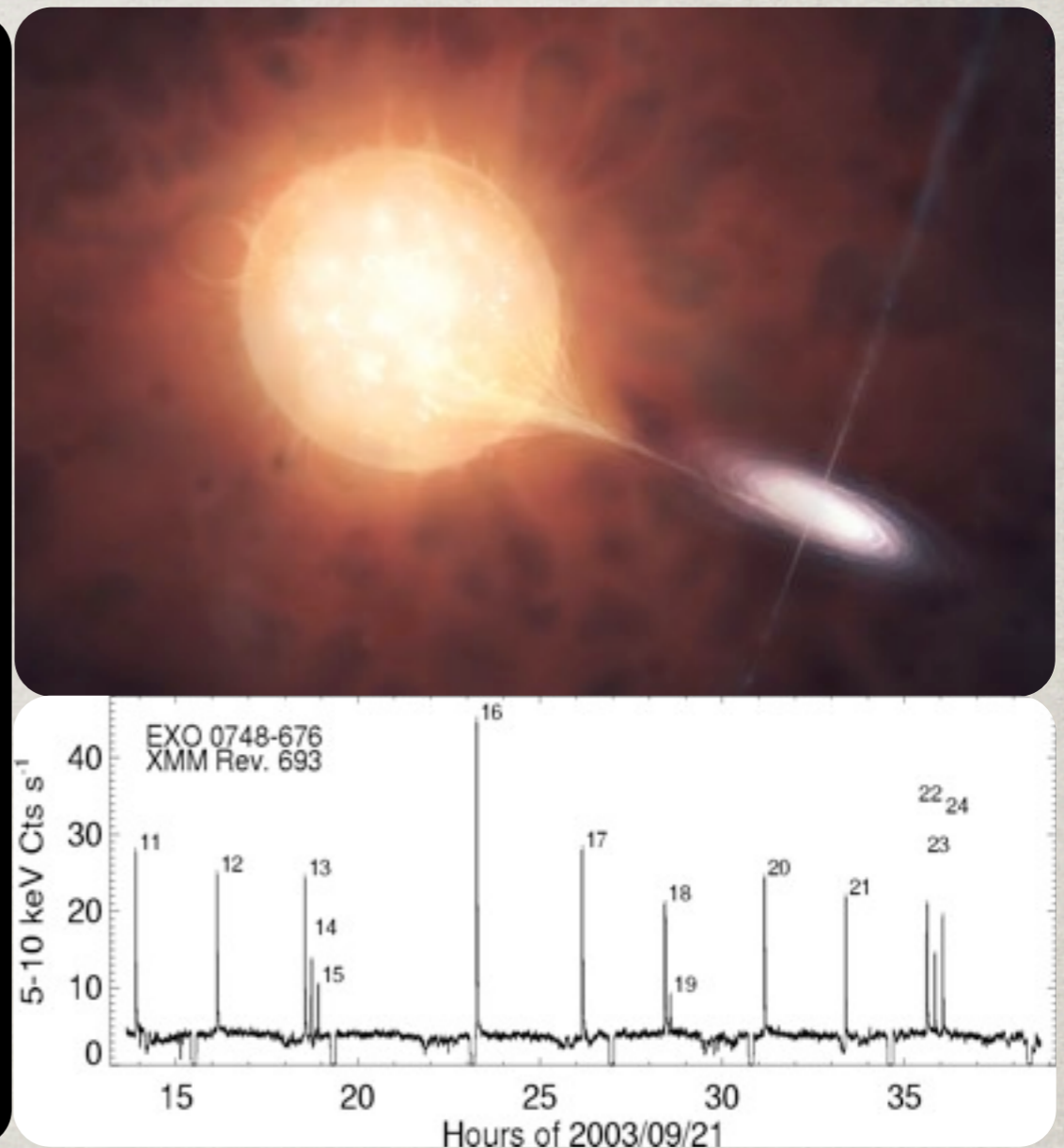
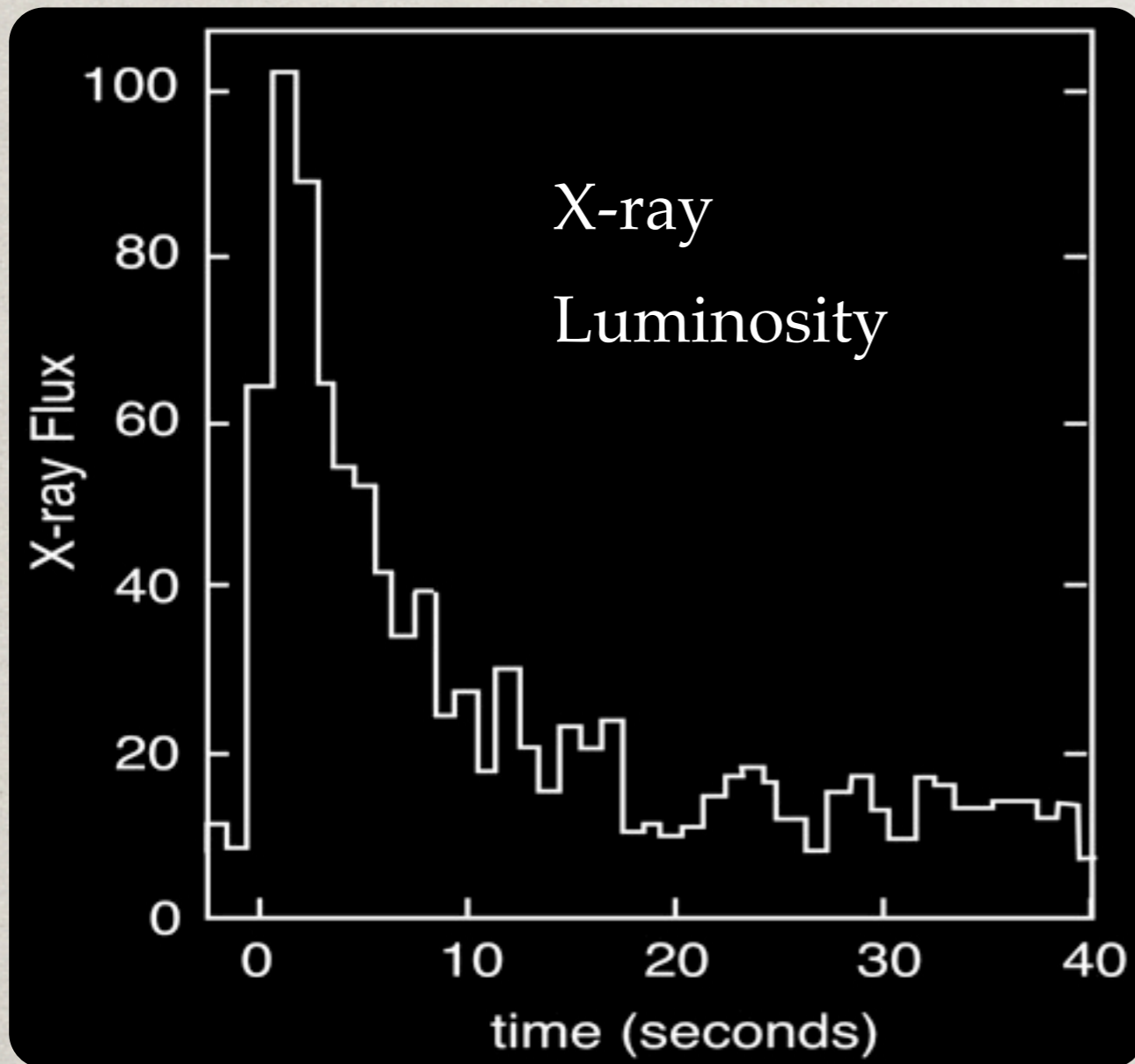
ACCRETION ONTO NEUTRON STARS



