

Specifications for pressure stability for the gas regulation system.

❖ Gain variation as a function of the gas pressure

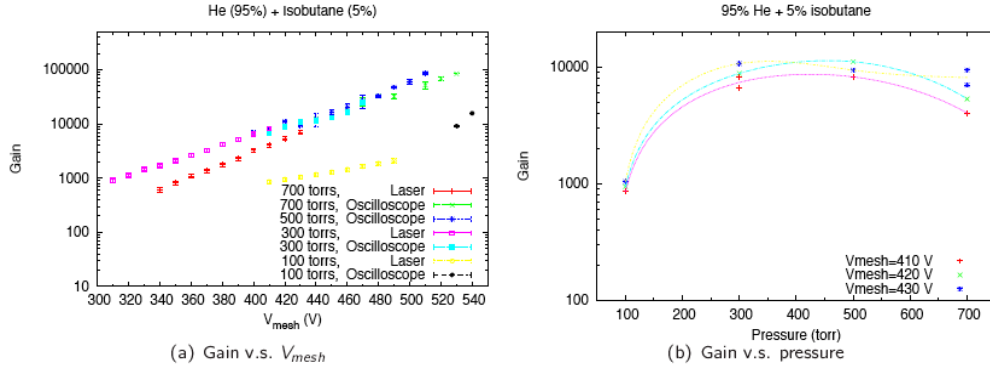


Figure 5: Gain as a function of pressure and V_{mesh} for the 95% He + 5% Isobutane mixture.

As one can see only in the low pressure region there is a strong dependence of gain.

Order of magnitude:

Factor 10 for $dp/p = 1$; resolution of micromegas about 1%; so dp/p giving a change of 1% corresponds to $dp/p = 0.1\%$ or $dp \sim 0.1 \text{ mbar}$

❖ Simulations for He+CO₂

Gas properties were simulated for the He+CO₂ mixture. We used the code [magboltz](#), where the Penning effect was taken into account. The ionization (aka Townsend) coefficient (α) [1/cm] was obtained and translated to the gain (g) using the equation $g = 2^{L\alpha}$, where L denotes the gap between the cathode and the anode. We employed 128 μm as L .

We estimated the gain fluctuation due to the uncertainties of (a) the total gas pressure (p_{tot}) and (b) the partial pressure of He (p_{He}), respectively.

➤ Fluctuation due to the total pressure

The change of p_{tot} primarily alters the Townsend coefficient α , but it also changes the gas density inside the gap. We can effectively take into account the latter effect by changing the gap size; $L(p_{tot}) = L_0 * (p_{tot} + \delta p_{tot})$.

In general, the Townsend coefficient α is nearly proportional to E/p_{tot} . On the other hand, $L(p_{tot})$ is proportional to p_{tot} . The number of ionization collision $L\alpha$ then remains almost unchanged because p_{tot} is cancelled between L and α , implying that the gain is rather stable with respect to p_{tot} .

The gain was calculated for different p_{tot} of 1, 0.5 and 0.2 atm, respectively (Fig. 1). The fluctuation of the gain was estimated for the respective cases by

changing the total pressure by +1%. As shown in Fig. 2, the resultant fluctuations are less than +-5%.

➤ Fluctuation due to the partial pressure of He

The simulated gains for 10% CO₂ and 11% CO₂ are shown in Fig. 1. The gain fluctuation was deduced based on the following equation;

$$\delta g/g = (1 - g(10\% \text{ CO}_2)/g(11\% \text{ CO}_2)) \text{ at a given } E/p_{\text{tot}}$$

The results are shown in Fig. 2. For the 1% accuracy of p_{He} with respect to p_{tot} , the magnitude of the gain fluctuation is about 30% at maximum.

Hence the gain fluctuation is dominated by the accuracy of the partial pressure of He. If the gain fluctuation is to be less than 1%, the partial pressure should have an accuracy of 0.03% with respect to the total pressure.

