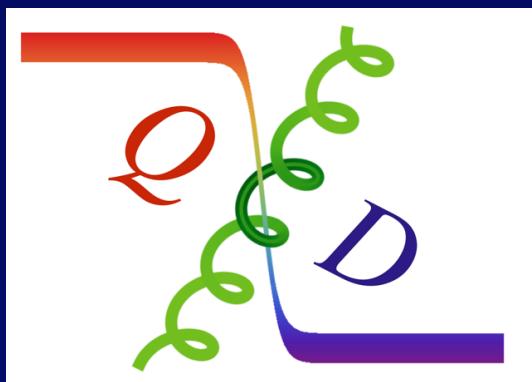


Proton Spin and Mass Decompositions and Neutrino-Nucleon Scattering from Lattice QCD

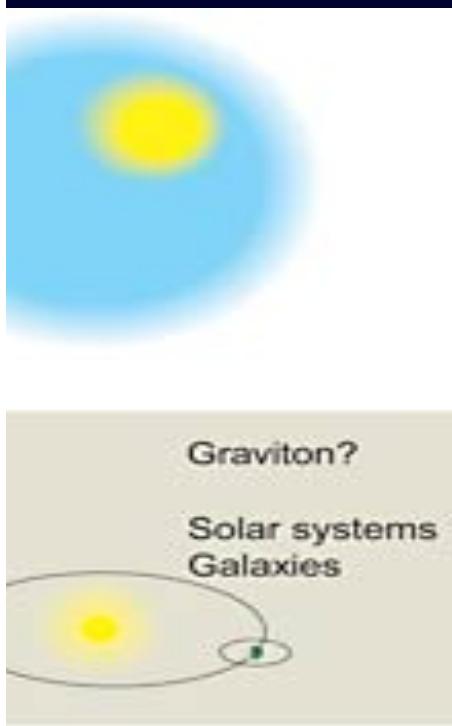
- Synopsis of Lattice QCD
- Proton Spin Decomposition -- Quark and Glue Spin, Orbital Angular Momentum
- Proton Mass Decomposition – Quark Condensate, Quark and Glue energies, and Trace Anomaly
- Neutrino-Nucleon Scattering -- Hadronic Tensor
 - Inverse Problem

χ QCD Collaboration



Michigan State University,
Feb. 11, 2020

Lattice QCD

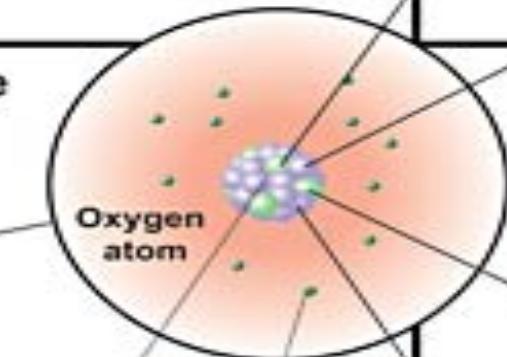
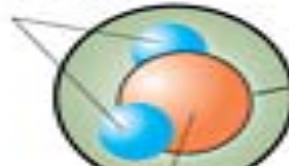


Gravity Force



Electromagnetic force

Hydrogen atom



Protons and Neutrons

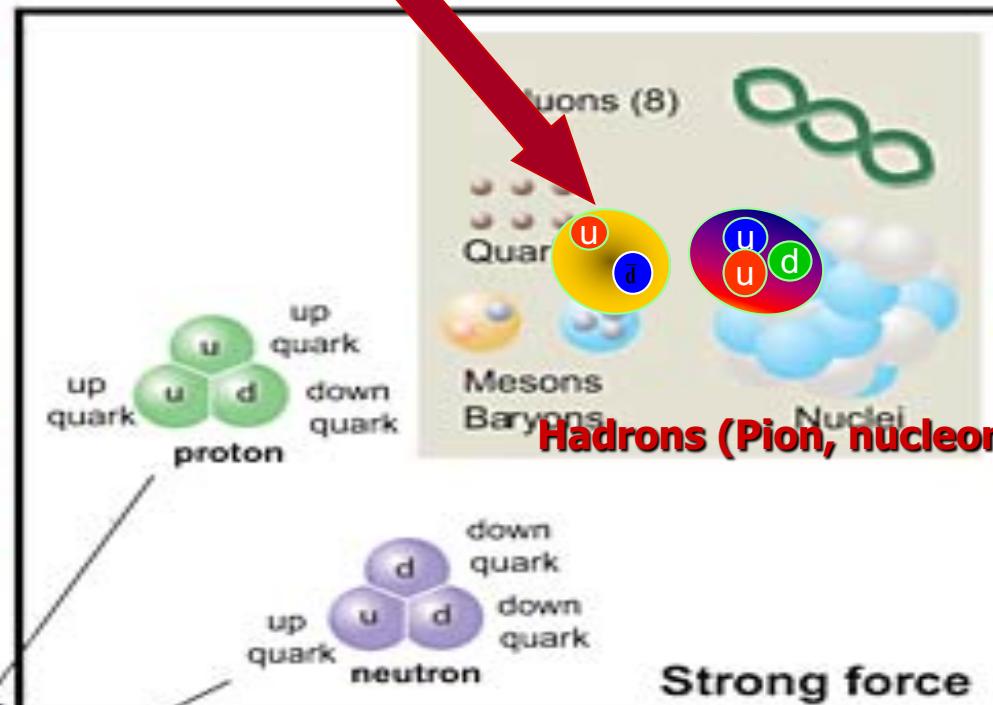
Electron

Matter molecule

Oxygen atom

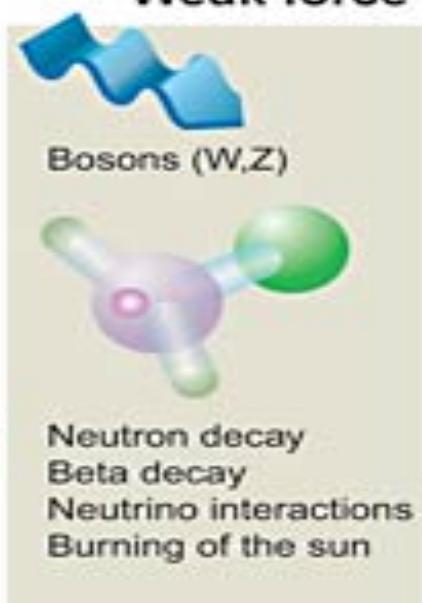
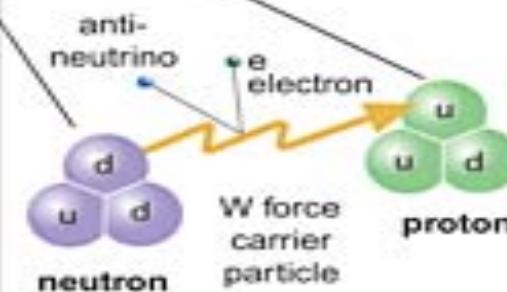


Atoms
Light
Chemistry
Electronics



Strong force

Weak force

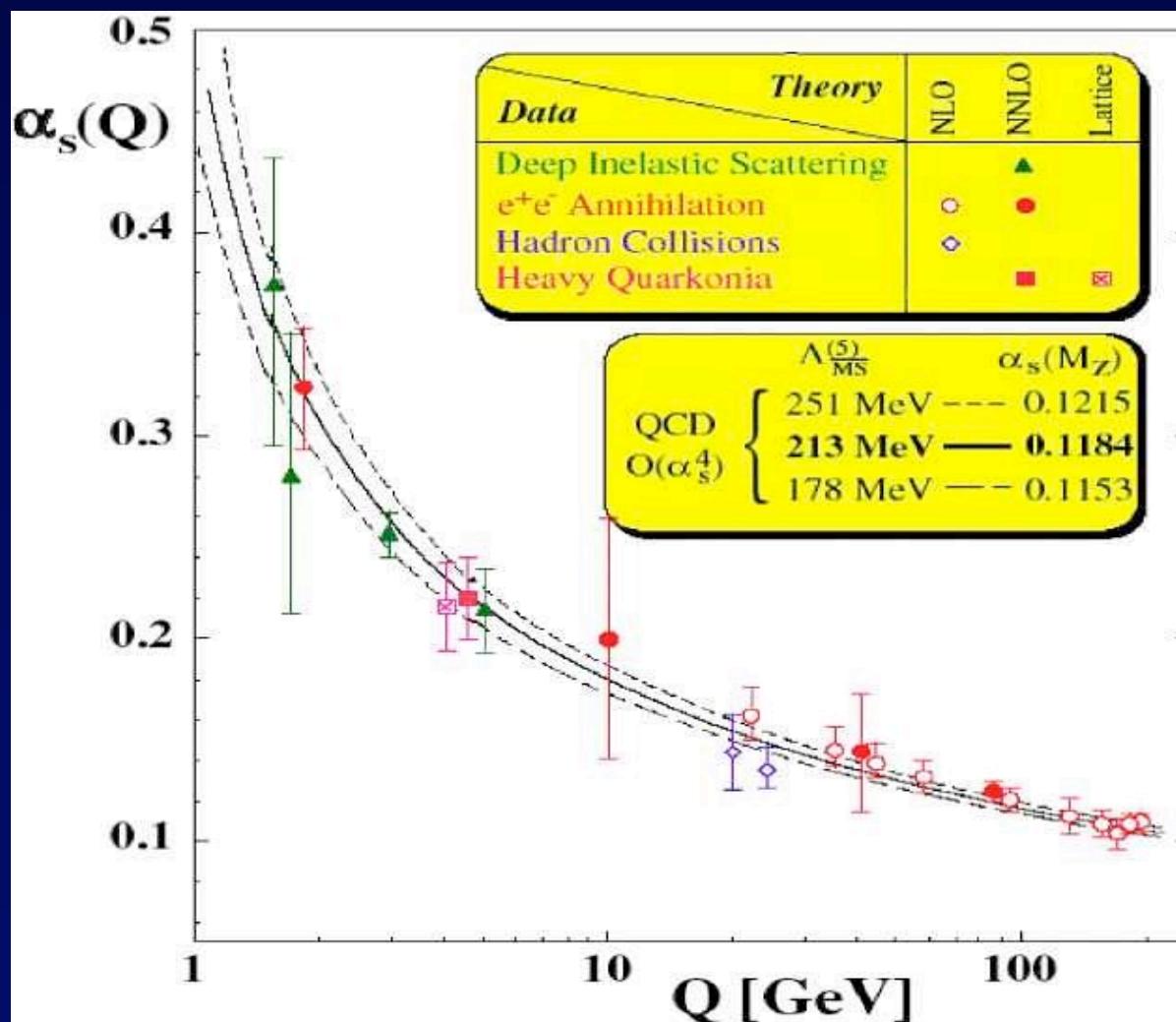


Theory of Strong Interactions : QCD

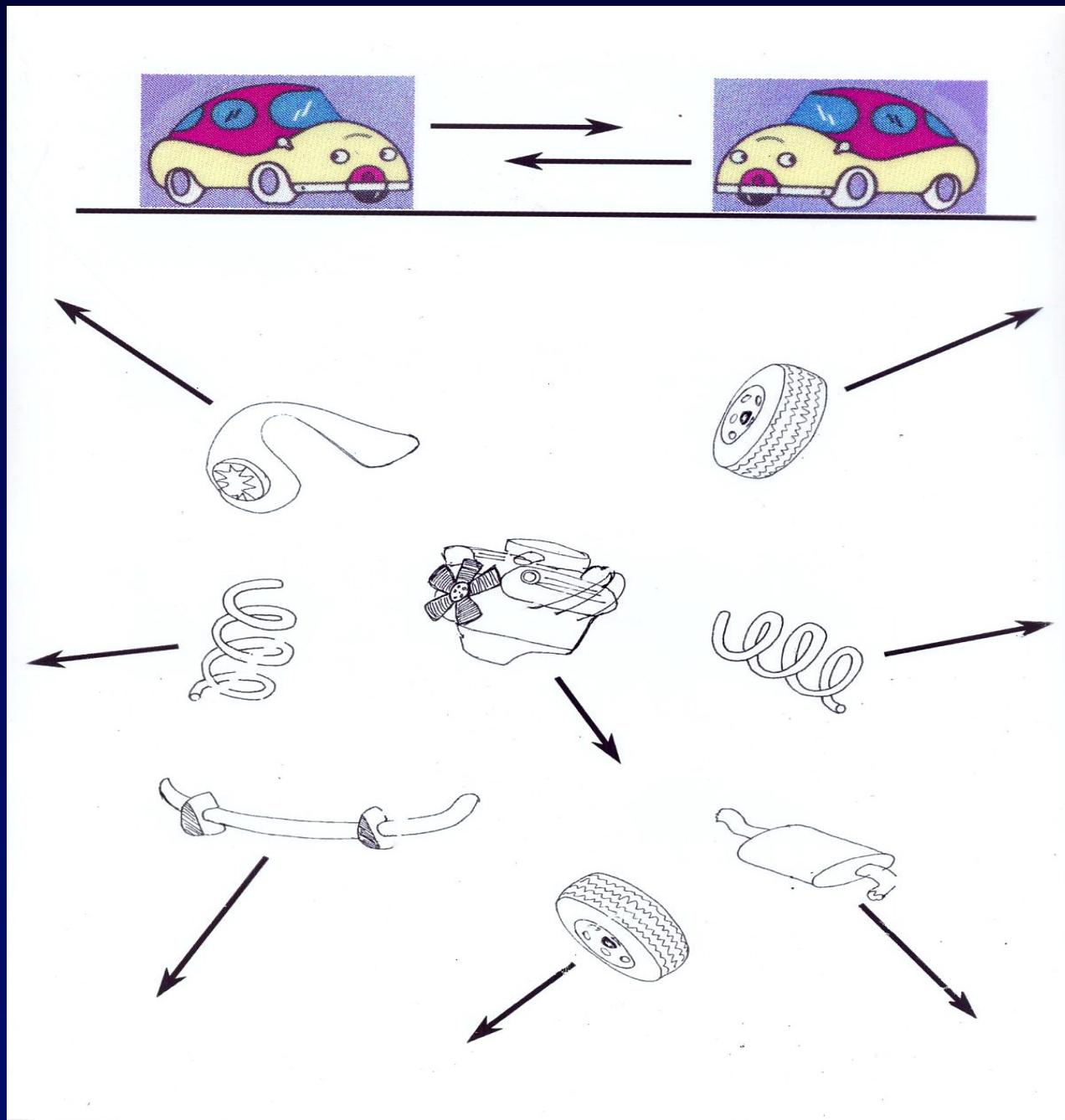
The force between quarks becomes weak for small quark separations--- a phenomenon known as

Asymptotic Freedom.