Accretion from a companion can also occur onto a neutron star. Such **accreting neutron stars** account for many **observed X-ray sources**.

Unlike for a white dwarf, **nuclear energy release** ($\sim 0.01 mc^2$) is **insignificant** compared to the **gravitational energy** ($\sim 0.1 mc^2$) release, unless it is **intermittent**.

X-RAY BURSTS

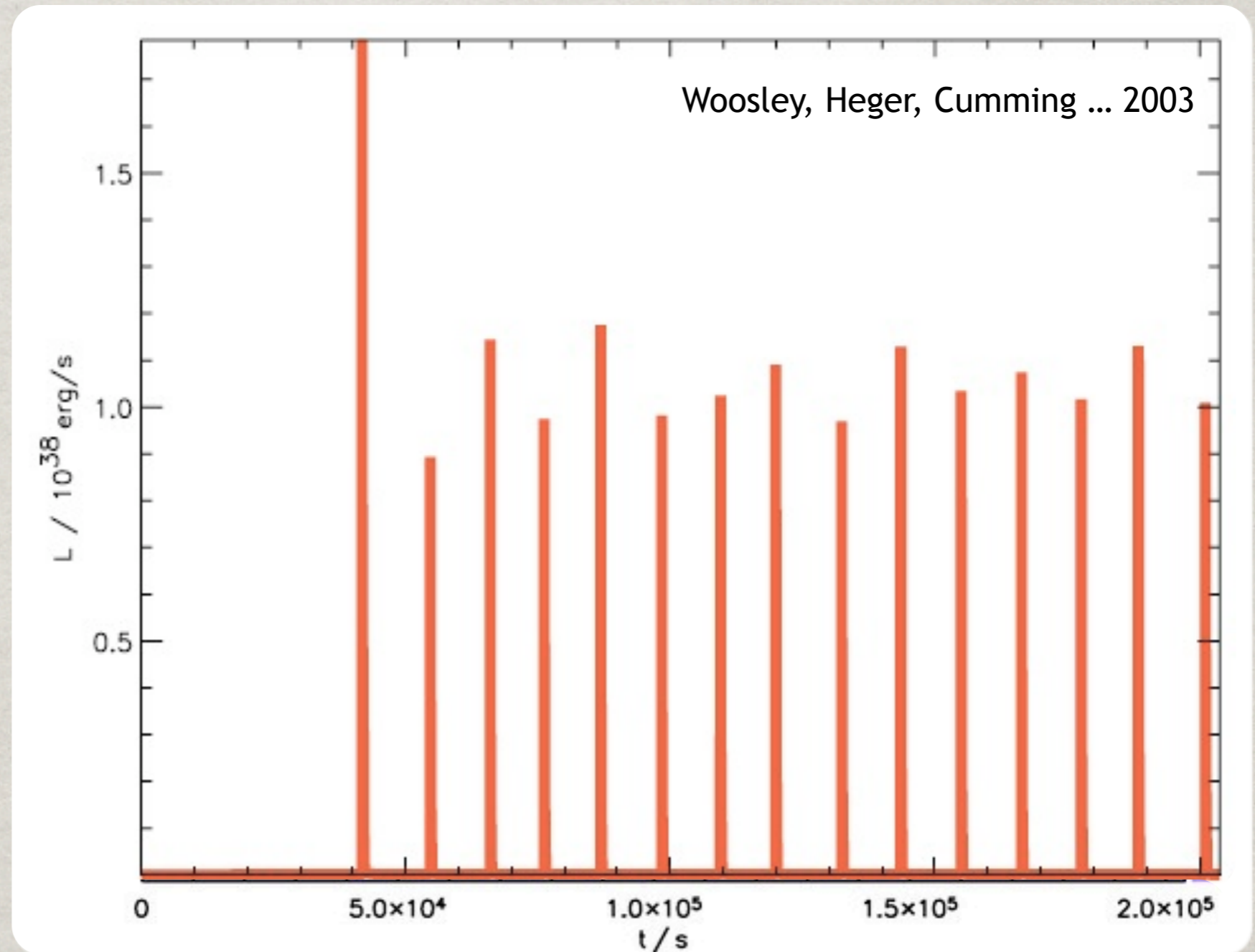


Some X-ray binaries periodically produce **seconds-long bursts** of X-rays ($\sim 10^{38}$ ergs s⁻¹) can recur hourly or daily.

These Type 1 X-ray bursts are due **unstable H-He burning** on the neutron star surface.