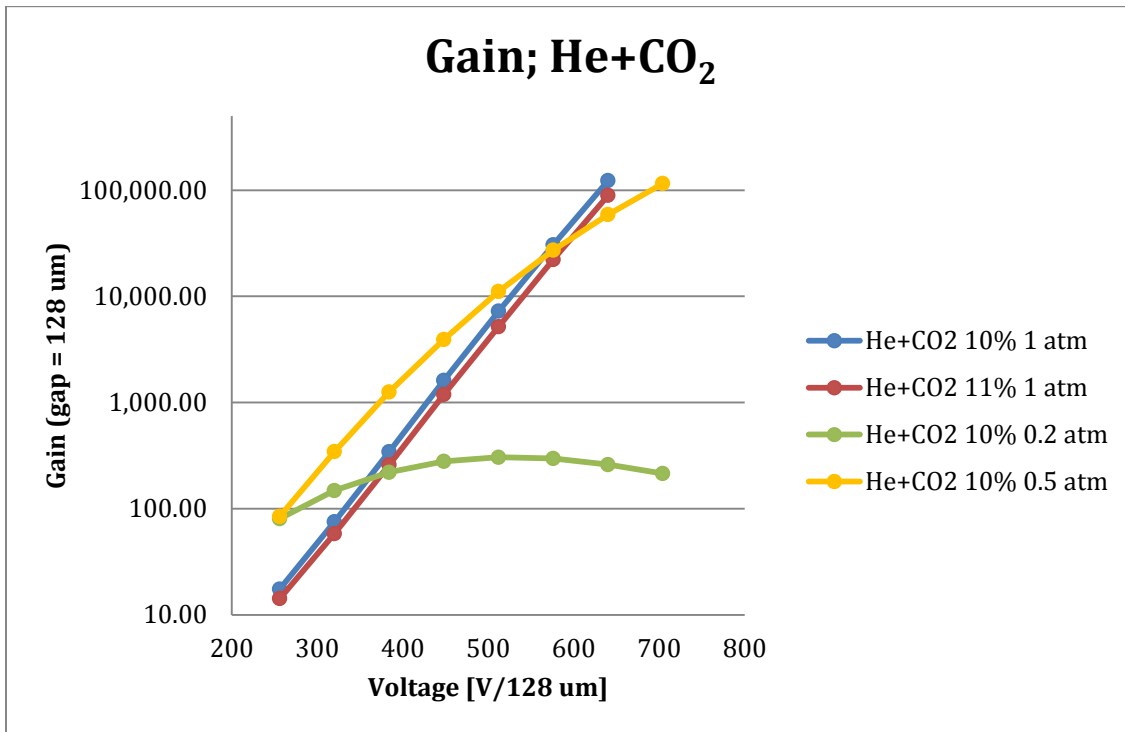


Figure 1: Simulated gain as a function of the field gradient.