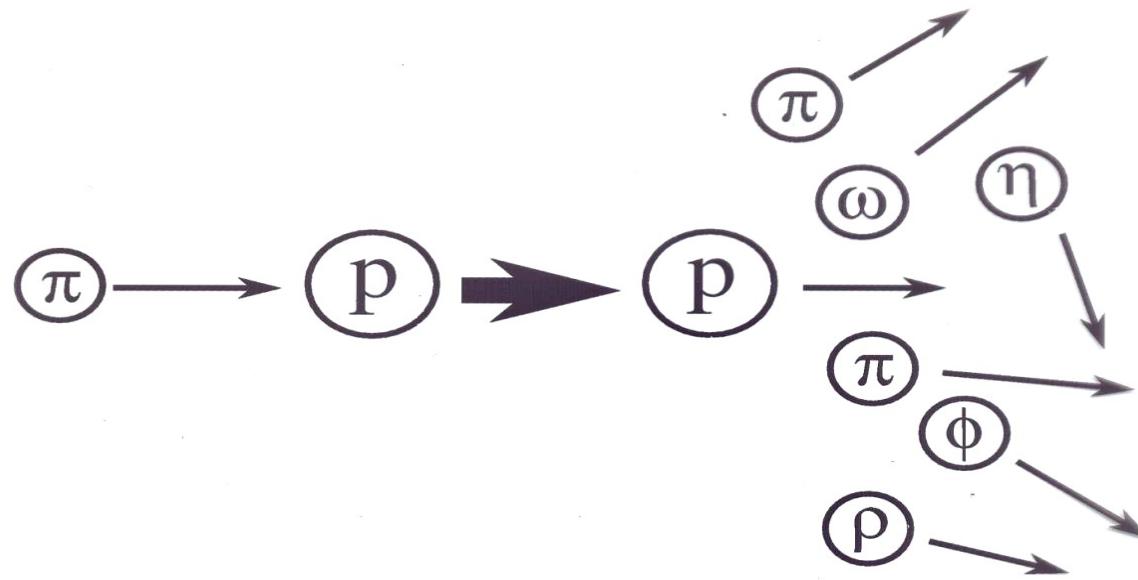
..... *Nobel Prize 2004*



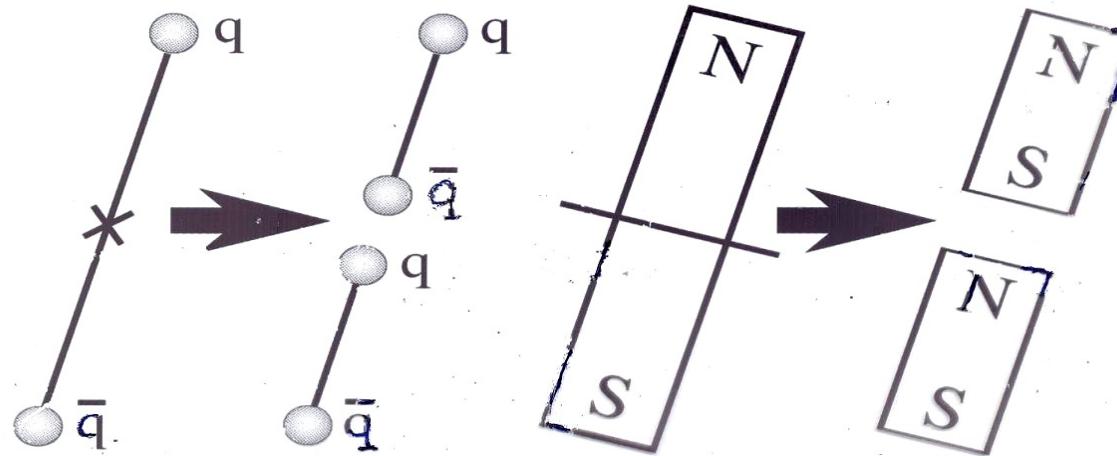
Confinement

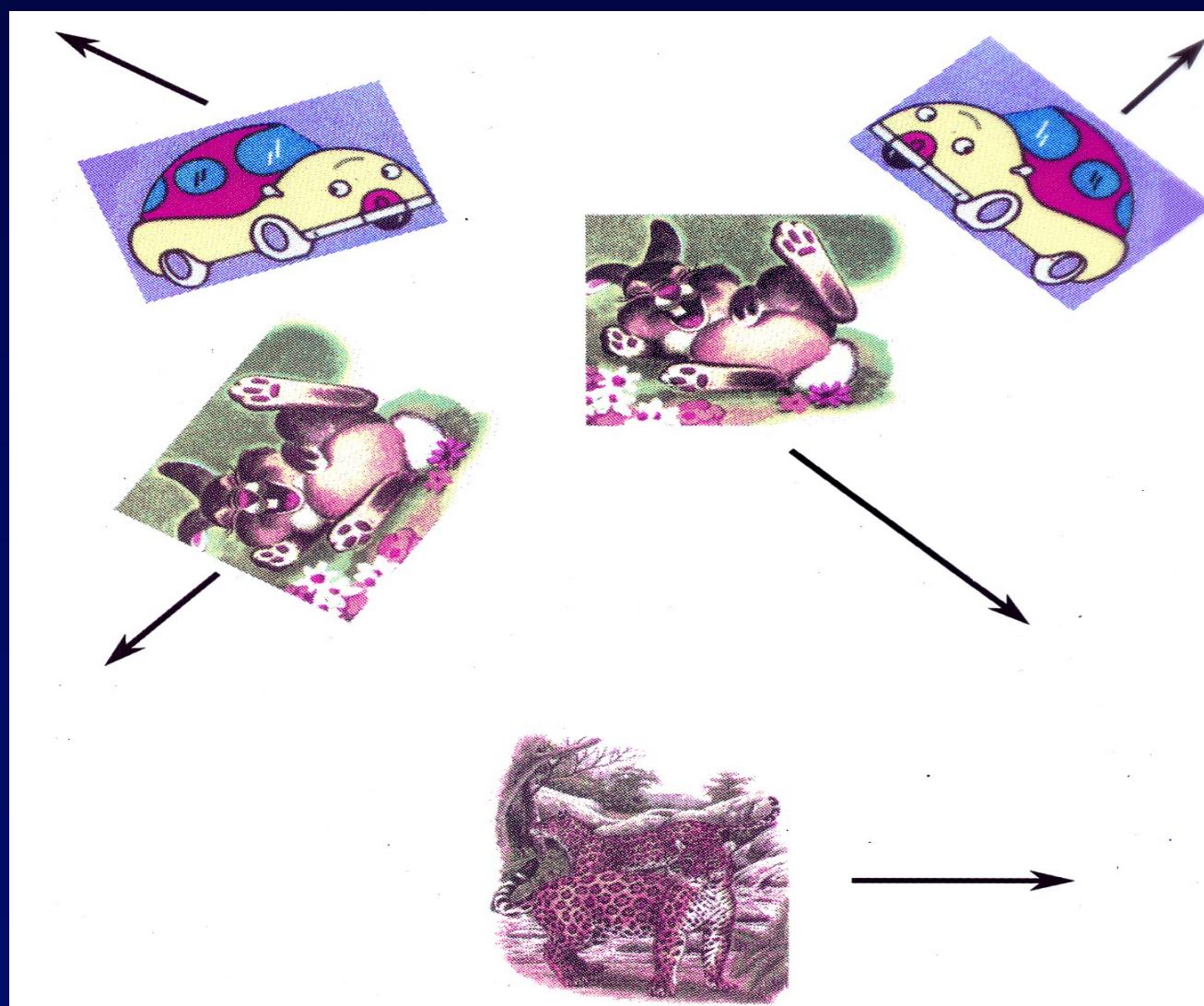
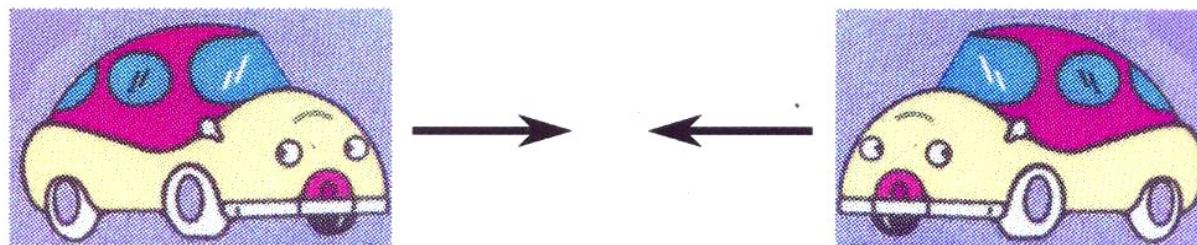


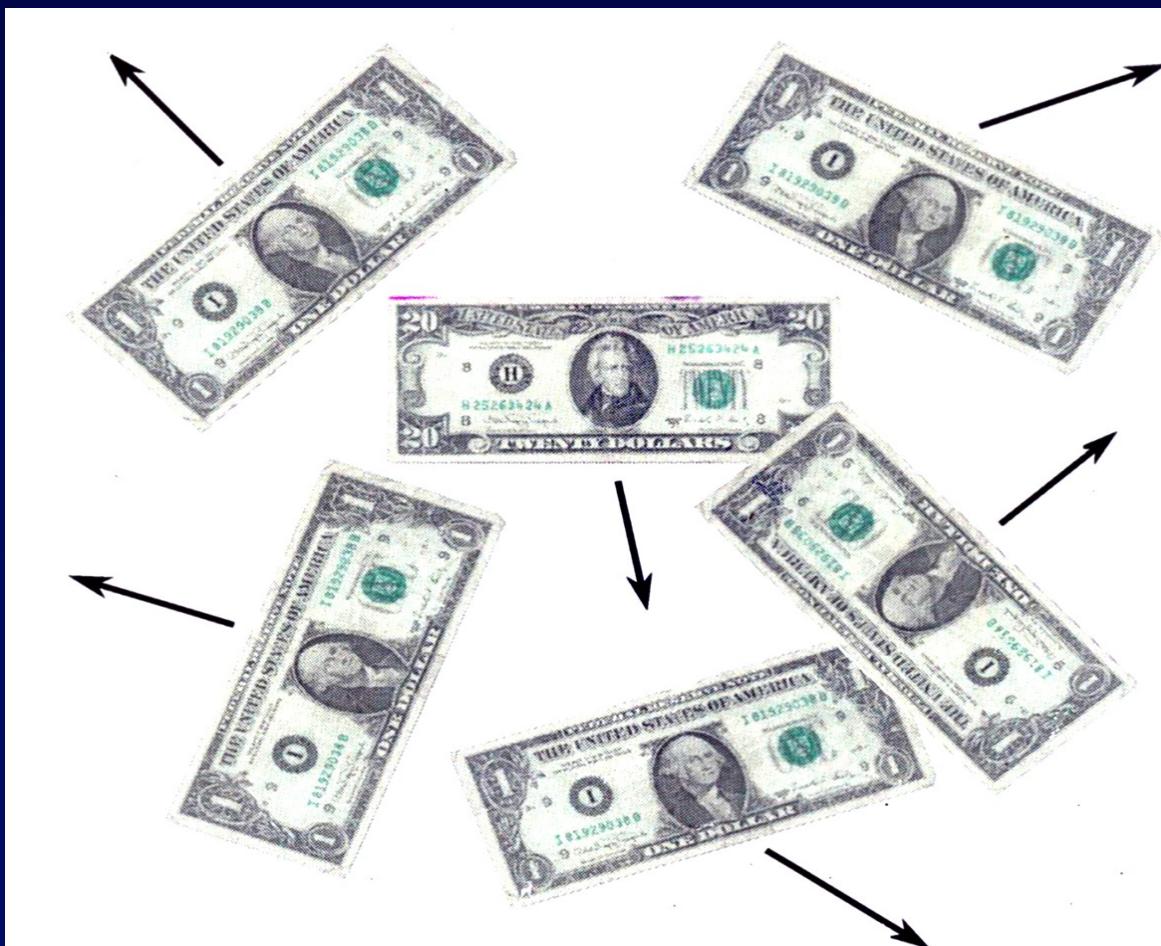
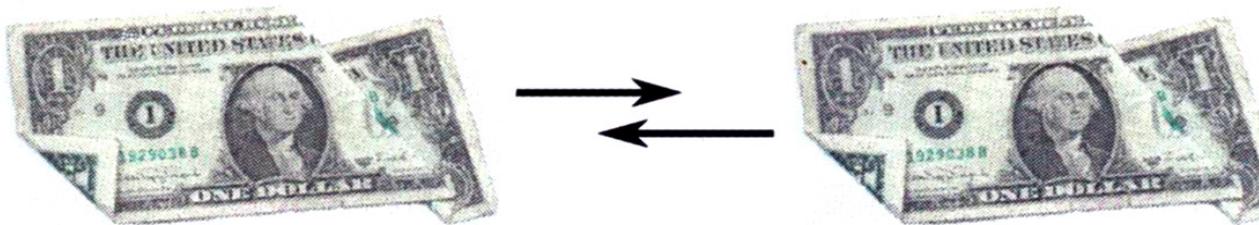
Quark Confinement



Analogy:





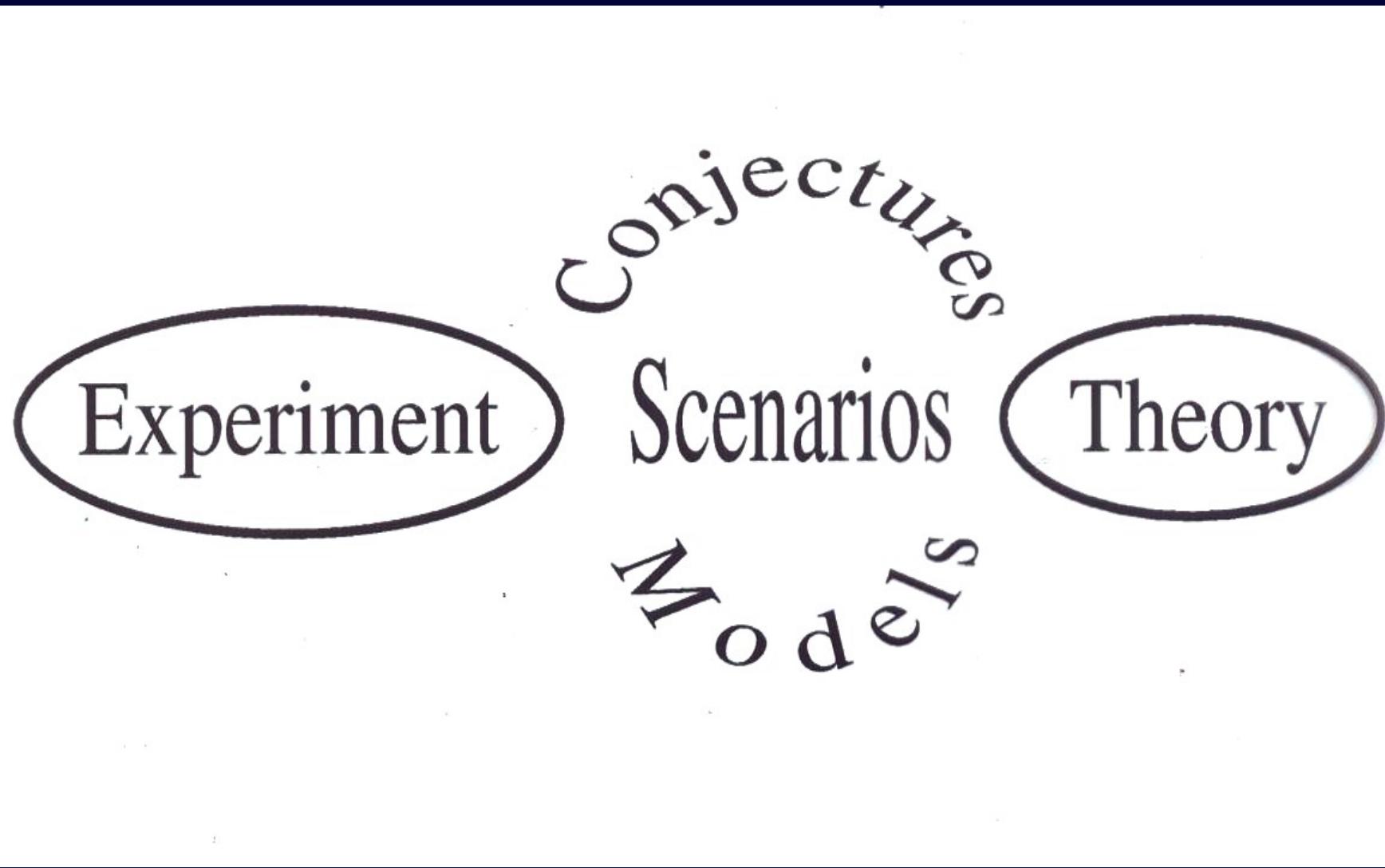


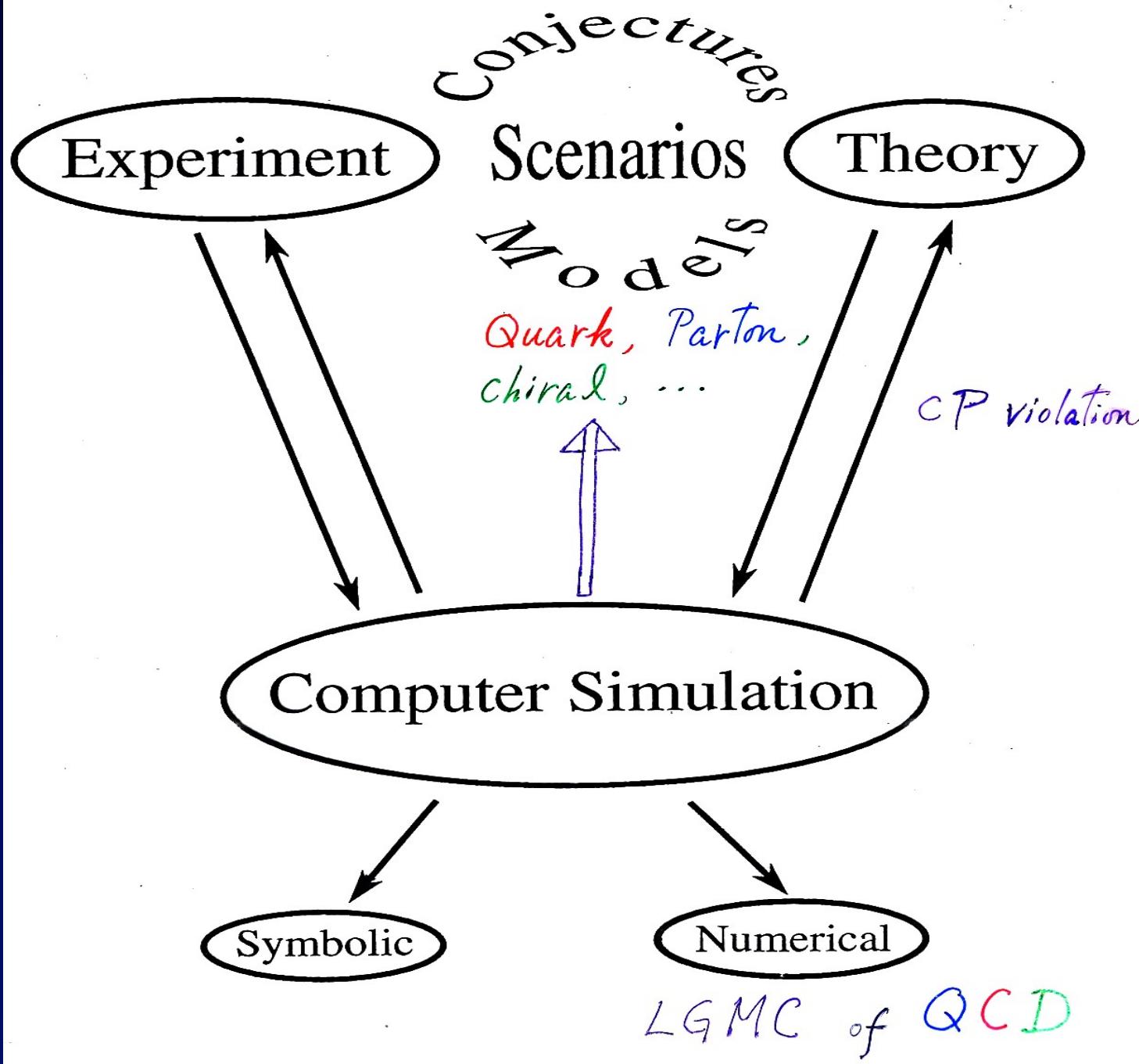
"The Understanding cannot *See*.
The Senses cannot *Think*.
By their union only can *Knowledge*
be produced."