BETTER MODELING FOR XRB

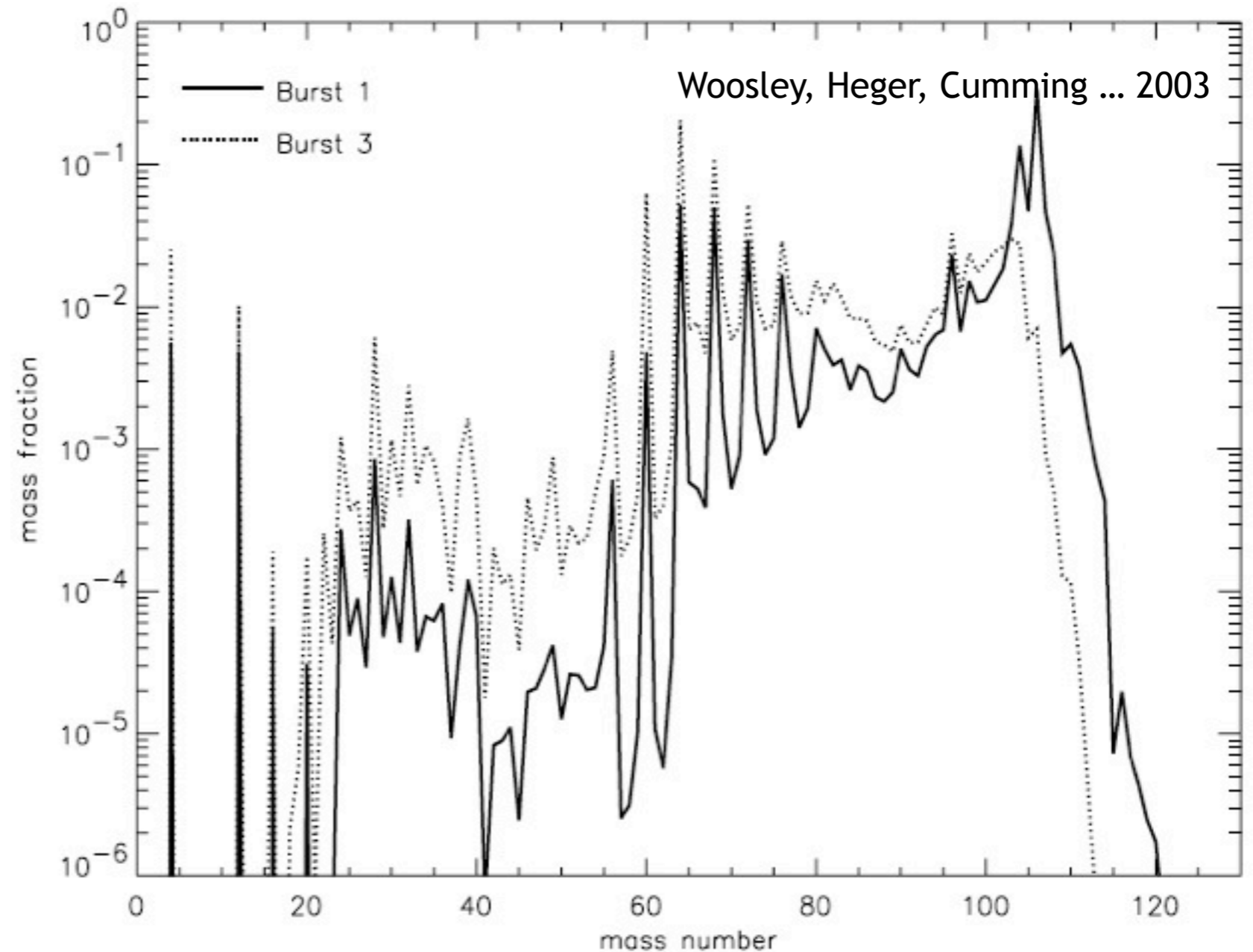
Single zone models have been replaced by true 1D hydrodynamic models which include large networks and General Relativity.



These improved models have taught us the importance of the ash to subsequent bursts, as the heavy element ash “dilutes” the accreted matter, weakening the burst.

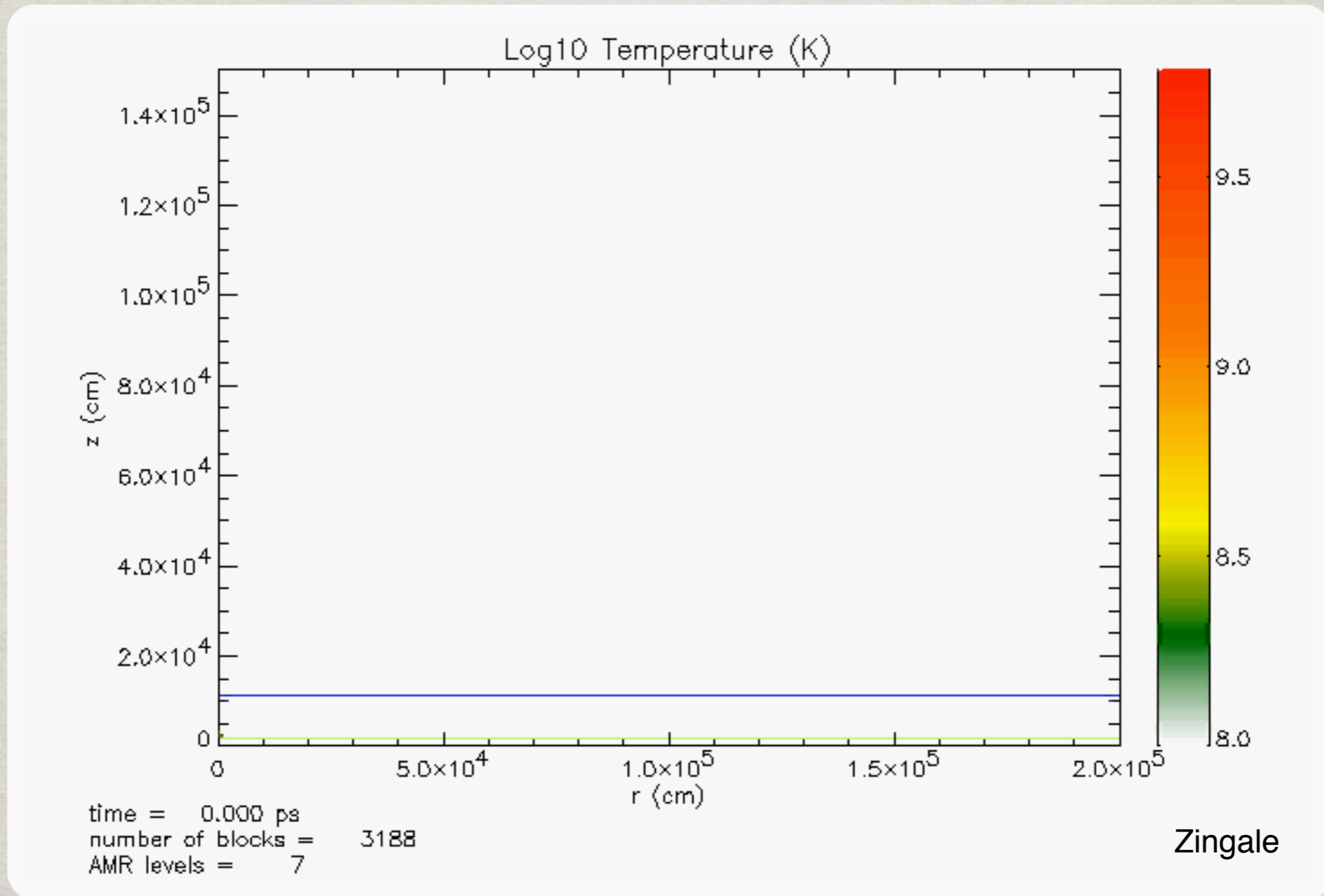
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XRB IN MULTI-D