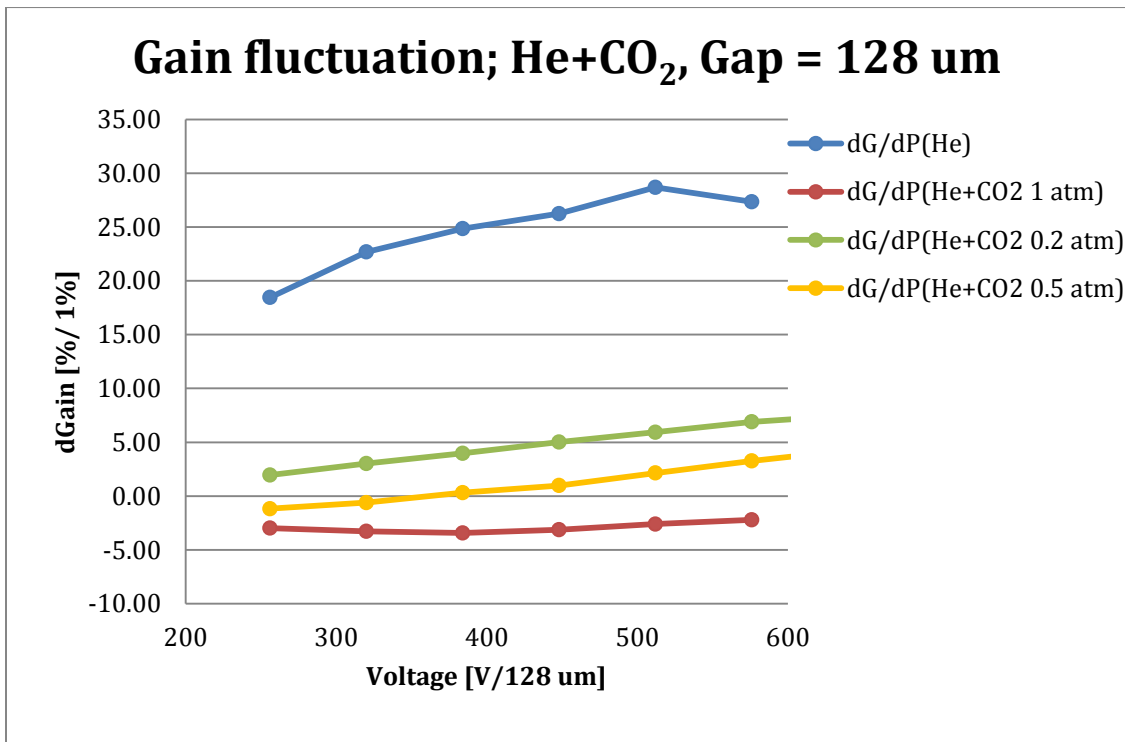
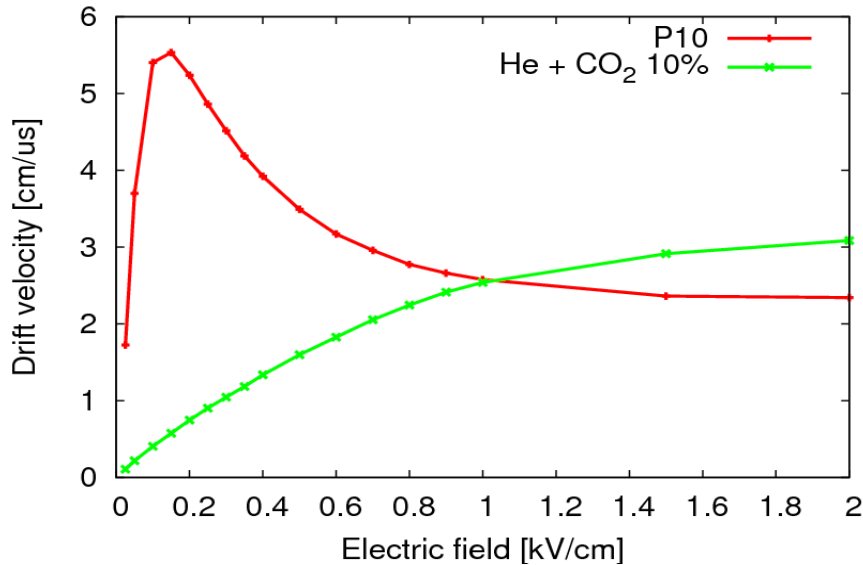


Figure 2: Gain fluctuation for (Blue) the partial pressure of He and (Red) the total pressure of He+CO₂

❖ Drift-time as a function of the pressure

Variations of drift-time as a function voltage/pressure will be monitored by the Laser system. As example we may consider the figure below:



In the case of P10 one will work the maximum, where dependence is small. For the He+CO₂, we want to work in the region 1kV/cm. For order of magnitude we have:

To good enough approximation, the drift velocity is proportional to the V, this is inversely proportional to p; in the middle of the detector, there drift time is ~ 50 cm/2(cm/ μ s) = 25 μ s; if we want a stability of 100 μ m, the drift time must be stable to $0.1/500=2/10,000 = 0.02\%$; so $dv/v \sim 2/10,000$ or $dp/p = 2/10,000$. if $p = 100$ mbar, $dp = 2 \cdot 10^{-2}$ mbar = 20 μ bar; this is close to the value given by Ana.