— Immanuel Kant

"*Imagination* is more
important than *Knowledge*."

— Albert Einstein

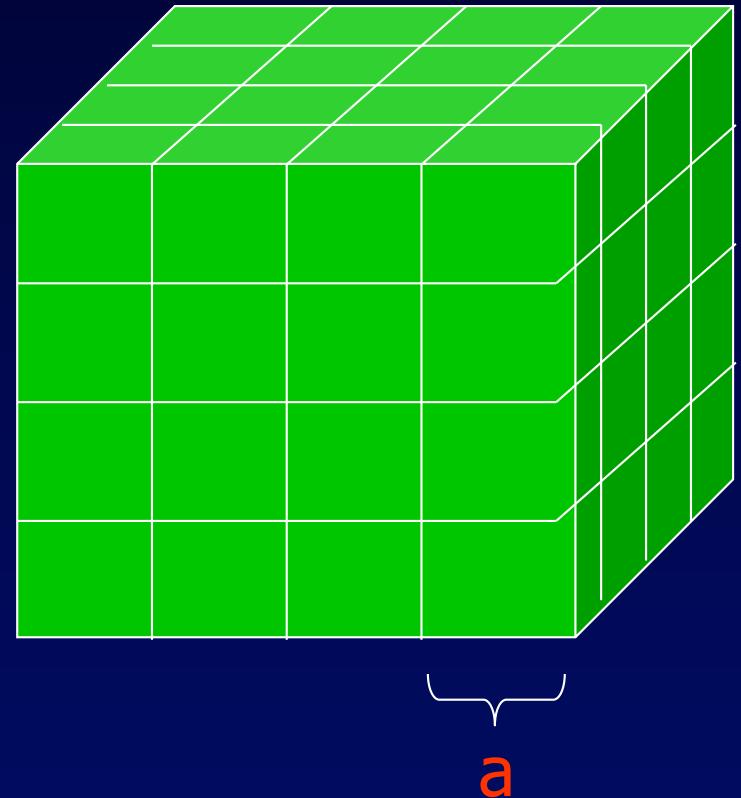




Lattice QCD

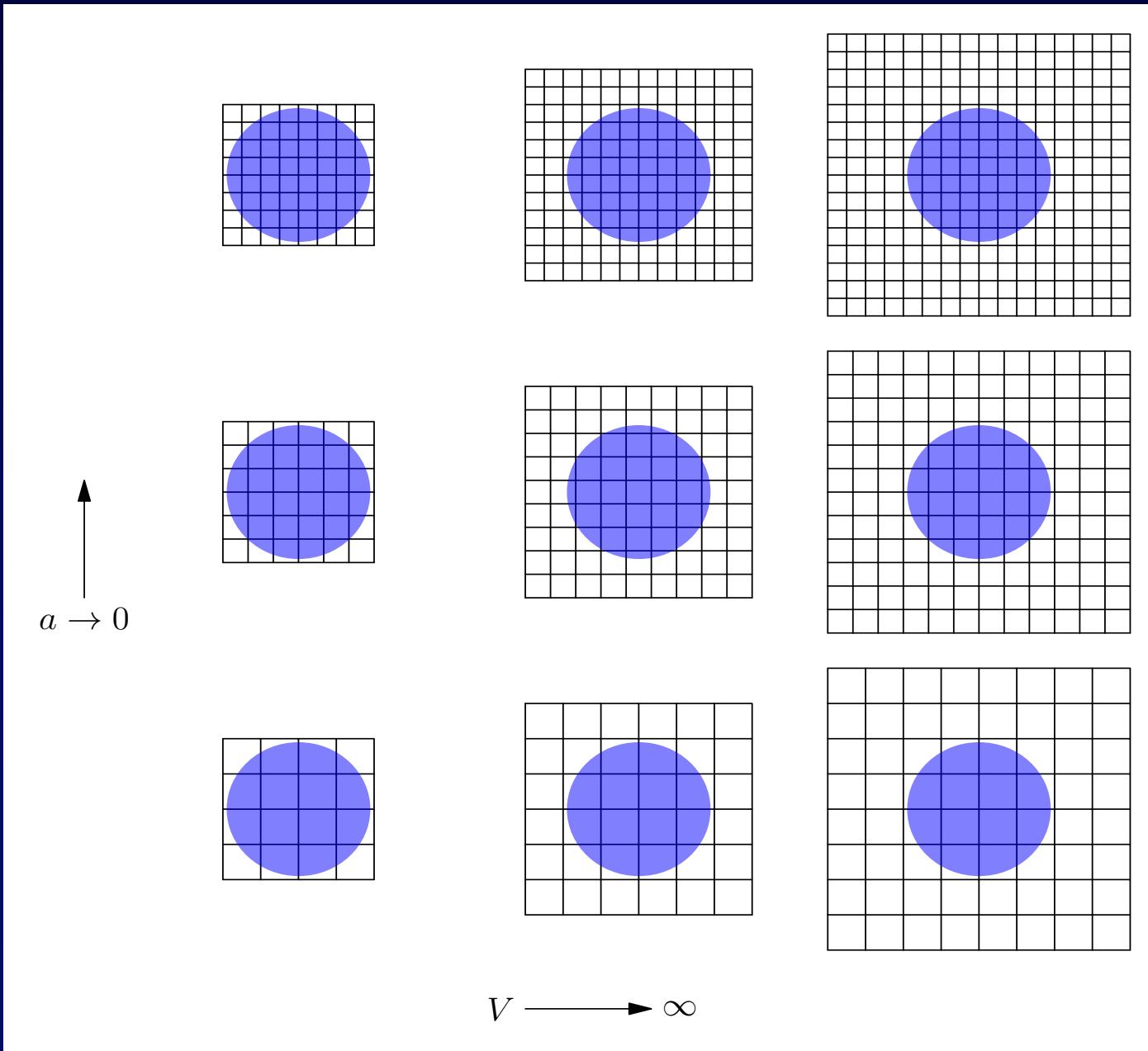
Why Lattice?

- Regularization
 - Lattice spacing a
 - Hard cutoff, $p \leq n/a$
 - Scale introduced (dimensional transmutation)
 - Renormalization
 - Perturbative
 - Non-perturbative
- Regularization independent Scheme
Schroedinger functional
Current algebra relations
- Numerical Simulation
 - Quantum field theory \rightarrow classical statistical mechanics
 - Monte Carlo simulation in Euclidean space (importance sampling)



$$e^{-S_G} \det M \geq 0$$

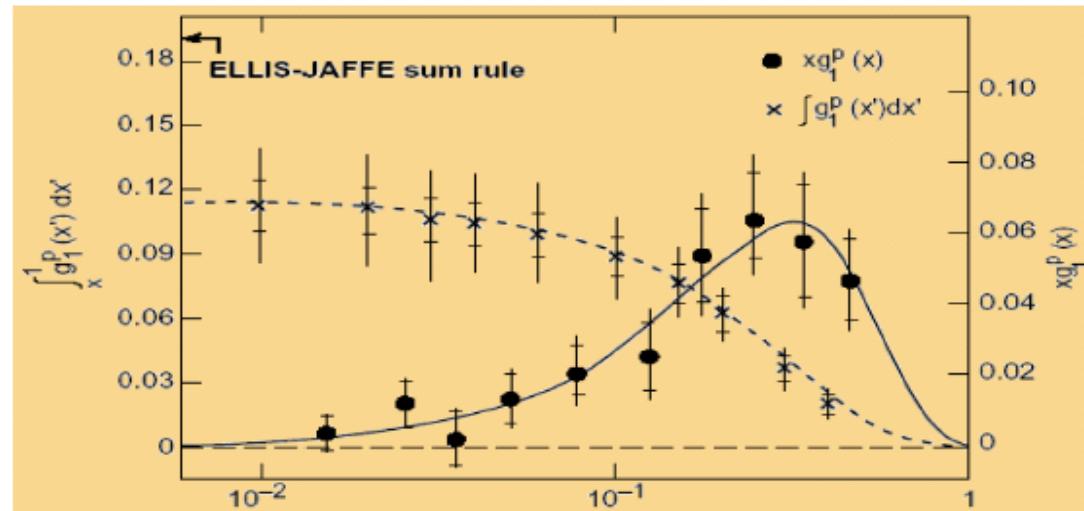
Continuum and Infinite Volume Limits at Physical Pion Mass (Systemic Errors)



Where does the spin of the proton come from?

Thirty years since the “spin crisis”

□ EMC experiment in 1988/1989 – “the plot”:



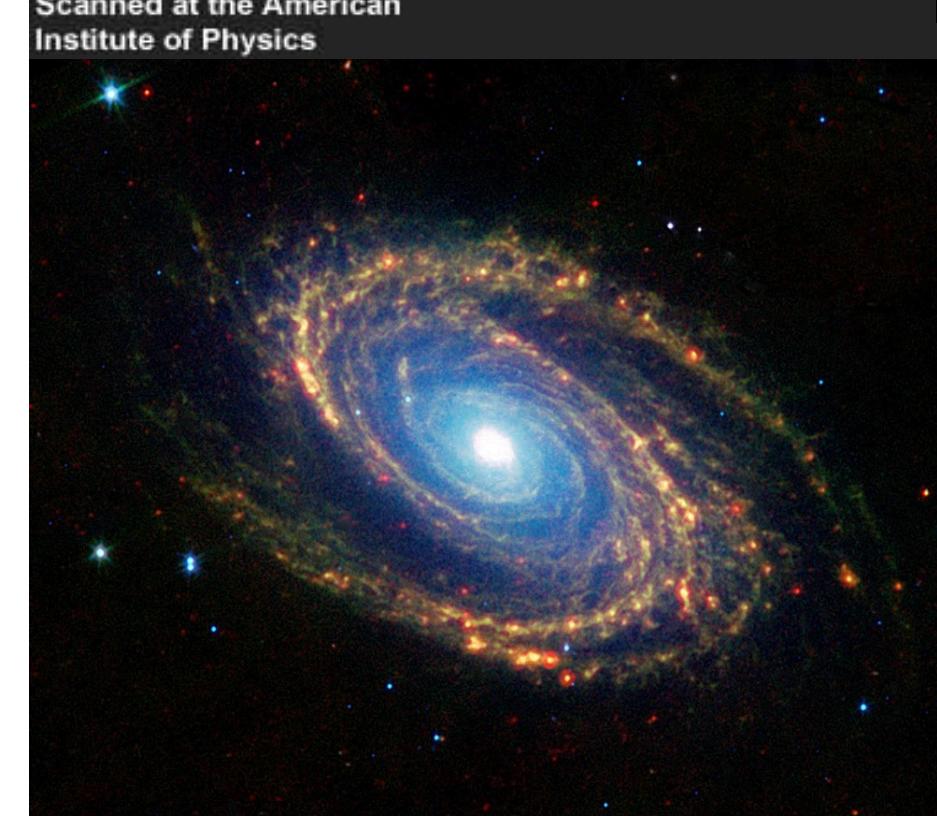
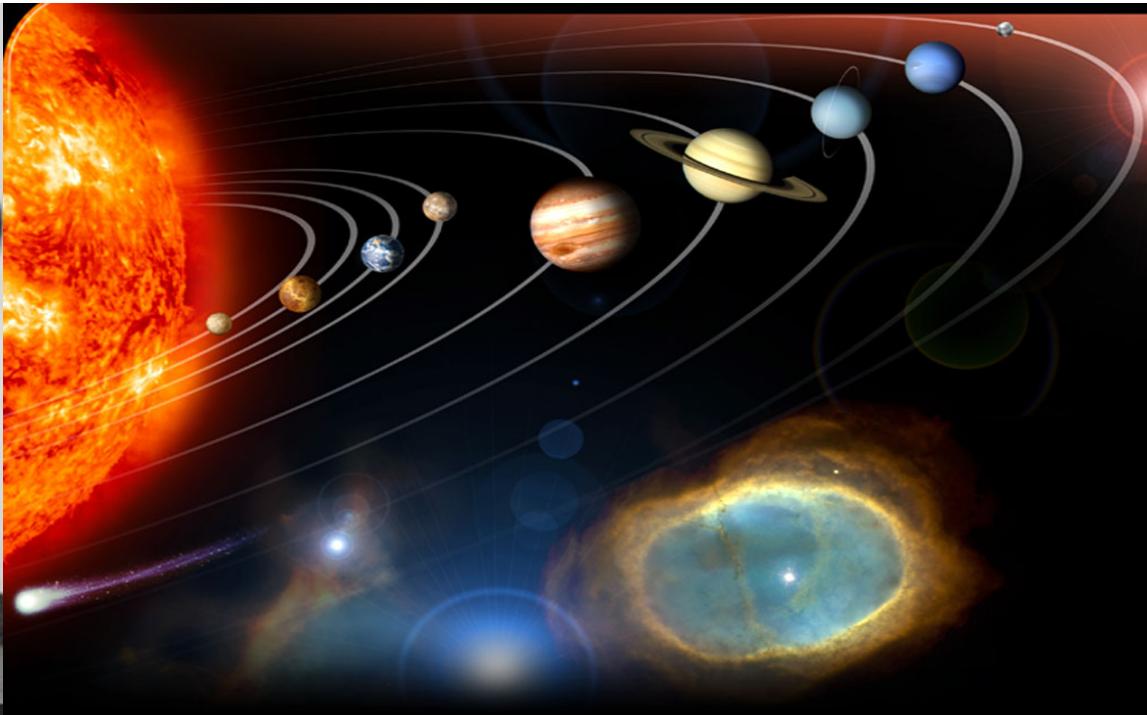
$$g_1(x) = \frac{1}{2} \sum e_q^2 [\Delta q(x) + \Delta \bar{q}(x)] + \mathcal{O}(\alpha_s) + \mathcal{O}(1/Q)$$

$$\Delta q = \int_0^1 dx \Delta q(x) = \langle P, s_{\parallel} | \bar{\psi}_q(0) \gamma^+ \gamma_5 \psi_q(0) | P, s_{\parallel} \rangle$$

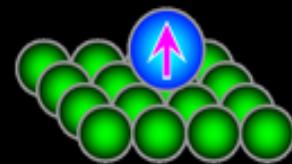
□ “Spin crisis” or puzzle: $\Delta\Sigma = \sum_q \Delta q + \Delta \bar{q} \sim 0.3$



Scanned at the American
Institute of Physics

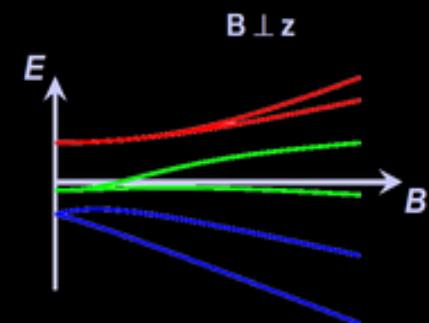
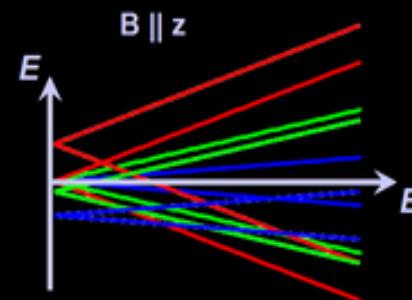