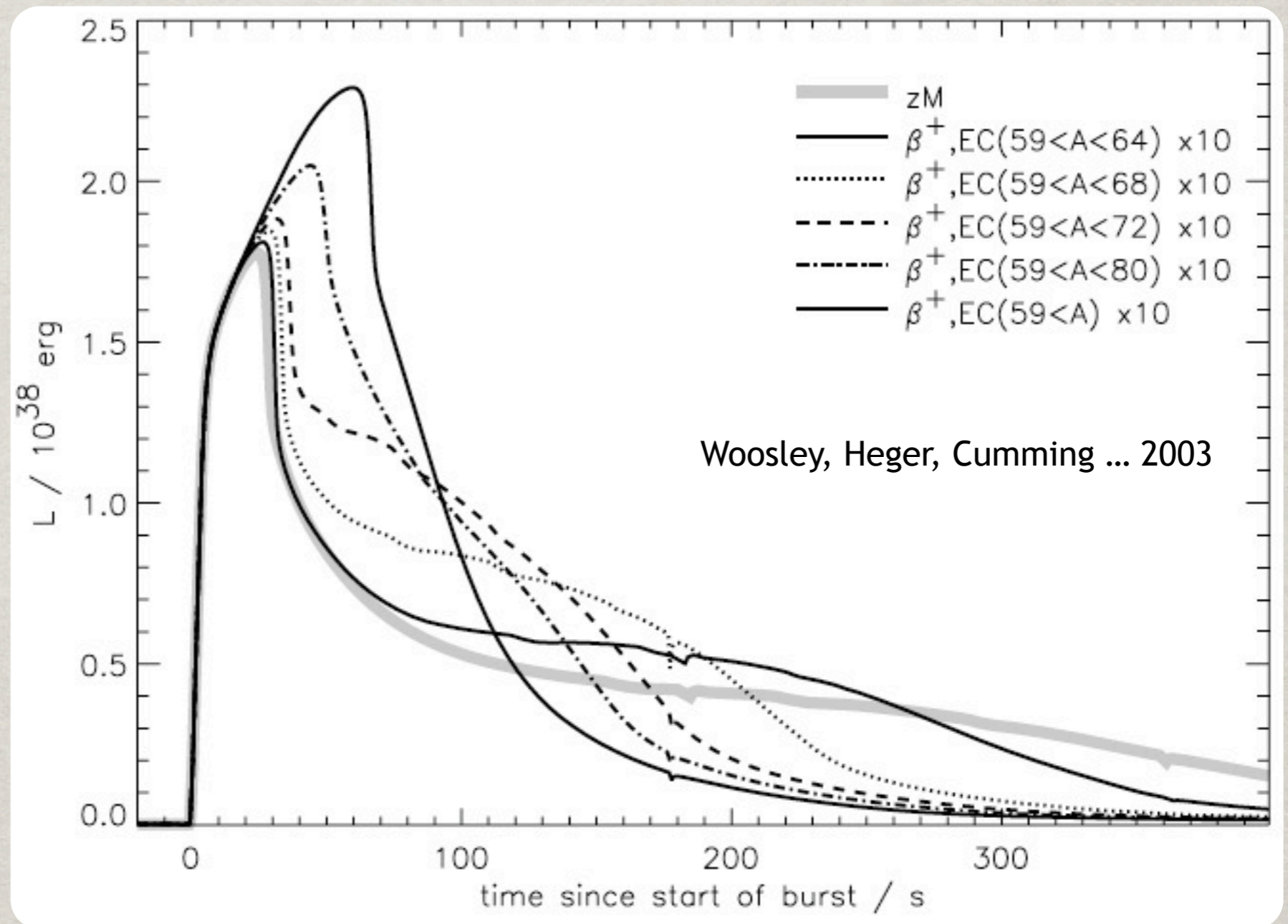


While **1D models** (by design) assume simultaneous ignition over **the entire surface**, ignition of the XRB takes **several hundred μ s** to **circle the neutron star**.

NUCLEAR PHYSICS INFLUENCE ON X-RAY BURSTS

Many of the rp-process waiting points are β^+ decays.