❖ Simulations for He+CO₂

The simulated electron drift velocity is plotted for different mixtures of 10% CO₂ and 11% CO₂, respectively, in Fig. 3. The fluctuation of the velocity was calculated for the case wherein the partial pressure of He or the total pressure of He+CO₂, respectively, changes by 1%. The results are shown in Fig. 4.

For the pressure accuracy of 1%, the magnitude of the velocity fluctuation is about $\pm 2\%$ at maximum around $E = 1$ kV/cm/atm, which is again governed by the accuracy of the partial pressure of He. A pressure accuracy of 0.01% yields a velocity fluctuation of less than $\pm 0.05\%$.

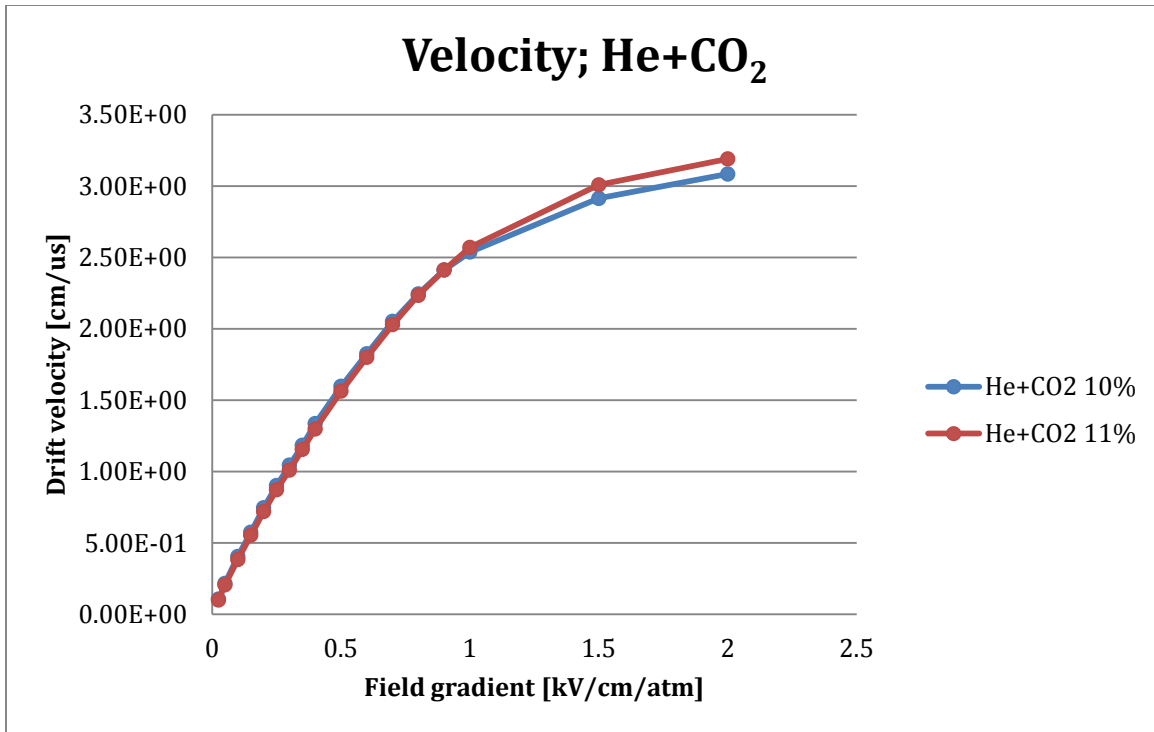


Figure 3: Simulated velocity of the electron drift.