Anisotropy at a surface



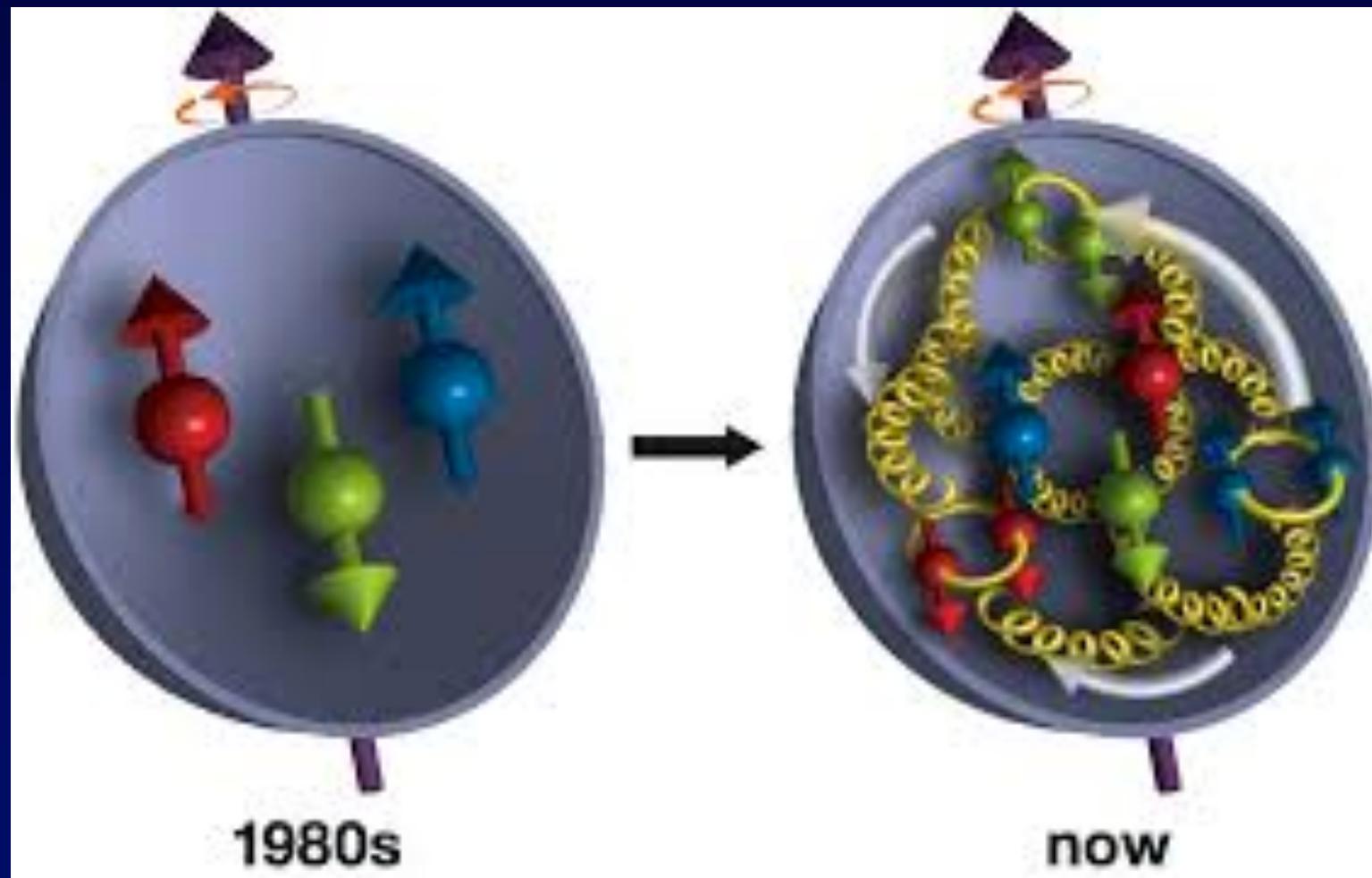
- Free atomic spin is rotationally invariant: all spin orientations are degenerate.
- Loss of rotational symmetry breaks degeneracy of spin orientations.

$$H = -g\mu_B \overset{\omega}{B} \cdot \overset{\omega}{S} + D S_z^2$$



Magnetic field dependence varies with angle of magnetic field.

Picture from quark model to QCD

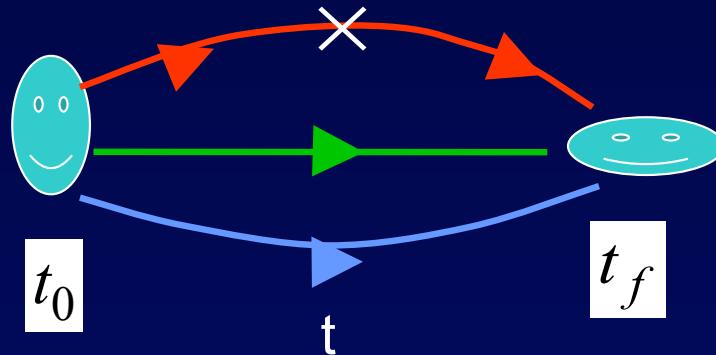


Lattice Calculations of Quark and Glue Spins

- Quark and Glue Momentum and Angular Momentum in the Nucleon

$$\sum_i w_i e^{-E_i t}$$

$$(\bar{u} \gamma_\mu D_\nu u + \bar{d} \gamma_\mu D_\nu d)(t)$$

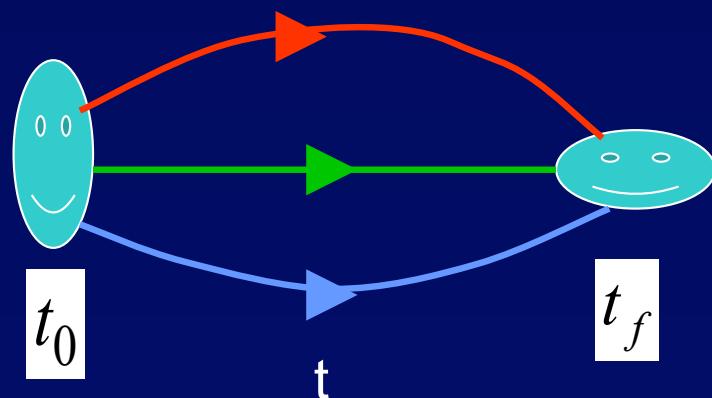


Connected
insertion (CI)

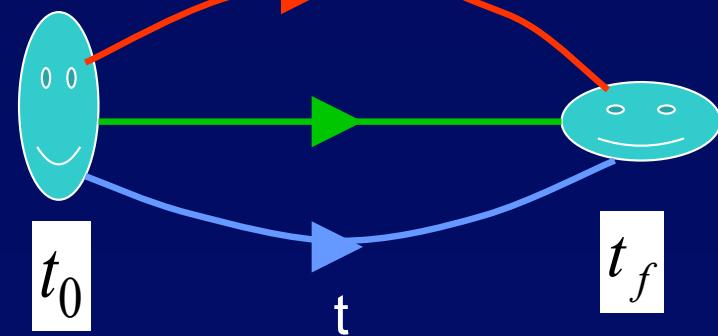
Disconnected
insertion (DI)



$$\bar{\Psi} \gamma_\mu D_\nu \Psi(t)(u, d, s)$$



$$F_{\mu\alpha} F_{\nu\alpha} - \frac{1}{4} \delta_{\mu\nu} F^2$$



Quark Spin Components $\overline{\text{MS}}$ (2 GeV)

g_A	$\Delta(u+d)$ CI	$\Delta(u/d)$ DI	Δs	Δu	Δd	$\Delta u - \Delta d$ (g_A^3)	$\Delta \Sigma$
PNDME			-0.053 (8)	0.777 (25)(30)	-0.438 (18)(30)	1.128 (27)(30)	0.286 (62)(72)
C. Alexandrou	0.598 (24)(6)	-0.077 (15)(5)	-0.042 (10)(2)	0.830 (26)(4)	-0.386 (16)(6)	1.216 (31)(7)	0.402 (34)(10)
χ QCD	0.580 (16)(30)	-0.070 (12)(15)	-0.035 (6)(7)	0.847 (18)(32)	-0.407 (16)(18)	1.254 (16)(30)	0.405 (25)(37)
NPPDFpol1.1 ($Q^2=10$ GeV 2)			-0.10 (8)	0.76 (4)	-0.41 (4)	1.17 (6)	0.25 (10)
DSSV ($Q^2=10$ GeV 2)			-0.012 +(56)-(62)	0.793 +(28)-(34)	-0.416 +(35)-(25)	1.209 +(45)-(42)	0.366 +(62)-(42)

PNDME, $N_F=2+1$, Clover fermion, multiple ensembles, $m_\pi = 315 - 135$ MeV

C. Alexandrou et al., $N_F=2$, twisted mass fermion, , $m_\pi = 131$ MeV, one lattice

χ QCD , $N_F=2+1$, Overlap fermion, , $m_\pi = 170, 290, 330$ MeV, 5 - 6 valence quarks for each of the three lattices → non-perturbative renormalization and normalization with anomalous Ward identity

Expt. $g_A^3 = 1.2723(23)$; CalLat: $g_A^3 = 1.271(13)$

Momenta and Angular Momenta of Quarks and Glue

- Energy momentum tensor operators decomposed in quark and glue parts gauge invariantly --- Xiangdong Ji (1997)

$$T_{\mu\nu}^q = \frac{i}{4} \left[\bar{\psi} \gamma_\mu \vec{D}_\nu \psi + (\mu \leftrightarrow \nu) \right] \rightarrow \vec{J}_q = \int d^3x \left[\frac{1}{2} \bar{\psi} \vec{\gamma} \gamma_5 \psi + \vec{x} \times \bar{\psi} \gamma_4 (-i \vec{D}) \psi \right]$$

$$T_{\mu\nu}^g = F_{\mu\lambda} F_{\lambda\nu} - \frac{1}{4} \delta_{\mu\nu} F^2 \rightarrow \vec{J}_g = \int d^3x \left[\vec{x} \times (\vec{E} \times \vec{B}) \right]$$

- Nucleon form factors

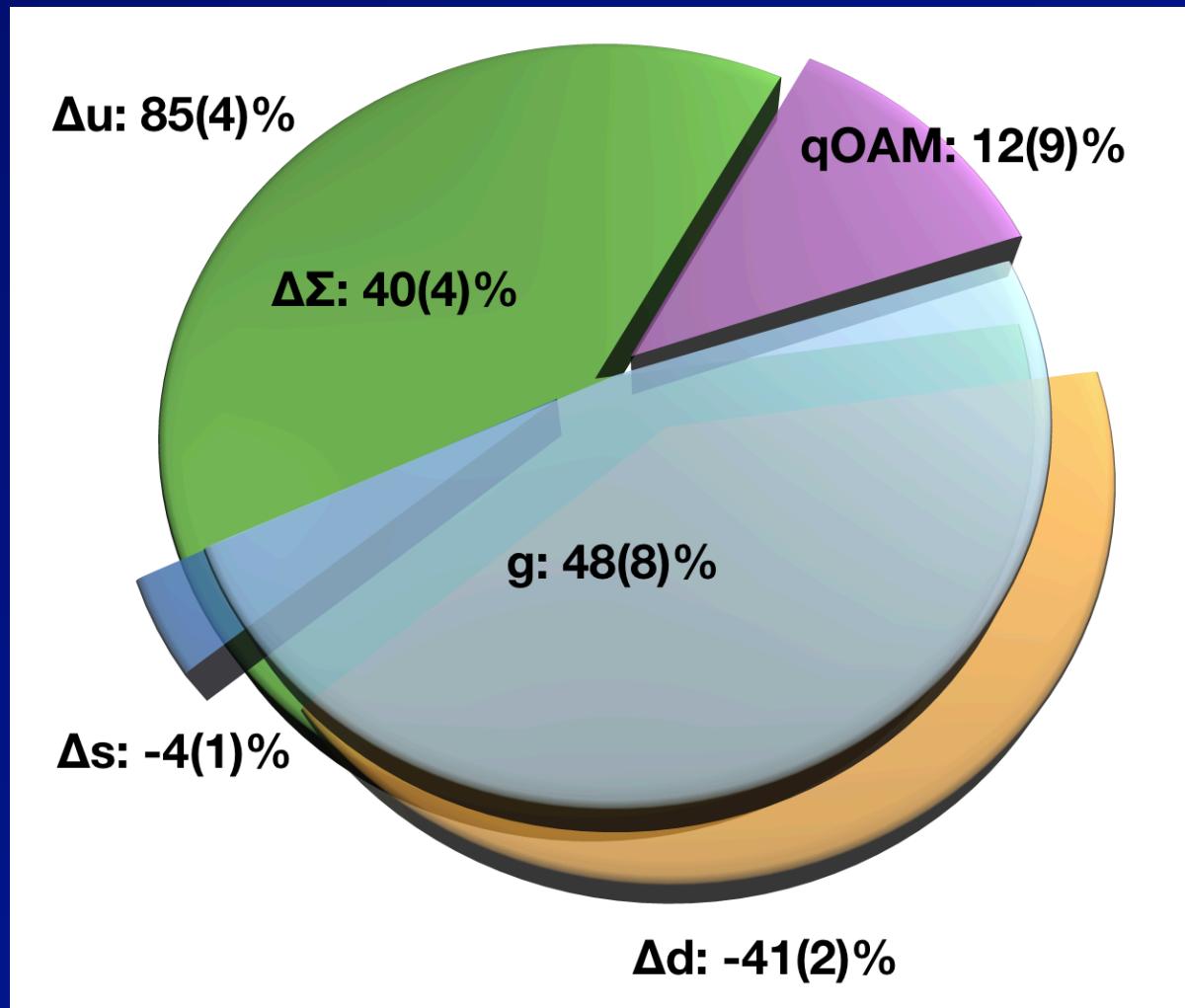
$$\begin{aligned} \langle p, s | T_{\mu\nu} | p' s' \rangle &= \bar{u}(p, s) [T_1(q^2) \gamma_\mu \bar{p}_\nu - T_2(q^2) \bar{p}_\mu \sigma_{\nu\alpha} q_\alpha / 2m \\ &\quad - i T_3(q^2) (q_\mu q_\nu - \delta_{\mu\nu} q^2) / m + T_4(q^2) \delta_{\mu\nu} m / 2] u(p' s') \end{aligned}$$

- Momentum and Angular Momentum

$$Z_{q,g} T_1(0)_{q,g} \left[\text{OPE} \right] \rightarrow \langle x \rangle_{q/g}(\mu, \bar{MS}), \quad Z_{q,g} \left[\frac{T_1(0) + T_2(0)}{2} \right]_{q,g} \rightarrow J_{q/g}(\mu, \bar{MS})$$

Proton Spin Decomposition (2+1 Flavor)