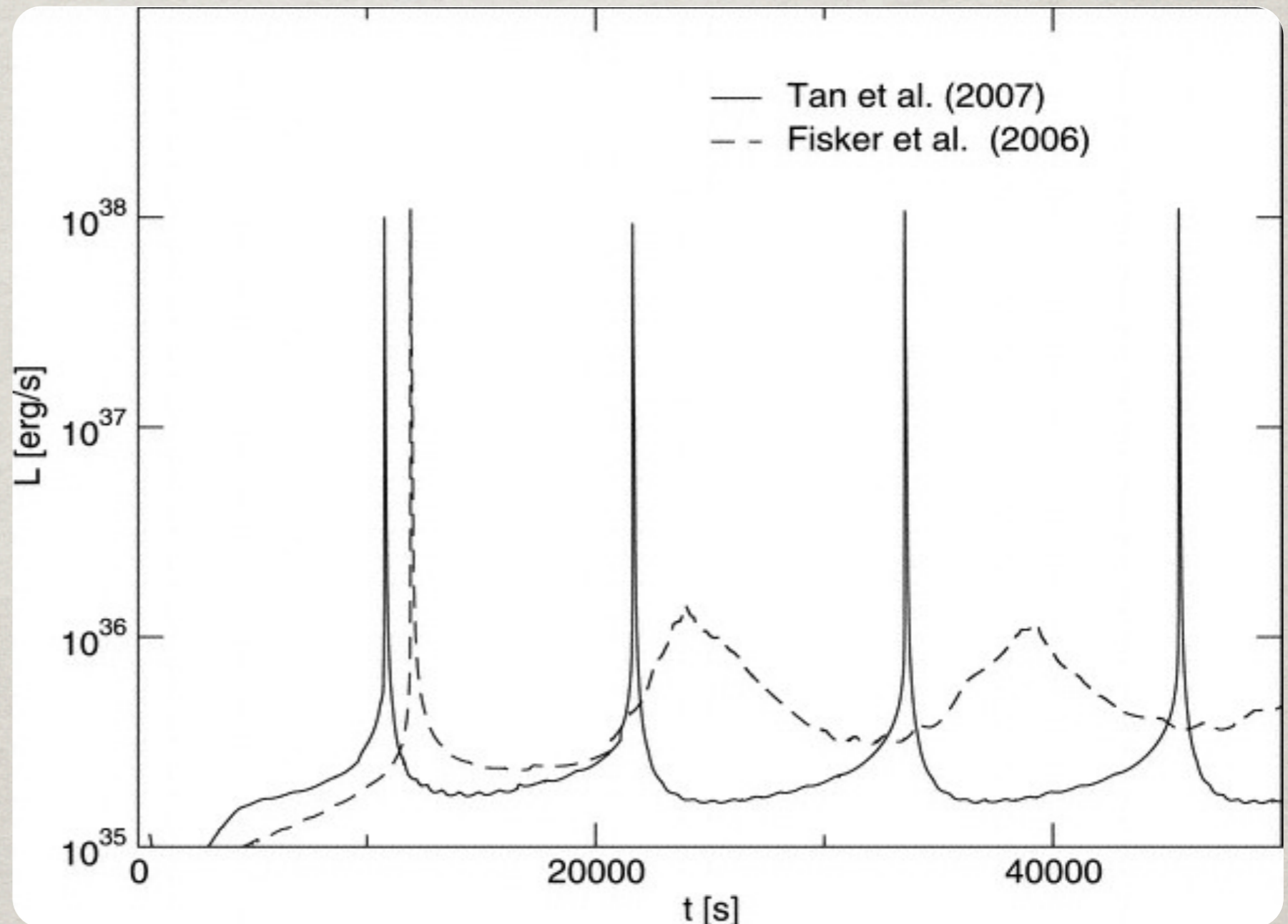
Variations of β^+ decays produce noticeable effects on XRB luminosity.



NUCLEAR PHYSICS INFLUENCE ON X-RAY BURSTS

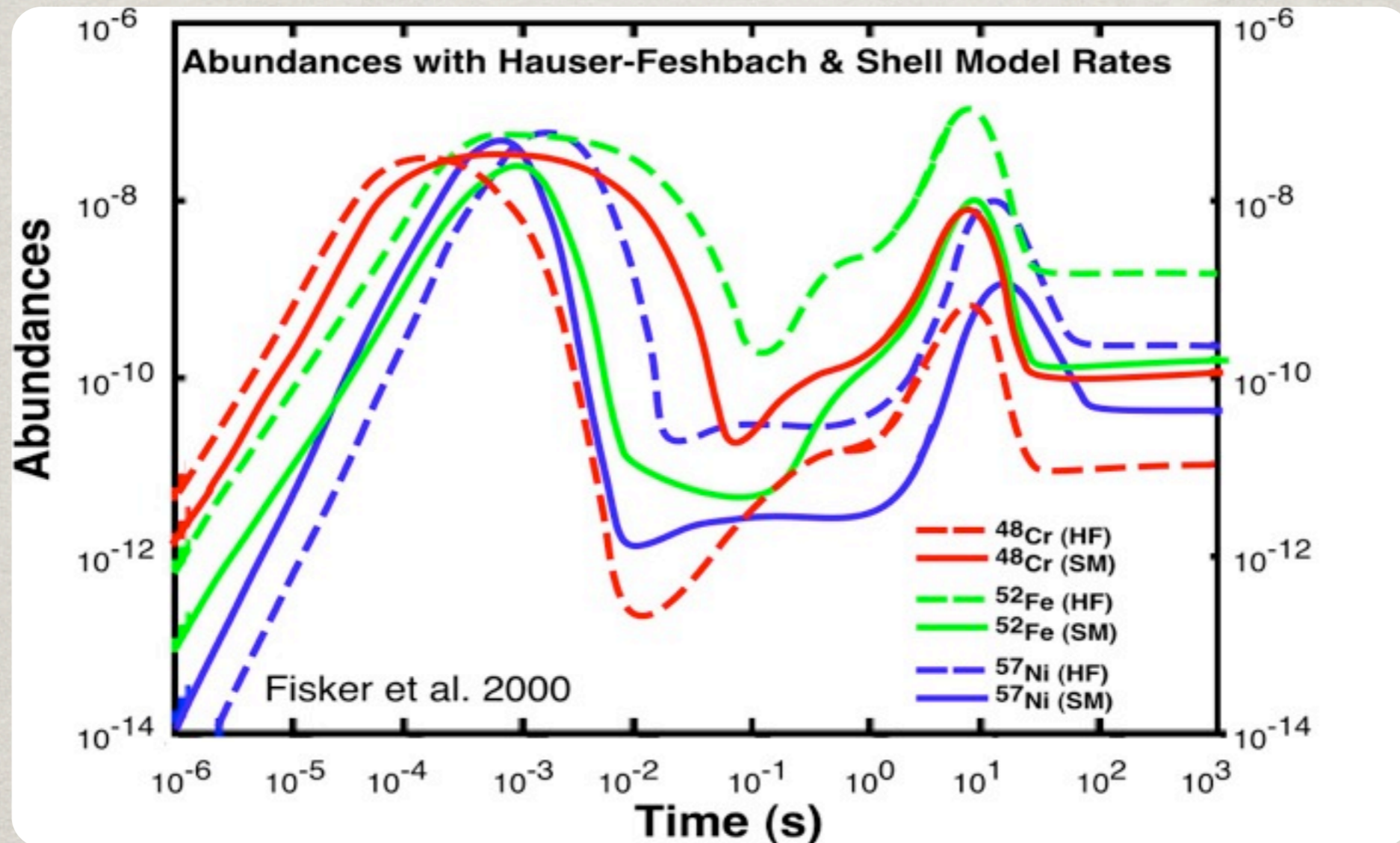
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XRB luminosity has also shown significant sensitivities to $^{15}\text{O}(\alpha,\gamma)^{19}\text{Ne}$ (Fisker, Tan, Görres & Wiescher 2007)