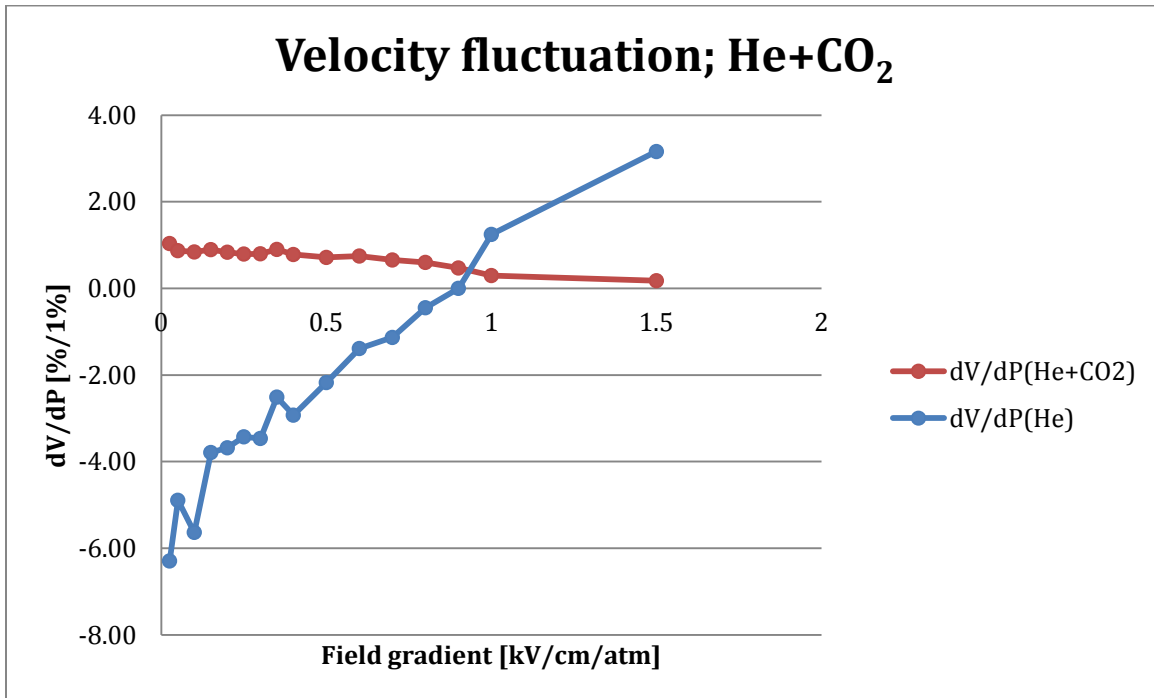


Figure 4: The fluctuation of the electron drift velocity for (Blue) the partial pressure of He and for (Red) the total pressure of He+CO₂

❖ Conclusion

The gain fluctuation is dominated by the partial pressure of He (p_{He}). A 0.03% accuracy of p_{He} , with respect to the total pressure, gives a gain fluctuation of 1%.

The drift velocity is dominated again by p_{He} . Given a 0.01% accuracy for p_{He} , the magnitude of the velocity fluctuation is less than $\pm 0.02\%$ around $E = 1 \text{ kV/cm/atm}$. The drift velocity will be monitored by a laser.

It should be examined if the pressure stability of $100 \mu\text{bar}$ (0.1% of 100 mbar) would result in a considerable cost decrease compared to that of $10 \mu\text{bar}$ (0.01% of 100 mbar). Otherwise one should adopt the value of $10 \mu\text{bar}$ as given by Ana.

❖ Appendix

The simulated gain was compared to the result obtained in the test bench of NSCL. The simulation well agrees with the measurement.