χ QCD Preliminary



Approximate by setting $T_2 = 0$

Symbols for Proton Spin Components



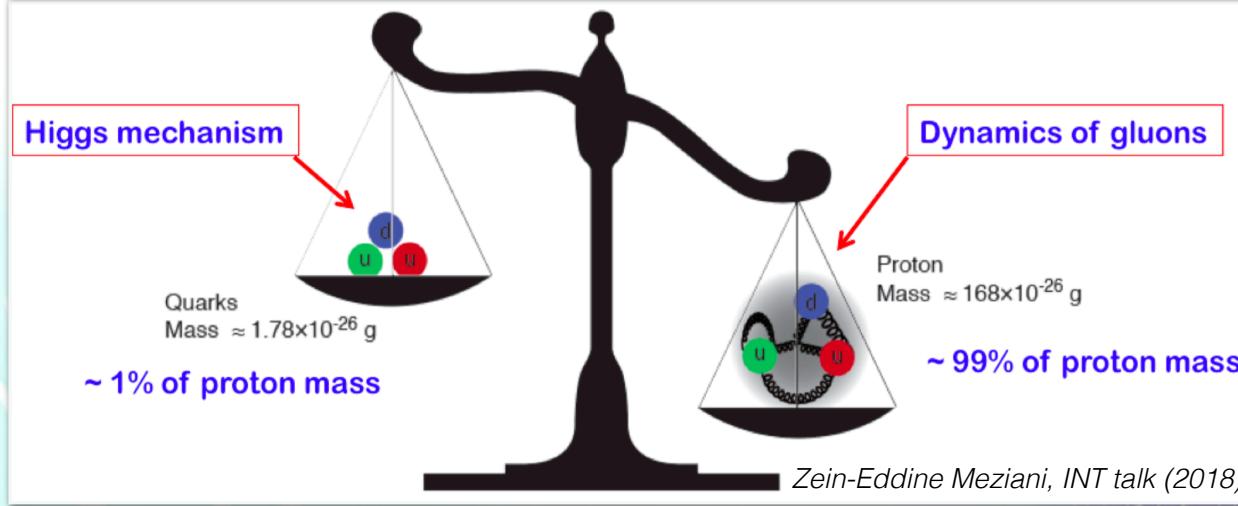
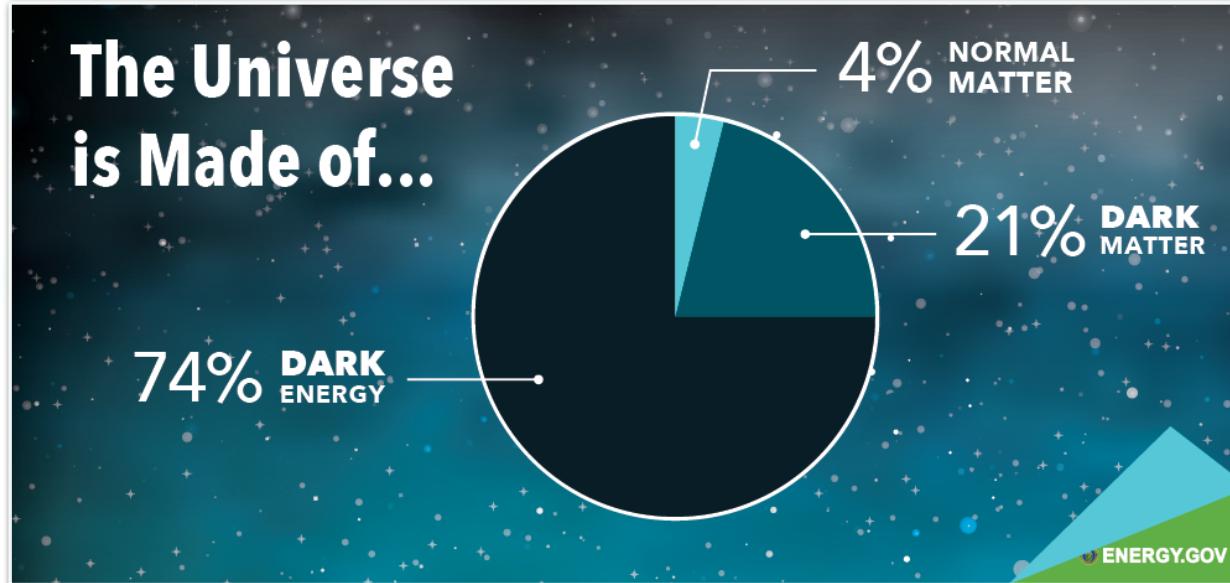
Skyrミon, a topological soliton



Trinacria – symbol of Sicily

Where does the proton mass come from?

Proton mass



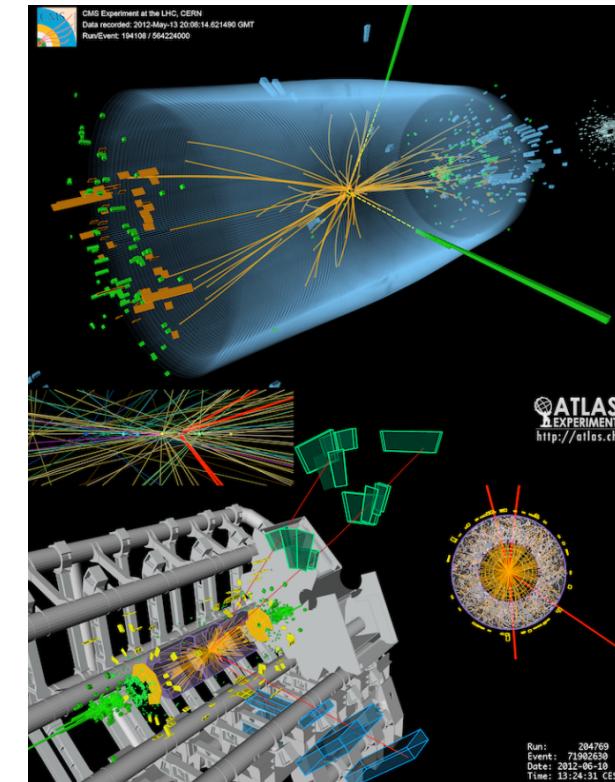
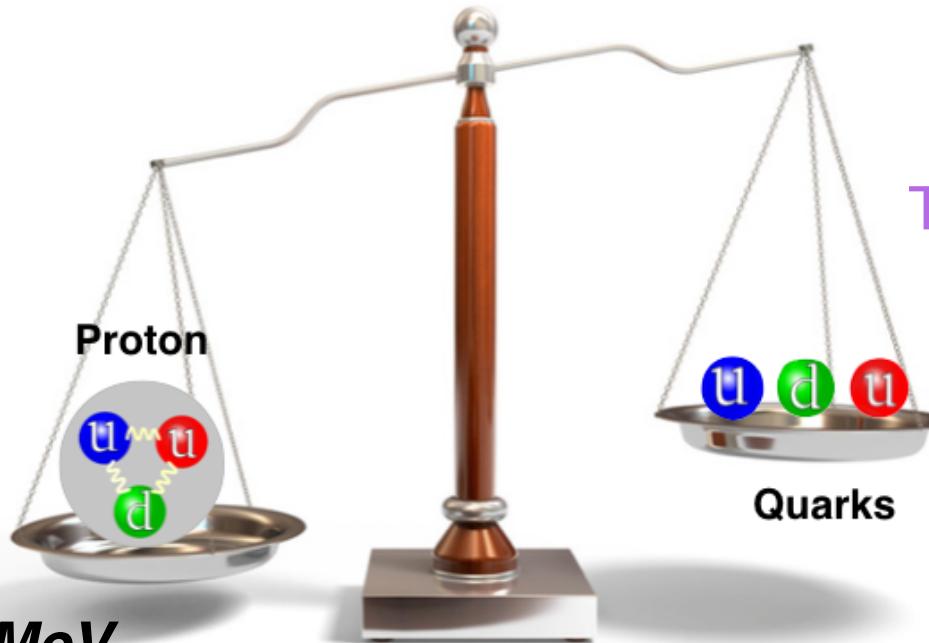
Motivation

Where does the proton mass come from, and how ?

But the mass of the proton is

938.272046(21) MeV.

~100 times of the sum of the quark masses!



The Higgs boson make the u/d quark having masses (2GeV MS-bar):

$$m_u = 2.08(9) \text{ MeV}$$
$$m_d = 4.73(12) \text{ MeV}$$

Laiho, Lunghi, & Van de Water,
Phys.Rev.D81:034503,2010

Quark and Glue Components of Hadron Mass

- Energy momentum tensor

$$T_{\mu\nu} = \frac{1}{4} \bar{\psi} \gamma_{(\mu} \vec{D}_{\nu)} \psi + G_{\mu\alpha} G_{\nu\alpha} - \frac{1}{4} \delta_{\mu\nu} G^2 \quad \langle P | T_{\mu\nu} | P \rangle = P_\mu P_\nu / M$$

- Trace anomaly

$$T_{\mu\mu} = -m(1+\gamma_m) \bar{\psi} \psi + \frac{\beta(g)}{2g} G^2$$

- Separate into traceless part $\bar{T}_{\mu\nu}$ and trace part $\hat{T}_{\mu\nu}$

$$\langle P | \bar{T}_{\mu\nu}^{q,g} | P \rangle = \langle x \rangle_{q,g} (\mu^2) (P_\mu P_\nu - \frac{1}{4} \delta_{\mu\nu} P^2) / M, \quad \langle x \rangle_q (\mu^2) + \langle x \rangle_g (\mu^2) = 1$$

$$\langle \bar{T}_{44} \rangle = -3/4M; \quad \langle \hat{T}_{\mu\mu} \rangle = -M$$

Proton mass decomposition

$$M = -\langle T_{44} \rangle = \langle H_m \rangle + \langle H_E \rangle(\mu) + \langle H_g \rangle(\mu) + \frac{1}{4}\langle H_a \rangle$$

$$M = -\langle \hat{T}_{\mu\mu} \rangle = \langle H_m \rangle + \langle H_a \rangle$$

X. Ji, PRL74:1071 (1995)

quark mass

$$\langle H_m \rangle = \sum_{u,d,s,\dots} \int d^3x m \bar{\psi} \psi$$

quark energy

$$\langle H_E \rangle = \frac{3}{4} \left(\langle x \rangle_q M - \langle H_m \rangle \right)$$

glue energy

$$\langle H_g \rangle = \frac{3}{4} \langle x \rangle_g M$$

anomaly

$$\langle H_a \rangle = \langle H_g^a \rangle + \langle H_m^\gamma \rangle$$

Ingredients

- ◆ proton mass
- ◆ scalar charge
- ◆ momentum fractions (both quark and glue)
- ◆ renormalization of momentum fractions including mixing

$$\langle x \rangle_{q,g} = \int_0^1 dx x f_{q,g}(x) = -\frac{\langle N | \frac{4}{3} \bar{T}_{44}^{q,g} | N \rangle}{M \langle N | N \rangle}$$

$$\bar{T}_{44}^q = \int d^3x \bar{\psi} \frac{1}{2} \left(\gamma_4 \overleftrightarrow{D}_4 - \frac{1}{4} \sum_{i=0,1,2,3} \gamma_i \overleftrightarrow{D}_i \right) \psi$$

$$\bar{T}_{44}^g = \int d^3x \frac{1}{2} (E^2 - B^2)$$

$$\langle H_g^a \rangle = \int d^3x \frac{-\beta(g)}{g} (E^2 + B^2)$$

$$\langle H_m^\gamma \rangle = \sum_{u,d,s,\dots} \int d^3x \gamma_m m \bar{\psi} \psi$$

Non-perturbative Renormalization and Mixing

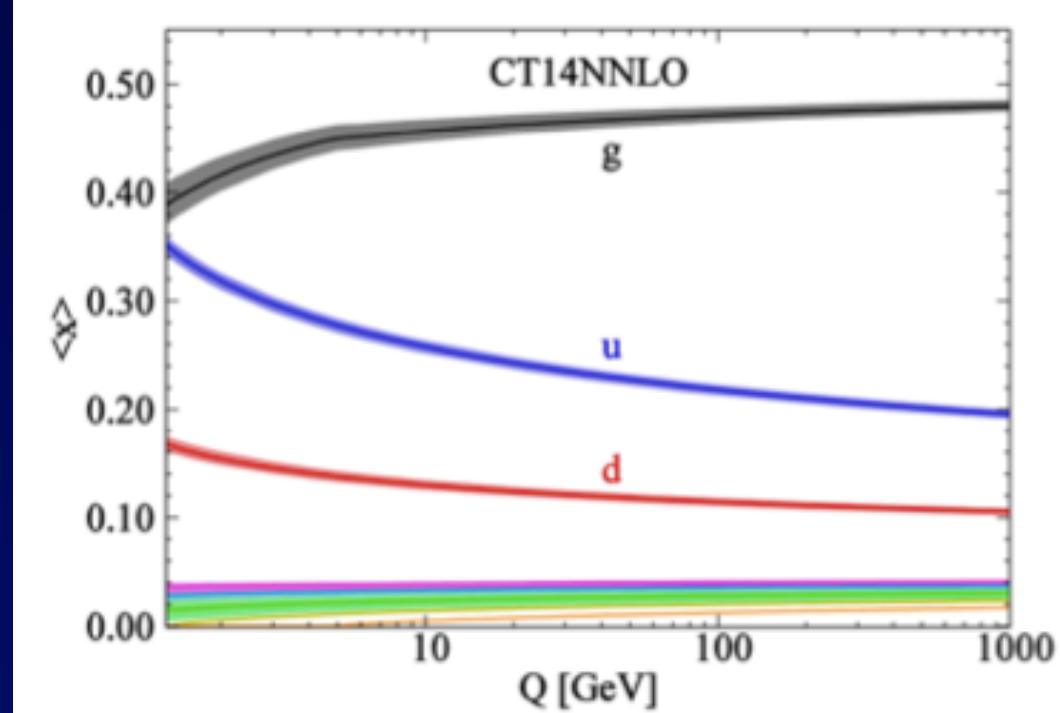
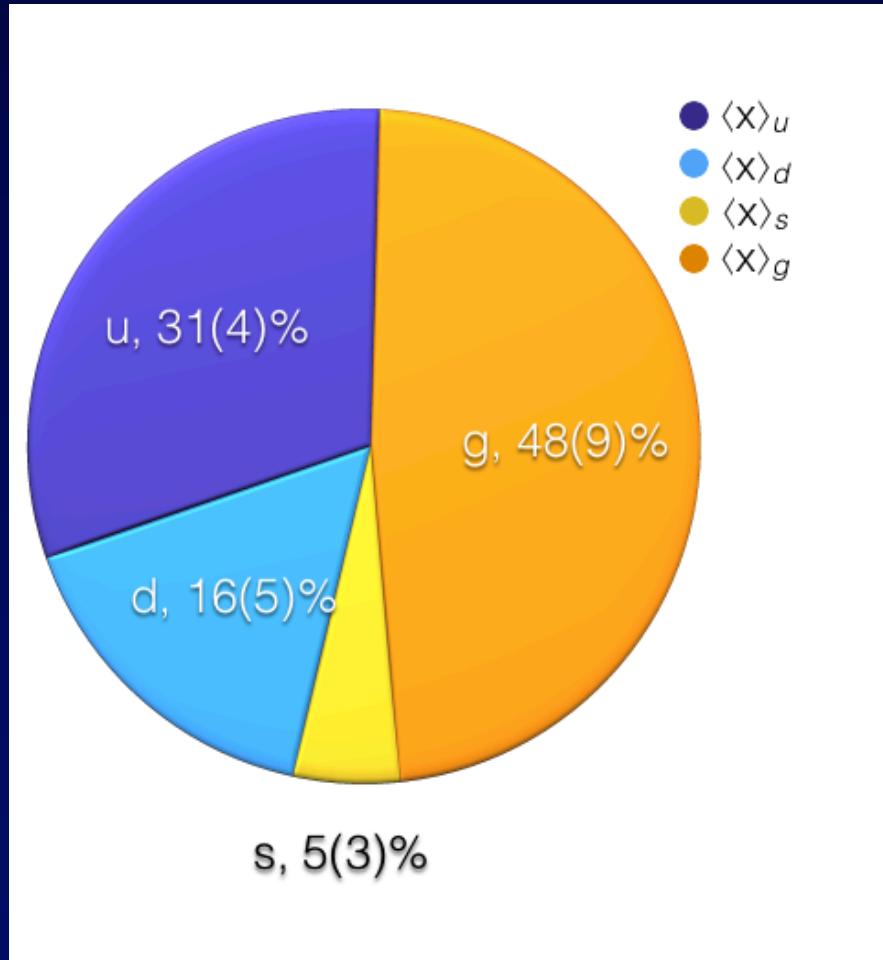
- Renormalized $\langle x \rangle_q$ and $\langle x \rangle_g$ in MS-bar at μ

$$\langle x \rangle_{u,d,s}^R = Z_{QQ}^{\overline{\text{MS}}}\langle x \rangle_{u,d,s} + \delta Z_{QQ}^{\overline{\text{MS}}}(\mu) \sum_{q=u,d,s} \langle x \rangle_q + Z_{QG}^{\overline{\text{MS}}}(\mu) \langle x \rangle_g, \quad \langle x \rangle_g^R = Z_{GQ}^{\overline{\text{MS}}}(\mu) \sum_{q=u,d,s} \langle x \rangle_q + Z_{GG}^{\overline{\text{MS}}} \langle x \rangle_g,$$

$$\begin{pmatrix} Z_{QQ}^{\overline{\text{MS}}}(\mu) + N_f \delta Z_{QQ}^{\overline{\text{MS}}}(\mu) & N_f Z_{QG}^{\overline{\text{MS}}}(\mu) \\ Z_{GQ}^{\overline{\text{MS}}}(\mu) & Z_{GG}^{\overline{\text{MS}}}(\mu) \end{pmatrix} \equiv \left\{ \left[\begin{pmatrix} Z_{QQ}(\mu_R) + N_f \delta Z_{QQ}(\mu_R) & N_f Z_{QG}(\mu_R) \\ Z_{GQ}(\mu_R) & Z_{GG}(\mu_R) \end{pmatrix} \right. \right. \\ \left. \left. \begin{pmatrix} R_{QQ}(\frac{\mu}{\mu_R}) + \mathcal{O}(N_f \alpha_s^2) & N_f R_{QG}(\frac{\mu}{\mu_R}) \\ R_{GQ}(\frac{\mu}{\mu_R}) & R_{GG}(\frac{\mu}{\mu_R}) \end{pmatrix} \right] \Big|_{a^2 \mu_R^2 \rightarrow 0} \right\}^{-1}$$

- Renormalization of glue operator in gluon propagator is very noisy \rightarrow Cluster Decomposition Error Reduction (CDER)