NUCLEAR PHYSICS INFLUENCE ON X-RAY BURSTS



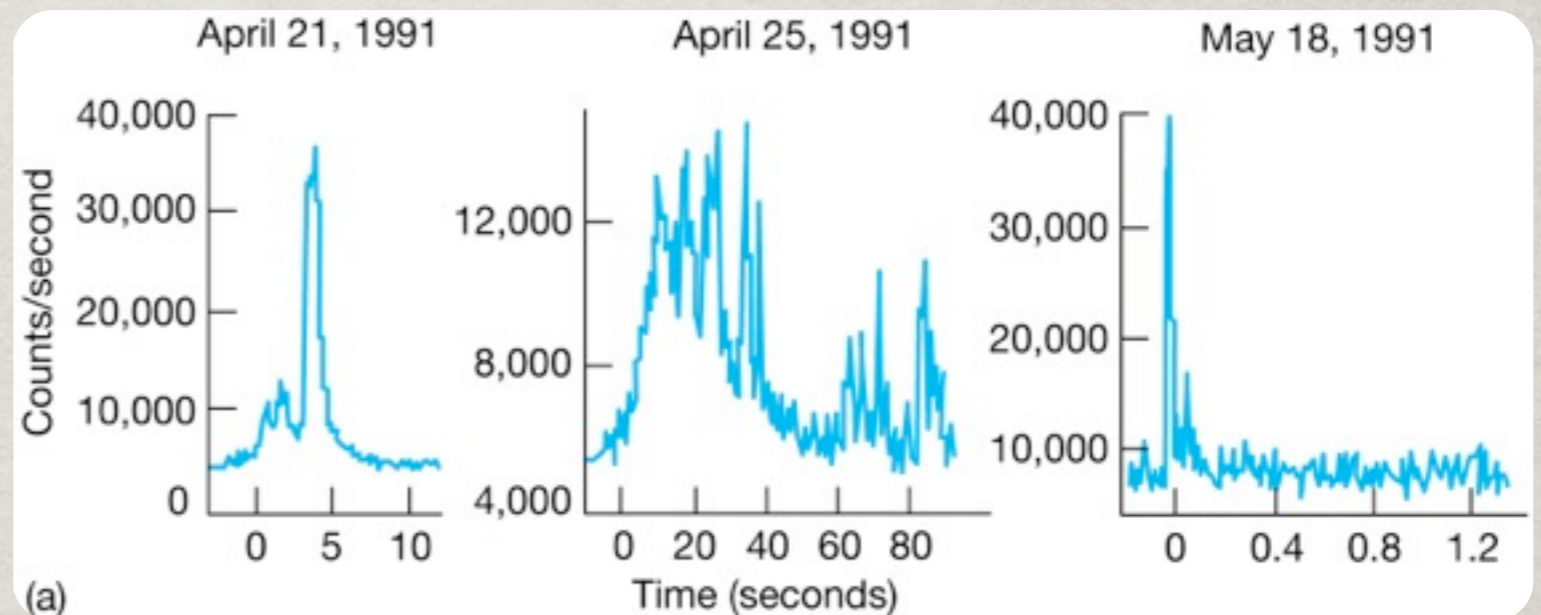
Nucleosynthesis in XRBs is also sensitive to nuclear reaction rates.

Improved **shell model calculations** for new rates change predictions of synthesized abundances **$\sim 10x$** .

GAMMA-RAY BURSTS

Intense bursts of gamma-rays are also observed lasting seconds and occurring roughly **once per day**.

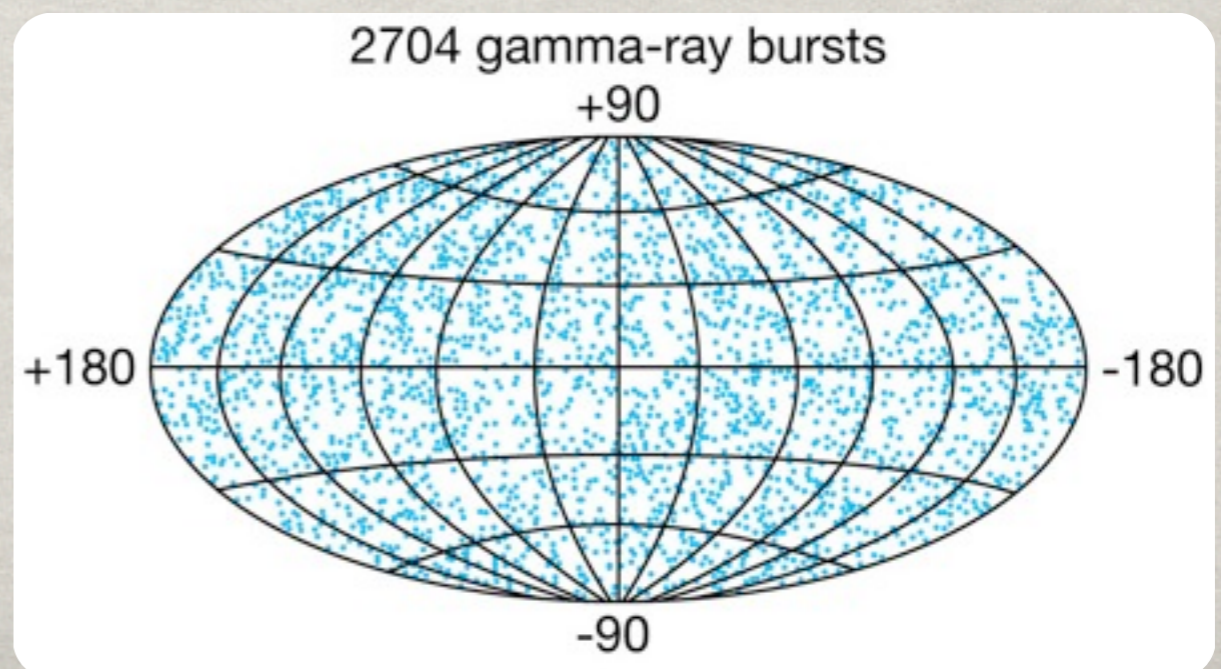
The bursts can be divided into **short** and **long duration** bursts ($t \sim 0.3$ s and ~ 30 s, respectively).



Maps of the burst's distribution on the sky show no

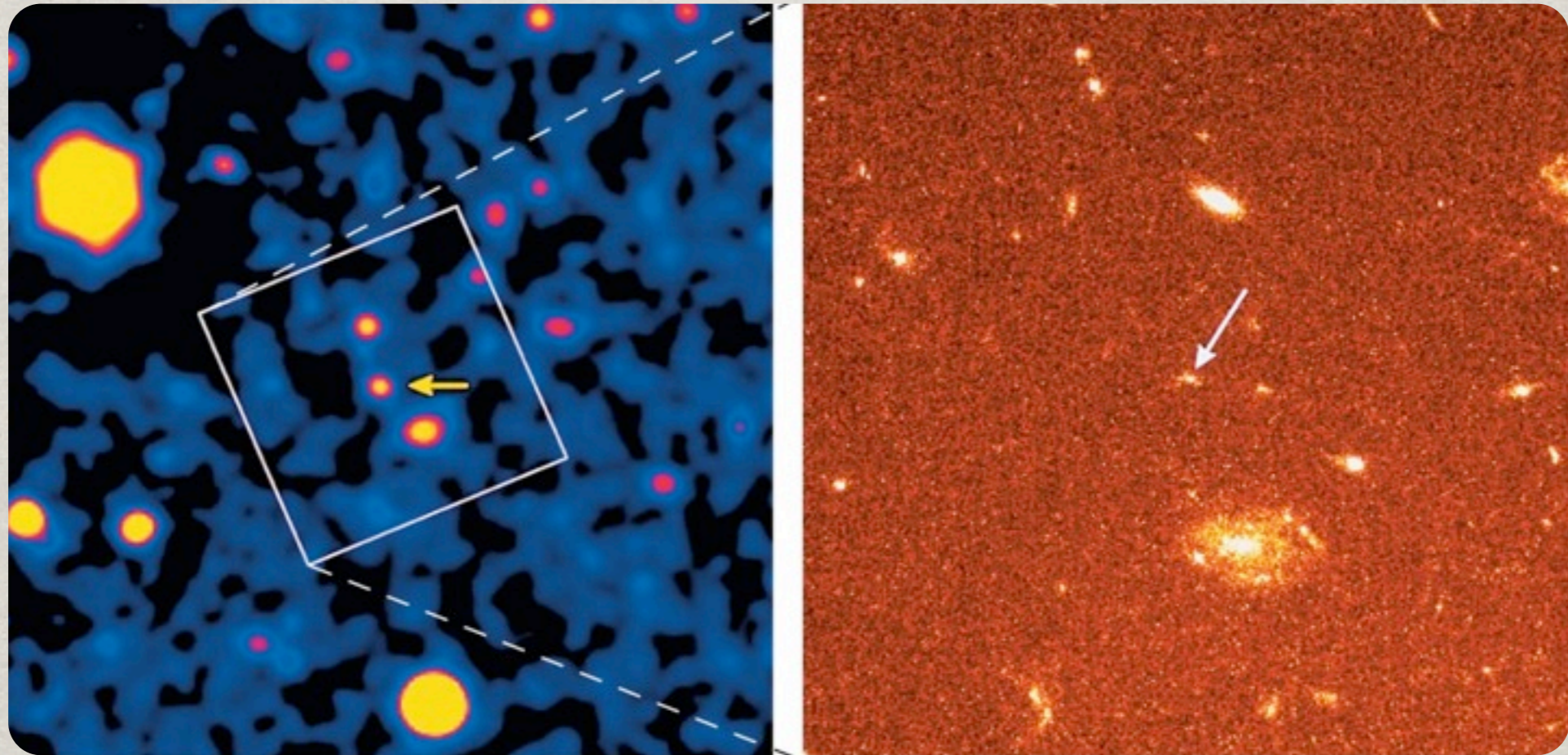
“**clumping**” of bursts anywhere, particularly not within the Milky Way.

Therefore, the bursts must originate from **outside our Galaxy**.



EXTRA-GALACTIC BURSTS

Distance measurements of **optical counterparts** of some gamma bursts have revealed coincidence with distant galaxies billions of light years away. Occasional spectral lines in the burst confirm the **large redshift**.



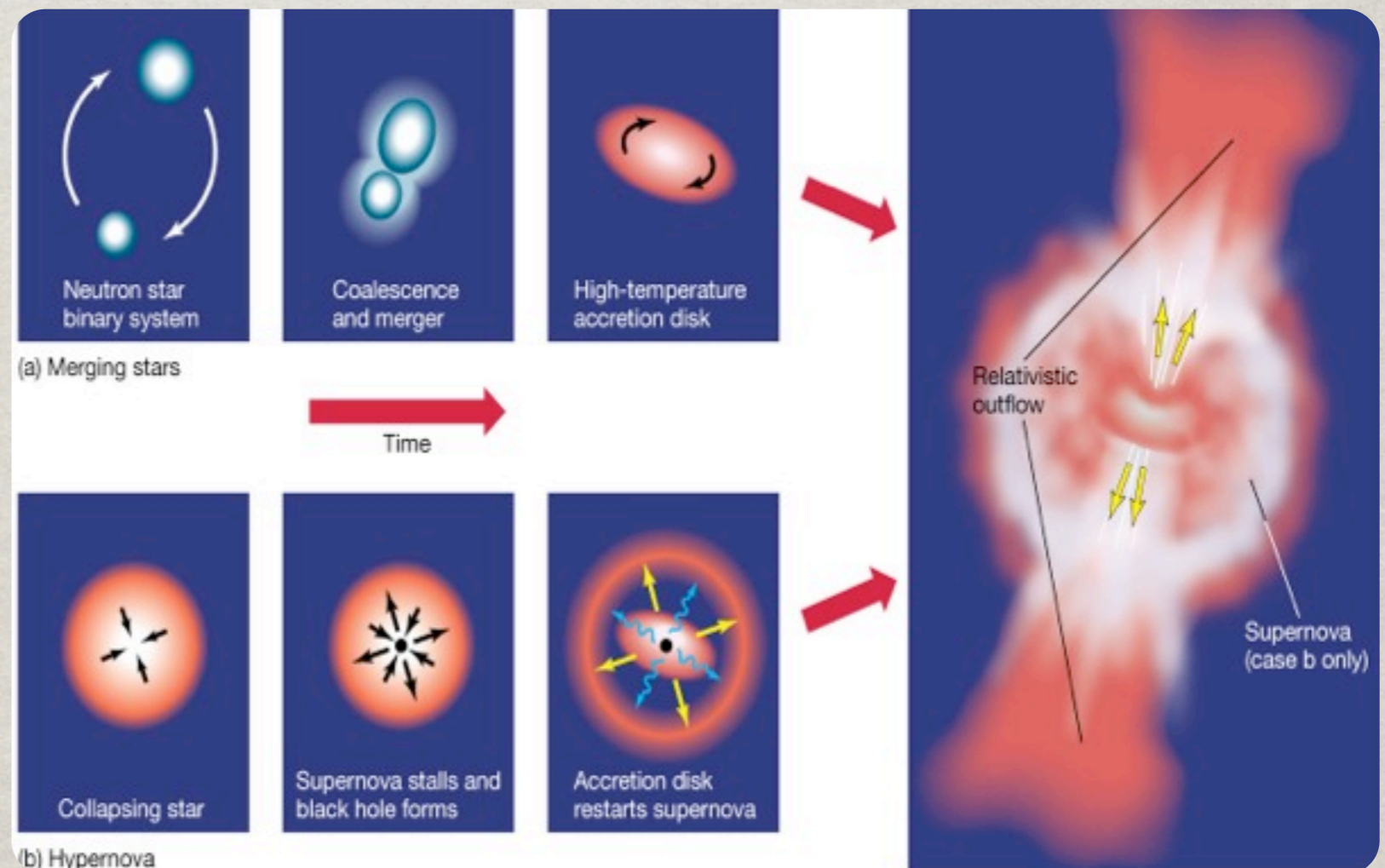
This requires **tremendous luminosity**, 10^{44} J, but confined to a narrow beam a **few degrees wide**.