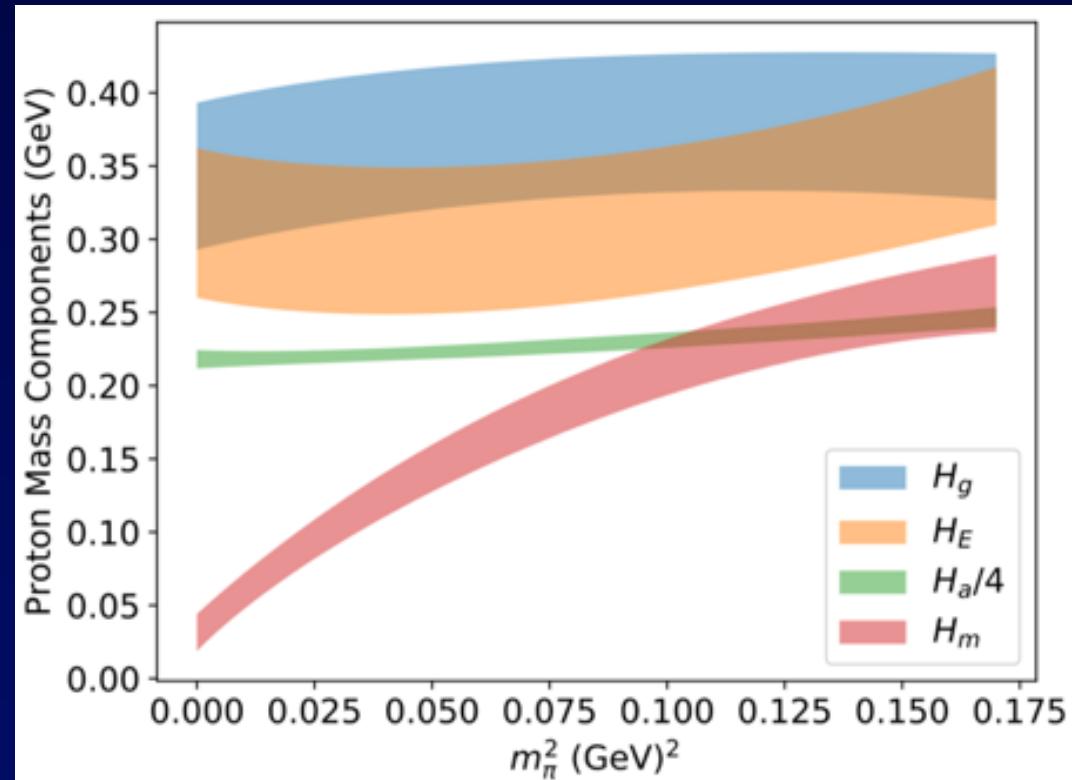
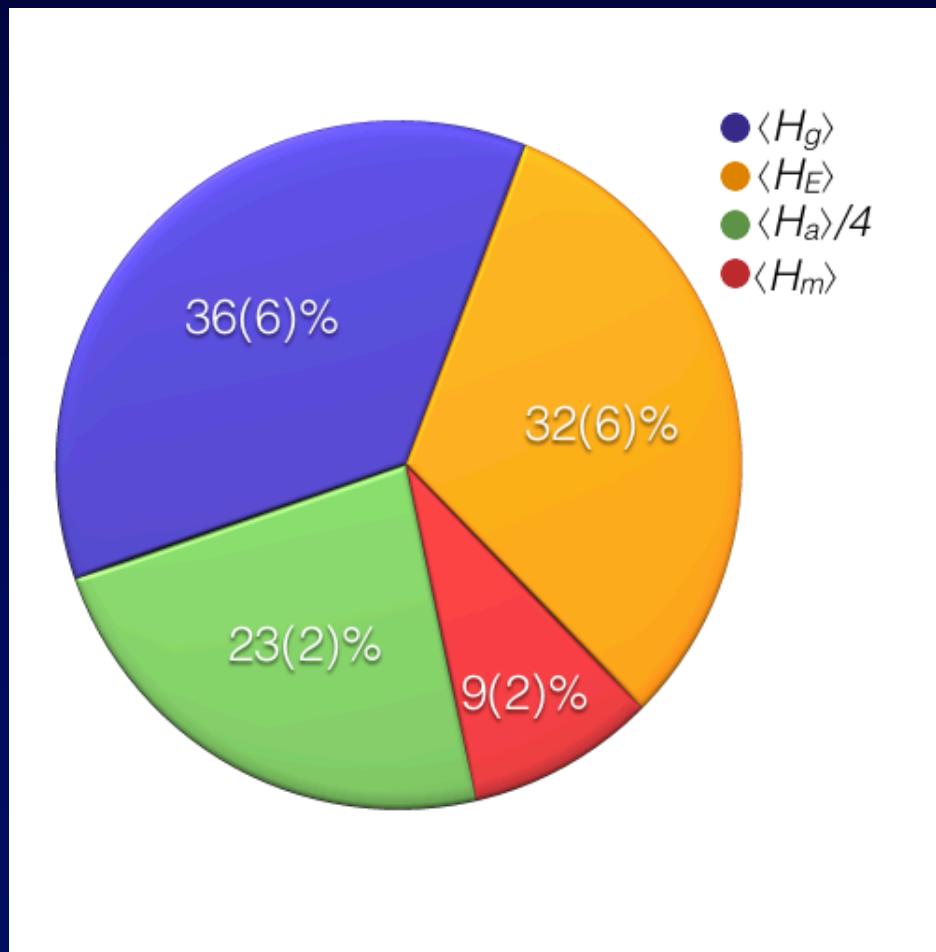
Comparison with Global Fitting of $\langle x \rangle$ MS-bar at 2 GeV



S. Dulat et al, Phys. Rev. D 93 (2016), 033006

χ QCD, preliminary

Proton Mass Decomposition



Y.B. Yang et al (χ QCD), PRL 121, 212001 (2018)

PRL as Editor's Suggestion

(Y.B. Yang, PRL 121, 212001 (2018))



11/26/2018

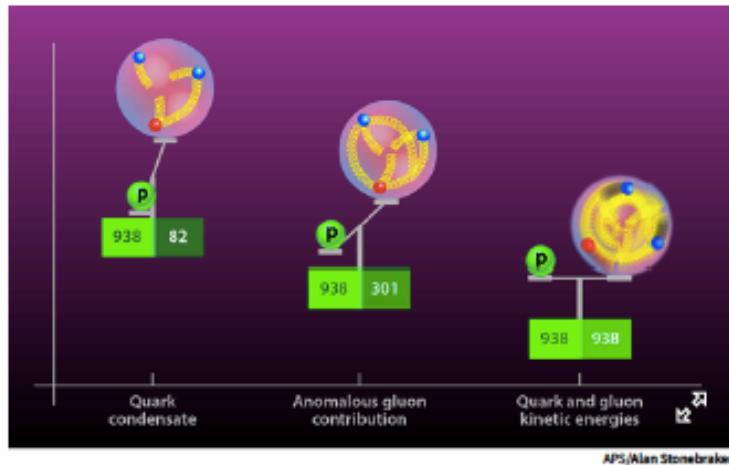
Physicists finally calculated where the proton's mass comes from | Science News

Viewpoint: Dissecting the Mass of the Proton

André Walker-Loud, Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, CA, USA

November 19, 2018 • Physics 11, 118

A calculation determines four distinct contributions to the proton mass, more than 90% of which arises entirely from the dynamics of quarks and gluons.



APS/Alan Stonebraker

Figure 1: The proton is comprised of two up quarks and one down quark, but the sum of these quark masses is a mere 1% of the proton mass. Using lattice QCD, Yang and colleagues determined the relative contributions of the four sources of the proton mass [1]. ... [Show more](#)

ScienceNews

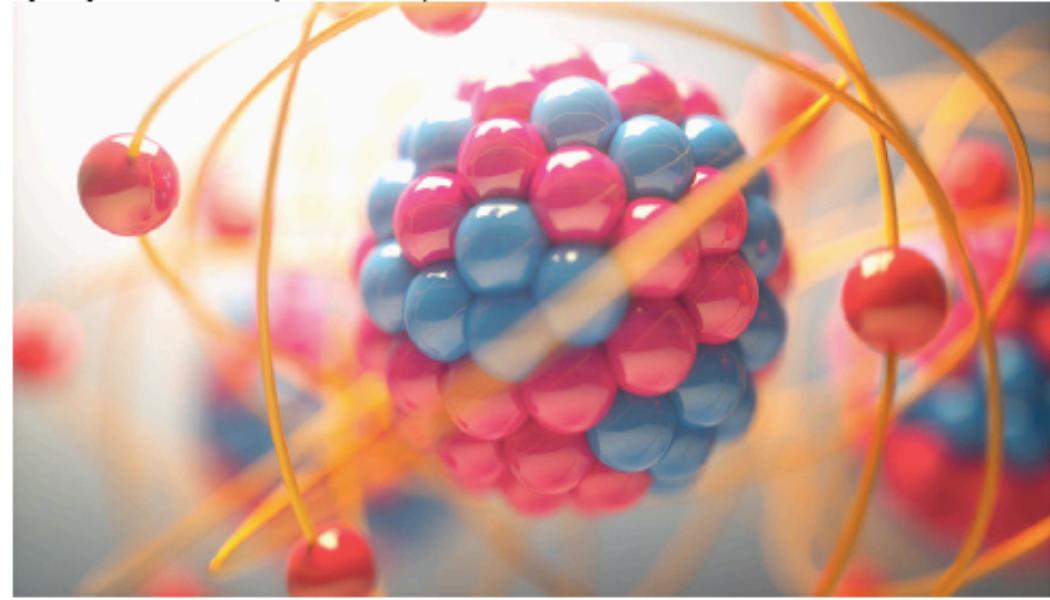
MAGAZINE OF THE SOCIETY FOR SCIENCE & THE PUBLIC

News: Particle Physics

Physicists finally calculated where the proton's mass comes from

Only 9 percent of the subatomic particle's bulk comes from the mass of its quarks

By Emily Conover 6:00am, November 26, 2018



MASSIVE UNDERTAKING Using a technique called lattice QCD, scientists figured out how protons (illustrated here in the nucleus of an atom) get their mass.

ktsdesign/Shutterstock

Hadronic Tensor in Euclidean Path-Integral Formalism

- Deep inelastic scattering
In Minkowski space

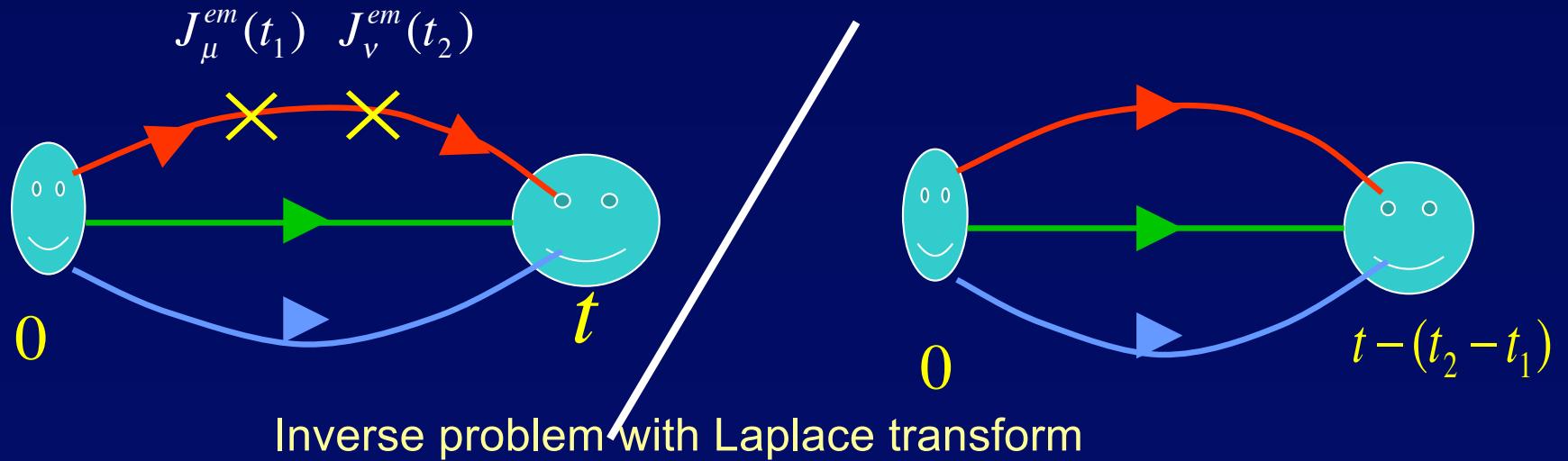
$$\frac{d^2\sigma}{dE'd\Omega} = \frac{\alpha^2}{q^4} \left(\frac{E'}{E}\right) l^{\mu\nu} W_{\mu\nu}$$

$$W_{\mu\nu}(\vec{q}, \vec{p}, v) = \frac{1}{\pi} \text{Im} T_{\mu\nu} = \langle N(\vec{p}) | \int \frac{d^4x}{4\pi} e^{iq\cdot x} J_\mu(x) J_\nu(0) | N(\vec{p}) \rangle_{\text{spin avg}}$$

$$= \frac{1}{2} \sum_n \int \prod_{i=1}^n \left[\frac{d^3 p_i}{(2\pi)^3 2E_{pi}} \right] (2\pi)^3 \delta^4(p_n - p - q) \langle N(\vec{p}) | J_\mu | n \rangle \langle n | J_\nu | N(\vec{p}) \rangle_{\text{spin avg}}$$

- Euclidean path-integral

KFL and S.J. Dong, PRL 72, 1790 (1994)
KFL, PRD 62, 074501 (2000)



Hadronic tensor on the lattice

four-point function with **3-dimensional Fourier transform**

$$C_4 = \sum_{\vec{x}_f} e^{-i\vec{p} \cdot \vec{x}_f} \sum_{\vec{x}_2 \vec{x}_1} e^{-i\vec{q} \cdot (\vec{x}_2 - \vec{x}_1)} \left\langle \chi_N(\vec{x}_f, t_f) J_\mu^\dagger(\vec{x}_2, t_2) J_\nu(\vec{x}_1, t_1) \bar{\chi}_N(\vec{0}, t_0) \right\rangle$$

Euclidean hadronic tensor defined as a function of time difference between the currents

$$\begin{aligned} \tilde{W}_{\mu\nu}(p, \vec{q}, \tau) &= \frac{E_p}{m_N} \frac{\text{Tr}[\Gamma_e C_4]}{\text{Tr}[\Gamma_e C_2]} \rightarrow \sum_{\vec{x}_2 \vec{x}_1} e^{-i\vec{q} \cdot (\vec{x}_2 - \vec{x}_1)} \langle p, s | J_\mu(\vec{x}_2, t_2) J_\nu(\vec{x}_1, t_1) | p, s \rangle \\ &= \sum_n A_n e^{-(E_n - E_p)\tau}, \quad \tau \equiv t_2 - t_1 \end{aligned}$$

Solving the **inverse problem** of a Laplace transform to get back to Minkowski space

$$\tilde{W}_{\mu\nu}(p, \vec{q}, \tau) = \int d\nu W_{\mu\nu}(p, \vec{q}, \nu) e^{-\nu\tau}$$

K.F. Liu and S. J. Dong, PRL 72, 1790 (1994)

K.-F. Liu, PRD 62, 074501 (2000)