GAMMA-RAY JETS

In both cases, the γ rays originate as **relativistic jets** created by accretion disks around newly formed **black holes**.

For the short bursts, binary **neutron stars merge** as a result of gravitational radiation, producing a black hole and jet.

Long bursts result when the black hole is the result of a **failed supernova**. The jets punch through the envelope of the star, later a disk wind drives off the envelope, causing a very bright supernova.

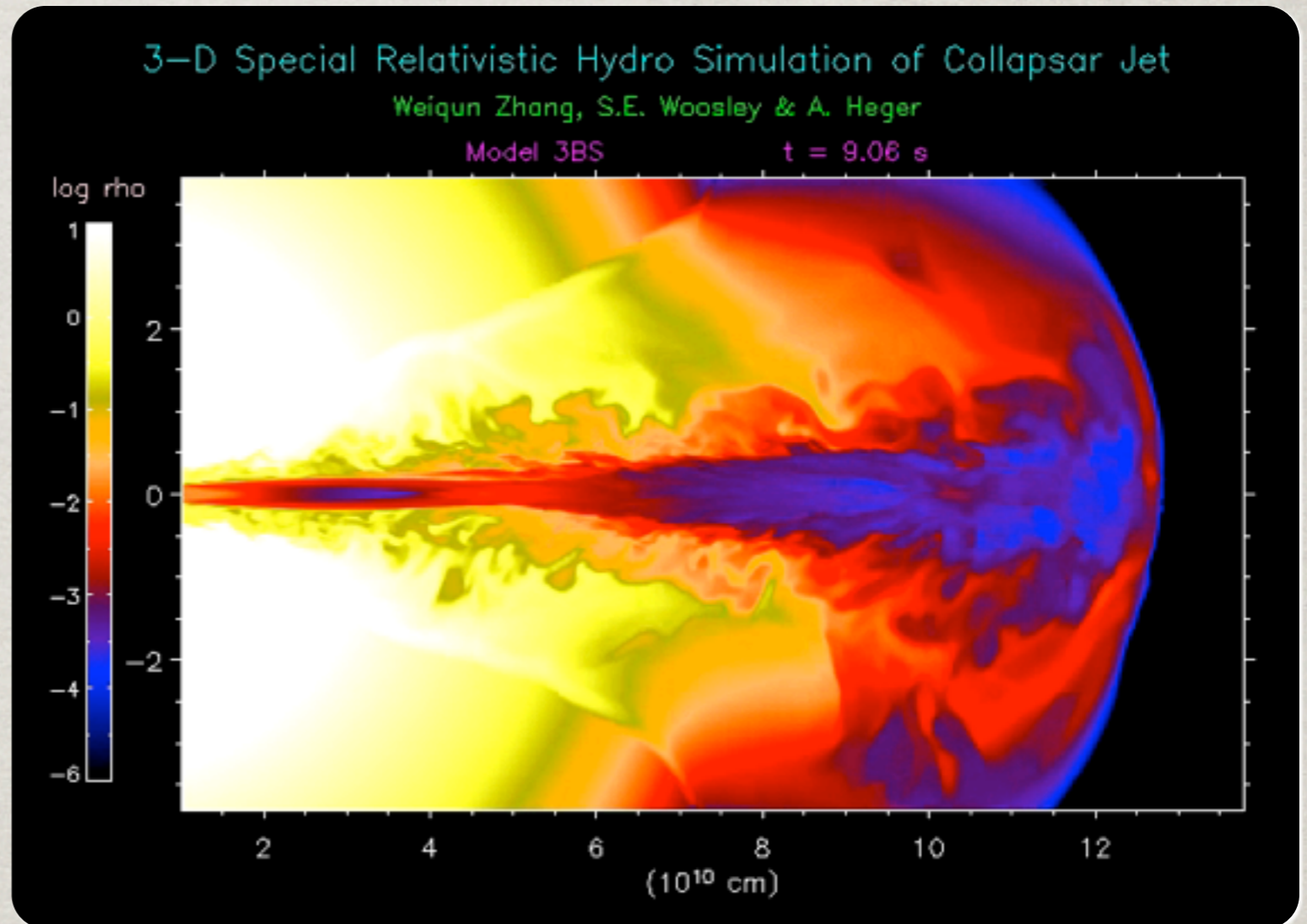


COLLAPSTAR OR MAGNETAR

Observations have tied some of these long GRBs to peculiar, **hyper-energetic** SN Ic.

One model is a **collapsar**, where an accretion disk forms around a newly-formed black hole in a **failed supernova** ($M > 30 M_{\odot}$), producing a jet which we see as the GRB.

Another model concerns the formation of a **magnetar**, an extremely magnetized neutron star.



SUMMARY

What role do star, supernovae, novae & X-ray bursts play in **cosmic nuclear evolution**?

*Core Collapse SN produce the **intermediate mass elements, O - Si- Ca**, and $\sim 1/2$ of **Iron Peak species**.

*Thermonuclear SN produce $\sim 1/2$ of the **Iron peak isotopes**.

*Stars produce **C, N & s-process**.

*Novae are likely responsible for odd mass isotopes of light elements like **C, N, O**.

Nuclear physics drives all of these events and their resulting nucleosynthesis.