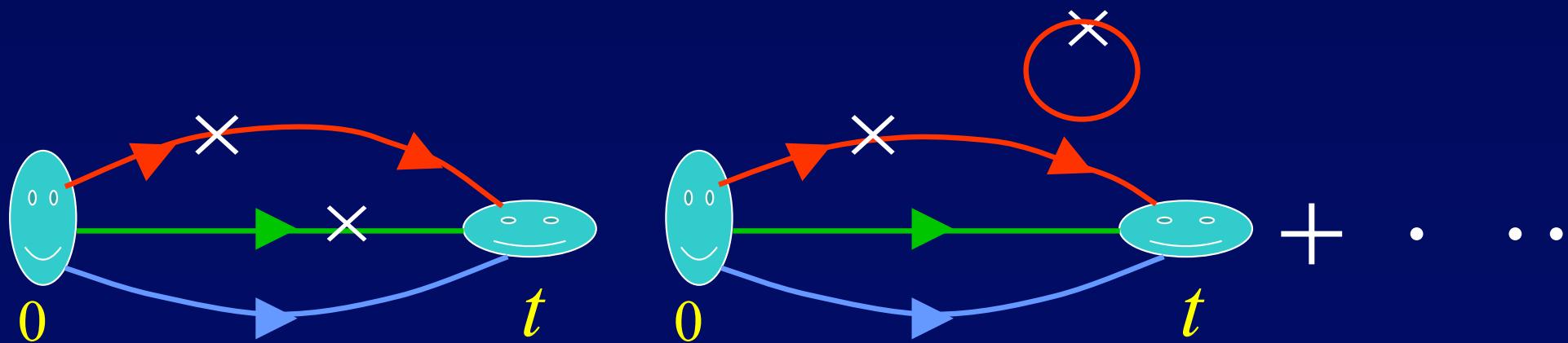
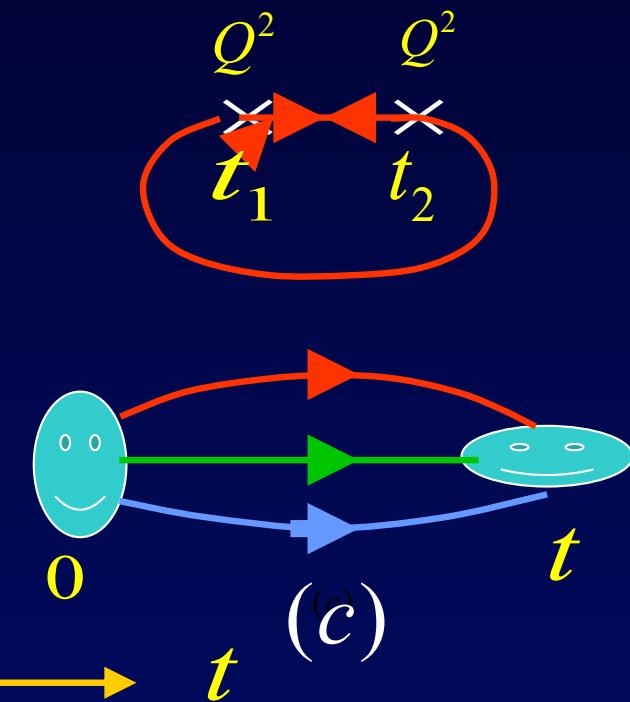
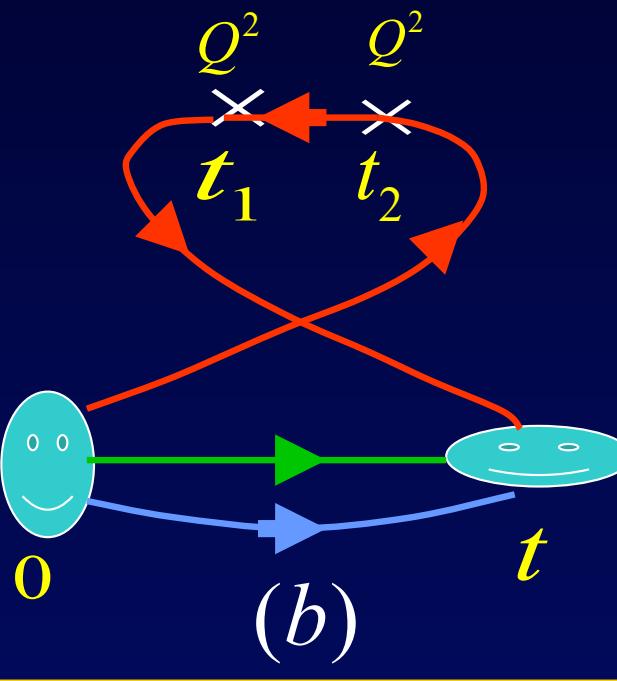
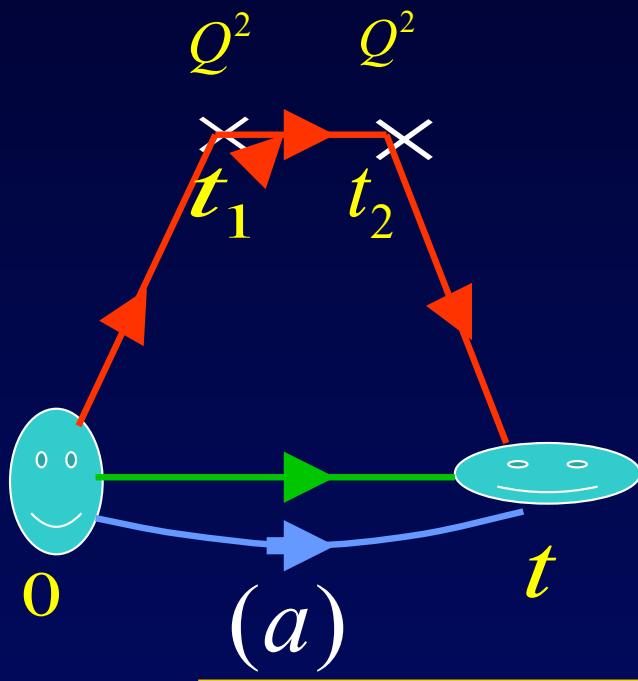
J. Liang et. al., EPJ Web Conf. 175, 14014 (2018)

J. Liang et. al., arXiv:1906.05312

$$q = q_V + q_{CS}$$

$$\bar{q}_{CS}$$

$$q_{DS} = (\neq ?) \bar{q}_{DS}$$

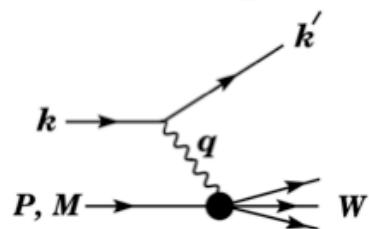


Cat's ears diagrams are suppressed by $O(1/Q^2)$.

Inverse problems are ubiquitous

- ◆ Extracting spectral functions from lattice data: $c_2(t) = \int d\omega e^{-\omega t} \rho(\omega)$
- ◆ Global fittings of PDFs: $F_i = \sum_a C_i^a \otimes f_a$ *Y. Burnier and A. Rothkopf, PRL 111, 182003 (2013)*
- ◆ Lattice calculation of Quasi-PDFs: $\tilde{q}(x, P_3) = \frac{2P_3}{4\pi} \sum_{z=-z_{\max}}^{z_{\max}} e^{-ixP_3 z} h_\Gamma(P_3, z)$ *J. Karpie et. al., JHEP04, 057 (2019)*
$$\tilde{q}(x, \mu^2, P^z) = \int_0^1 \frac{dy}{y} Z\left(\frac{x}{y}, \frac{\mu}{P^z}\right) q(y, \mu^2) + \mathcal{O}\left(\Lambda^2/(P^z)^2, M^2/(P^z)^2\right)$$
 X. Xiong et. al., PRD90:014051 (2014)
- ◆ Lattice cross sections: $\sigma_n(\omega, \xi^2, P^2) = \sum_a \int_{-1}^1 \frac{dx}{x} f_a(x, \mu^2) \times K_n^a(x\omega, \xi^2, x^2 P^2, \mu^2) + O(\xi^2 \Lambda_{\text{QCD}}^2)$ *Y.-Q. Ma and J.-W. Qiu, PRL 120, 022003 (2018)*
- ◆ Lattice calculation of Pseudo-PDFs: $\mathfrak{M}_R(\nu, \mu^2) \equiv \int_0^1 dx \cos(\nu x) q_v(x, \mu^2)$ *K. Orginos et al., PRD96, 094503 (2017)*

Sketch the hadronic tensor

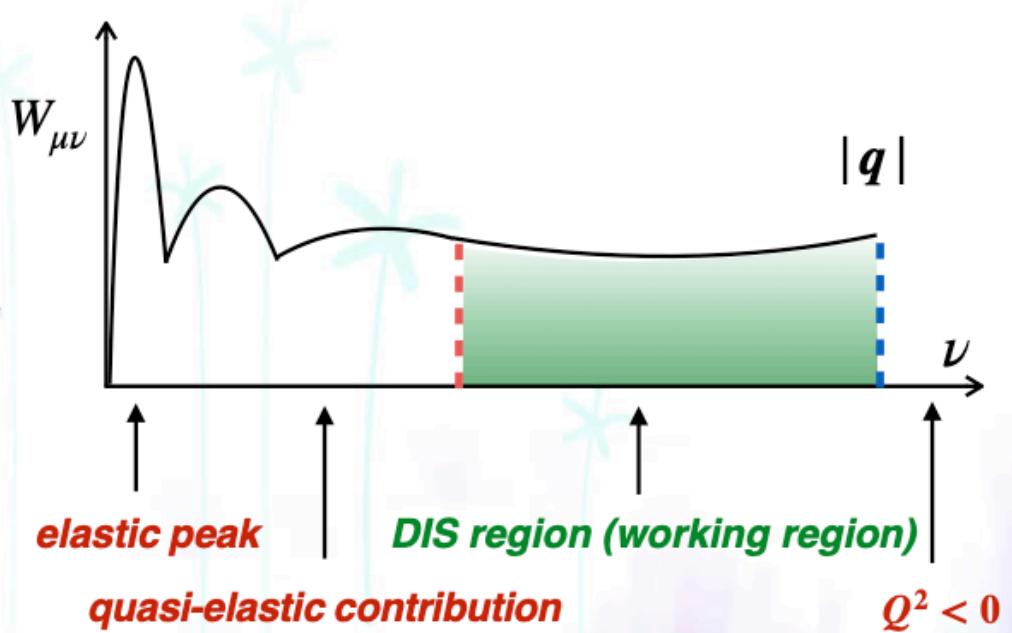
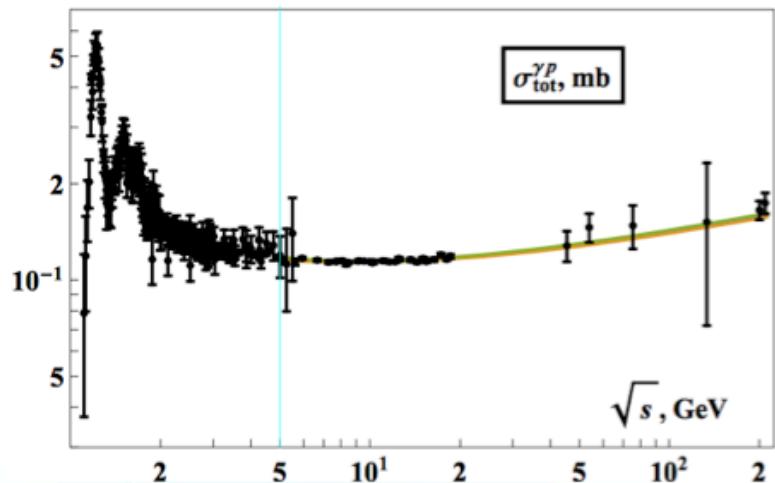


$$Q^2 = -q^2 = \mathbf{q}^2 - \nu^2$$

$$x = \frac{Q^2}{2m\nu}$$

fixed by momentum setup

comes from solving the inverse problem



$\nu > (E_{n=0} - E_p) + \Delta E$ (away from the elastic peak)

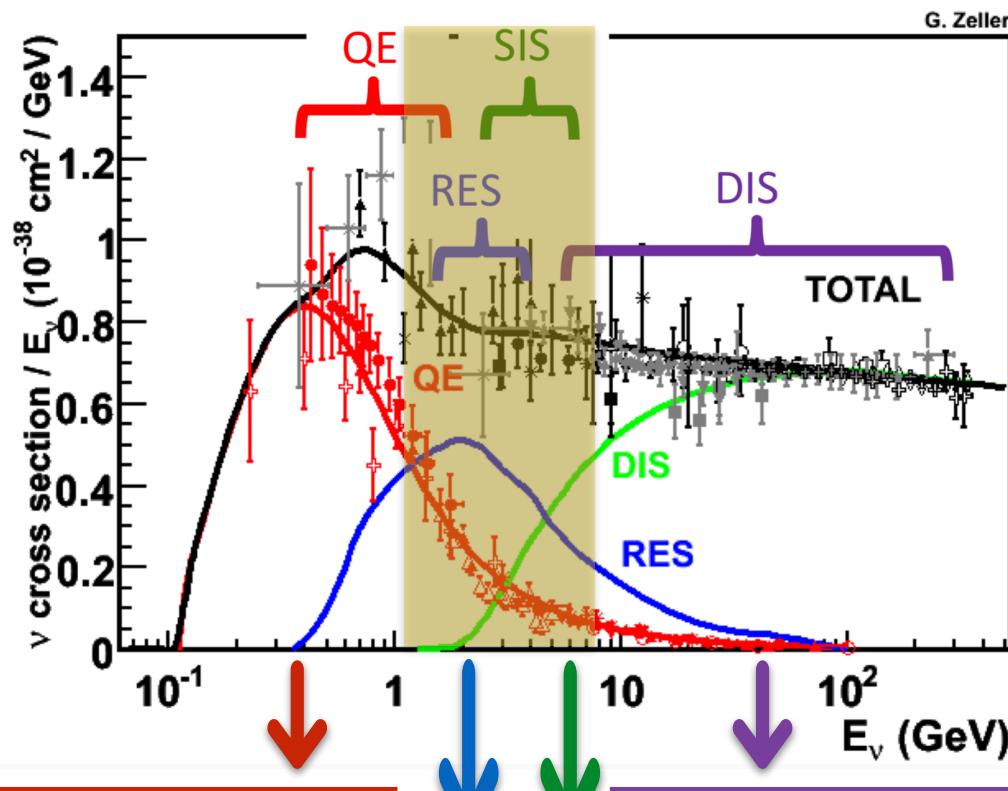
$\nu < |q|$ (physical x and Q^2)

Hadronic tensor and neutrino-nucleus scattering

- ◆ New long-baseline neutrino experiments are in preparation: T2K, NOvA, PINGU, ORCA, Hyper-Kamiokande, DUNE...

- ◆ Be
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- ◆ Ch
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elastic form factors

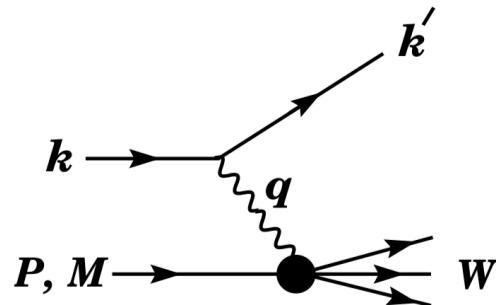
Parton distribution functions

inclusive hadronic tensor!

inclusive hadronic tensor!

J.A. Formaggio and G.P. Zeller, RMP84, 1307 (2012)

Hadronic tensor



for lepton-nucleon scatterings

$$\frac{d^2\sigma}{dxdy} = \frac{2\pi ya^2}{Q^4} \sum_j \eta_j L_j^{\mu\nu} \mathbf{W}_{\mu\nu}^j$$

$$W_{\mu\nu}(p, \vec{q}, \nu) = \frac{1}{4\pi} \int d^4z e^{iq\cdot z} \left\langle p, s \left| \left[J_\mu^\dagger(z) J_\nu(0) \right] \right| p, s \right\rangle \sim \text{Im} [T_{\mu\nu}]$$

$$W_{\mu\nu} = \left(-g_{\mu\nu} + \frac{q_\mu q_\nu}{q^2} \right) F_1(x, Q^2) + \frac{\hat{P}_\mu \hat{P}_\nu}{P \cdot q} F_2(x, Q^2)$$

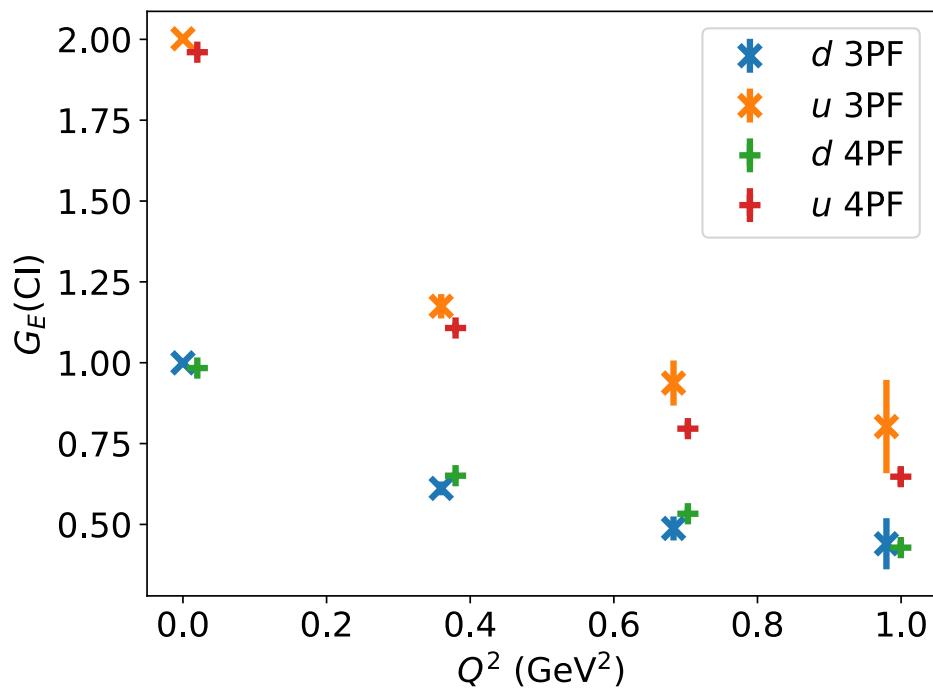
The hadronic tensor and structure functions encode the non-perturbative nature of the nucleon.

- ◆ Elastic form factors: $\mathbf{F}_2^{\text{el}} = \delta(q^2 + 2m_N\nu) \frac{2m_N}{1 - q^2/4m_N^2} \left(\mathbf{G}_E^2(\mathbf{q}^2) - \frac{q^2}{4M_N^2} \mathbf{G}_M^2(\mathbf{q}^2) \right)$
- ◆ Neutrino-nucleon scattering
- ◆ PDFs from DIS: $\mathbf{F}_i = \sum_a C_i^a \otimes \mathbf{f}_a$

Elastic FFs

$$\tilde{W}_{44}(\vec{p}, \vec{q}, \tau) = \sum_{\vec{x}_2 \vec{x}_1} e^{-i\vec{q} \cdot (\vec{x}_2 - \vec{x}_1)} \langle p, s | J_\mu(\vec{x}_2, t_2) J_\nu(\vec{x}_1, t_1) | p, s \rangle = \sum_n A_n e^{-\nu_n \tau}$$

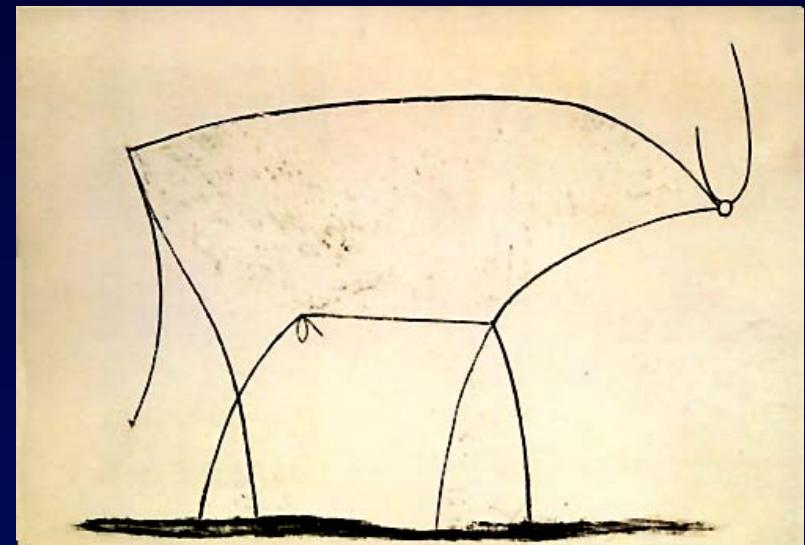
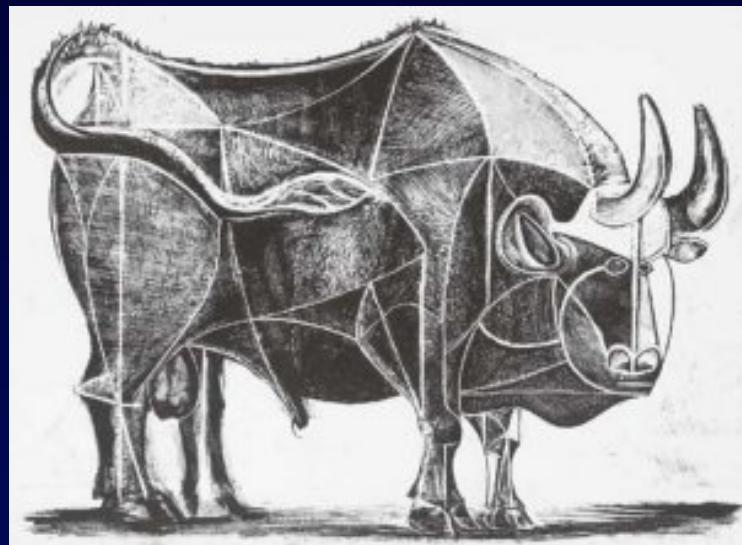
$$A_0 = \langle p, s | J_4(\vec{q}) | n=0 \rangle \langle n=0 | J_4(-\vec{q}) | p, s \rangle = G_E^2(Q^2)$$



FFs extracted from 3-point functions and 4-point functions show consistency.

(32IF lattice, $a \sim 0.063$ fm, pion mass ~ 370 MeV)

Le Taureau of Pablo Picasso (1945)



5th stage

11th stage

Dynamical chiral fermion
Physical pion mass
Continuum limit
Infinite volume limit



Quenched approximation

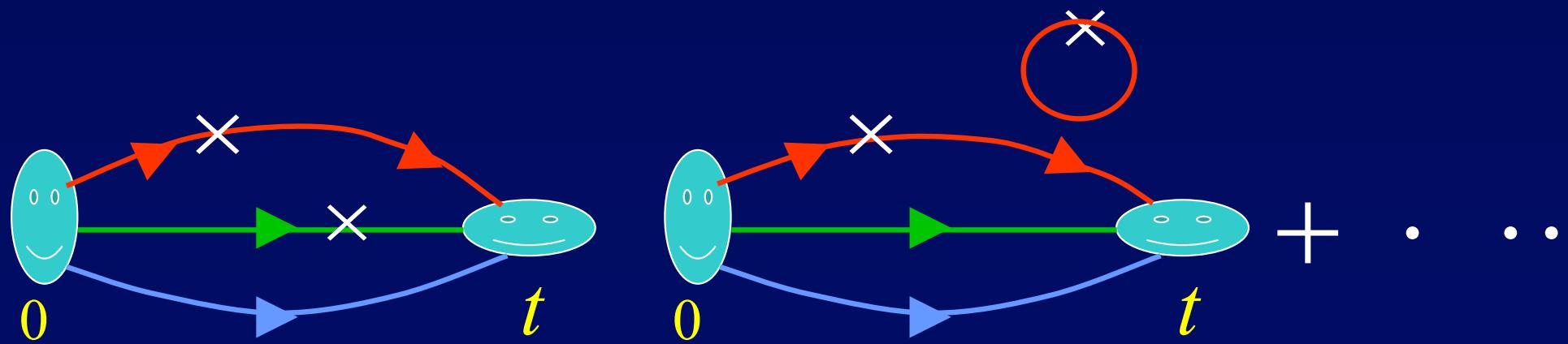
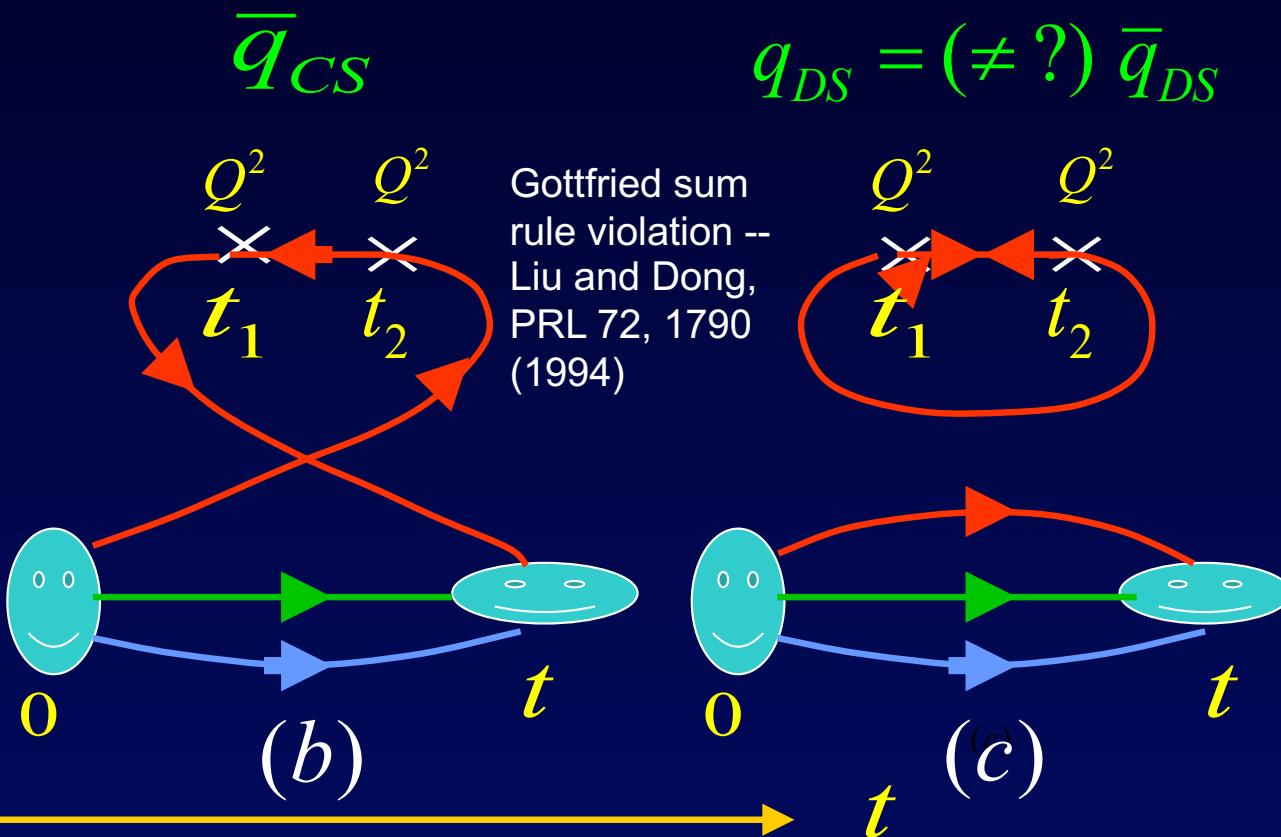
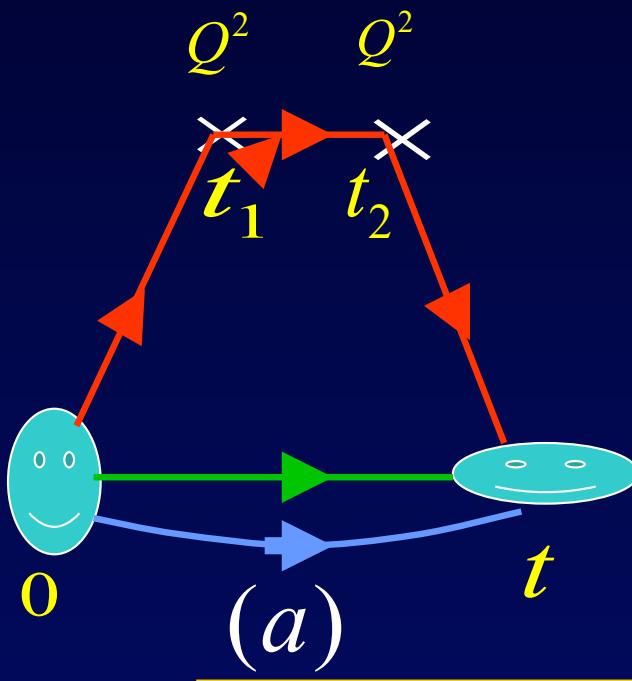
Summary and Challenges

- Together with experiments on LHC and EIC and global fitting of PDF, lattice QCD calculations of hadron structure (proton spin and mass decomposition, moments of PDFs, form factors, etc.) can advance our understanding of the nucleon properties in more detail.
- Hadronic tensor calculation from lattice QCD involves a numerically challenging inverse problem. It can address the low-energy neutrino-nucleon elastic scattering, the inelastic resonance region, the shallow inelastic as well as the deep inelastic region. The latter will require large lattices with lattice spacing as small as 0.02 fm. Together with experiments from DUNE, they can help us understand the salient properties of the neutrinos.

Imagination may not be relevant.
Knowledge cannot predict.
By simulating *imagination* and
testing against *knowledge* only can
reality be re-created.

Hadronic Tensor

$$q = q_V + q_{CS}$$



Cat's ears diagrams are suppressed by $O(1/Q^2)$.

$$q_i^- = q_i^{v+cs} - \bar{q}_i^{cs} + q_i^{ds} - \bar{q}_i^{ds} \equiv q_i^v + q_i^{ds} - \bar{q}_i^{ds}$$

Correlators

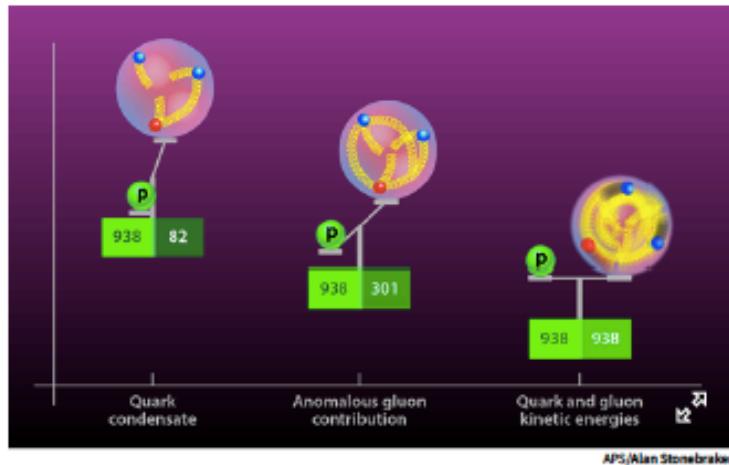
- One point collerators → $\langle \bar{\psi}\psi \rangle, \langle G_{\mu\nu} G_{\mu\nu} \rangle \dots$
- Two point correlators → hadron masses, decay constants ...
- Three point correlators → parton moments form factors, nucleon matrix elements ...
- Four point correlators → hadronic tensor, PDF, GPD, TMD, nEDM...

Viewpoint: Dissecting the Mass of the Proton

André Walker-Loud, Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, CA, USA

November 19, 2018 • Physics 11, 118

A calculation determines four distinct contributions to the proton mass, more than 90% of which arises entirely from the dynamics of quarks and gluons.



APS/Alan Stonebraker

Figure 1: The proton is comprised of two up quarks and one down quark, but the sum of these quark masses is a mere 1% of the proton mass. Using lattice QCD, Yang and colleagues determined the relative contributions of the four sources of the proton mass [1]. ... [Show more](#)

ScienceNews

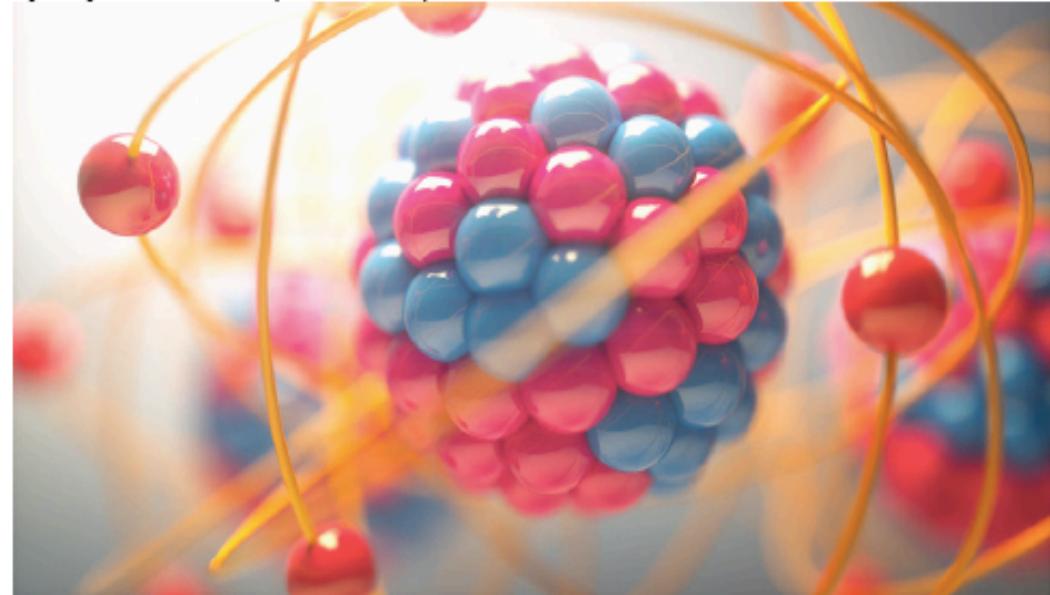
MAGAZINE OF THE SOCIETY FOR SCIENCE & THE PUBLIC

News: Particle Physics

Physicists finally calculated where the proton's mass comes from

Only 9 percent of the subatomic particle's bulk comes from the mass of its quarks

By Emily Conover 6:00am, November 26, 2018



MASSIVE UNDERTAKING Using a technique called lattice QCD, scientists figured out how protons (illustrated here in the nucleus of an atom) get their mass.

ktsdesign/Shutterstock